

Supporting Information.

Spotlight on Excitonic Coupling in Polymorphic and Textured Anilino Squaraine Thin Films

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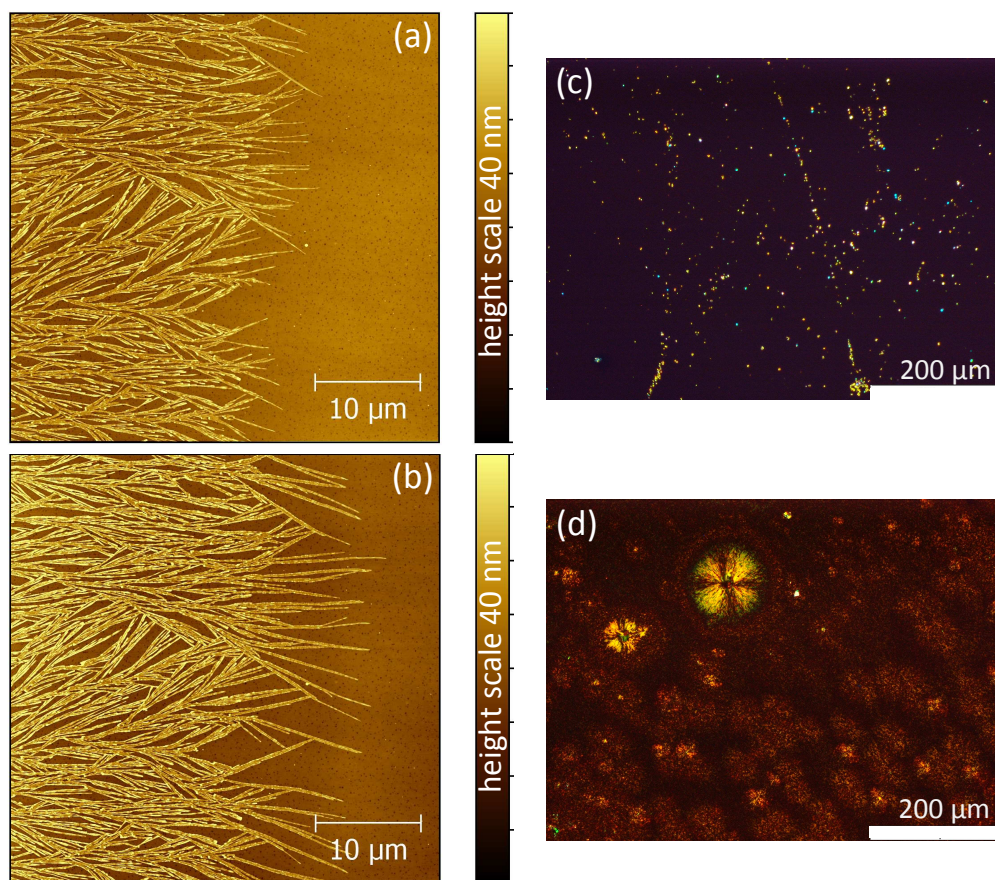


Figure S1: AFM images monitoring the growth of a sunflower in an amorphous SQIB film, time between (a) and (b) are approx. 42 hours. Optical microscopy images between crossed polarizers of (c) a fresh SQIB thin film annealed at 60 °C and (d) the same sample after two months storage.

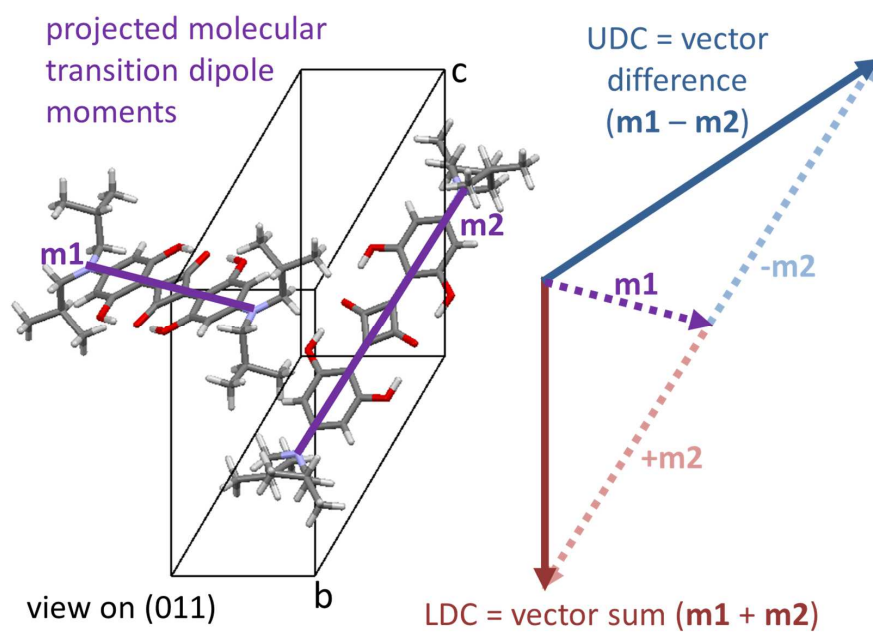


Figure S2: Graphical vector addition (red arrow, LDC) and subtraction (blue arrow, UDC) of the molecular transition dipole moments (violet color) for the monoclinic polymorph (sunflowers).

