

Supporting Information for “Kohn–Sham decomposition in real-time time-dependent density-functional theory: An efficient tool for analyzing plasmonic excitations”

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S1 Derivation of Eq. (9) within the PAW formalism

Within the PAW formalism [P. E. Blöchl, *Phys. Rev. B* **50