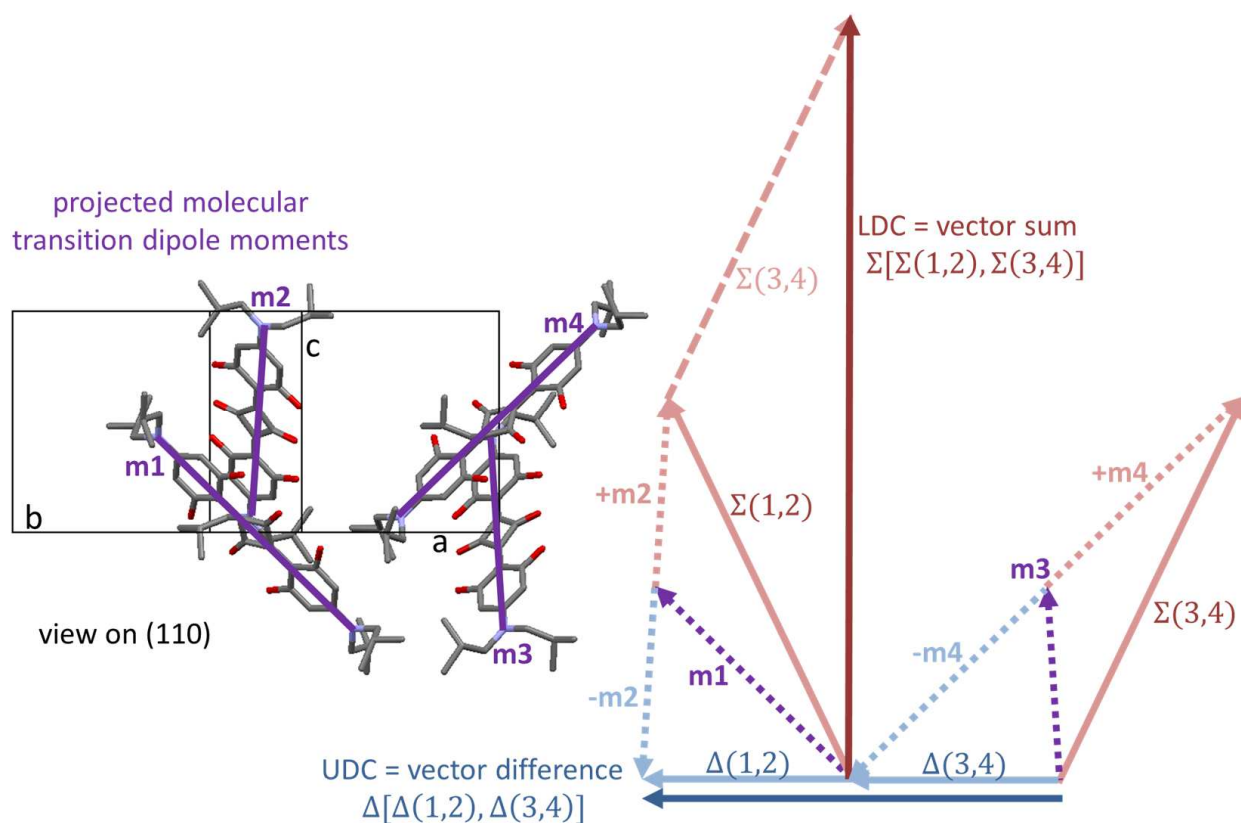


Figure S3: Graphical vector addition (red arrow, LDC) and subtraction (blue arrow, UDC) of the molecular transition dipole moments (violet color) for the orthorhombic polymorph (platelets). At first, the vector sums Σ [differences Δ] of the two pairs of stacked molecules are calculated: $\Sigma(1, 2) = m1 + m2$ [$\Delta(1, 2) = m1 - m2$] and $\Sigma(3, 4) = m3 + m4$ [$\Delta(3, 4) = m3 - m4$]. Secondly, the vector sum [difference] of the stacked pairs are graphically added [subtracted]: $\Sigma(1, 2) + \Sigma(3, 4) = \text{LDC}$ [$\Delta(1, 2) - \Delta(3, 4) = \text{UDC}$], resulting in the overall Davydov components.

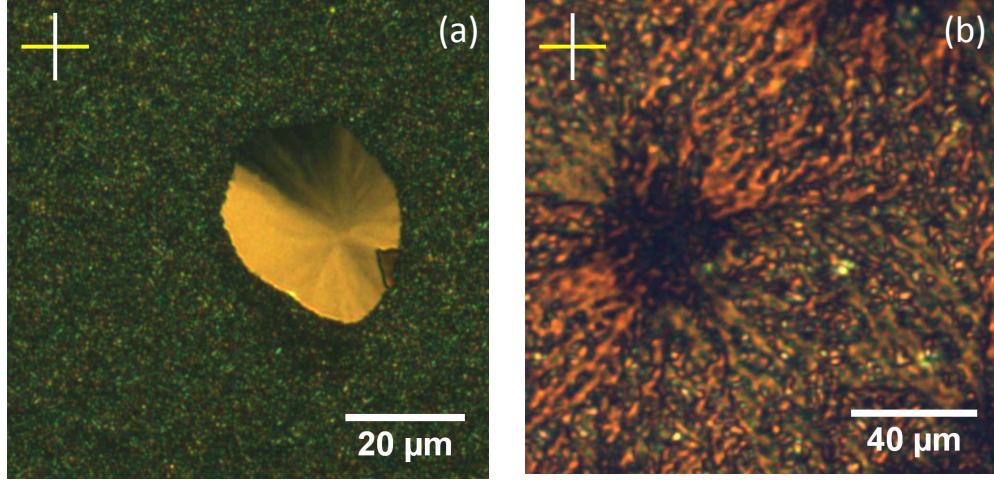


Figure S4: Birefringence of a single SQIB platelet (a) and a sunflower (b). The cross indicates the direction of polarizer (white) and analyzer (orange). Screenshots of the movie files.

In Figure S4, screenshots of two movies are displayed. In (a) a golden platelet is rotated between two crossed polarizers under a polarized light microscope. The images are observed in transmission. The polarizer and analyzer directions are depicted by a white and orange line, respectively, in the upper left. The same is shown in (b) for a sunflower. The movies are available as downloads.

A total of $N = 73$ images are recorded over a rotation angle ϕ of 360° . From the intensity variation $I_n^{x,y}$ ($n = 0 \dots N - 1$) of each pixel at image position (x, y) , the discrete Fourier coefficients $\tilde{I}_\gamma(x, y)$ are determined using

$$\tilde{I}_\gamma(x, y) = \frac{1}{N} \sum_{n=0}^{N-1} I_n^{x,y} e^{-i2\pi\gamma n/N} \quad . \quad (1)$$

From this, the extinction angle ϕ_{ext} , i.e. the angle for which the intensity is smallest, is extracted via

$$\phi_{\text{ext}}(x, y) = \frac{1}{2} \arg \left(\tilde{I}_{\gamma=4}^*(x, y) + i \left| \tilde{I}_{\gamma=4}(x, y) \right| \right) + \frac{\pi}{8} \quad . \quad (2)$$

Eventually, a value of $\pi/2$ is added to select the tangential extinction angle over the radial one.

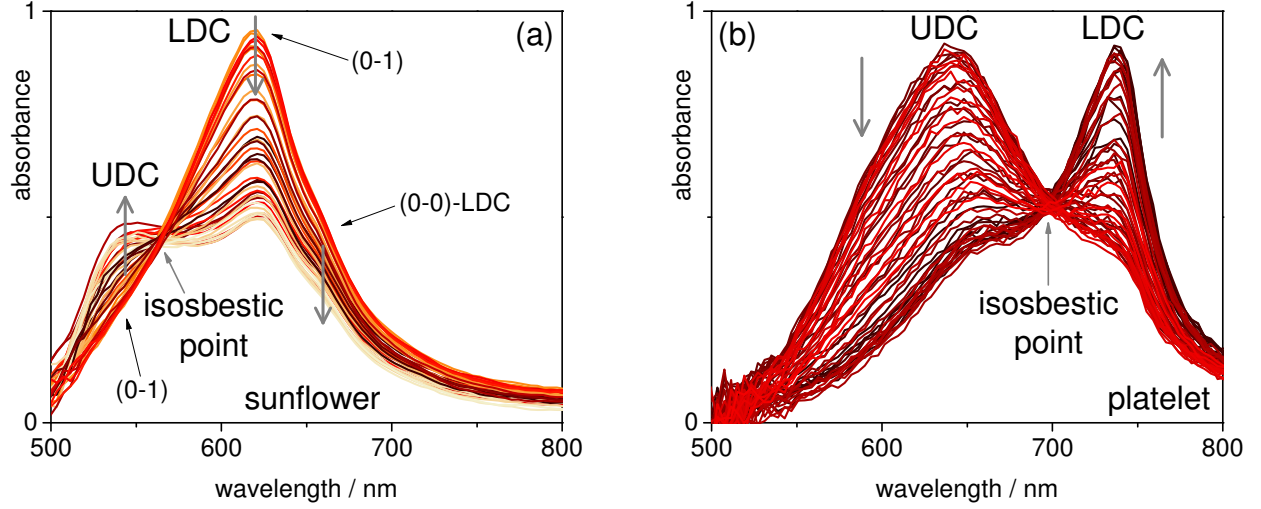


Figure S5: Absorption spectra recorded with a single polarizer dependent on the rotation angle of the polarizer (a) within a typical sunflower, and (b) on a typical platelet area. UDC = Upper Davydov Component; LDC = Lower Davydov Component. Note that for the sunflower in (a) the oscillator strength of the (0-1)-LDC might appear stronger than it actually is due to spectral overlap with the not resolved (0-0)-UDC.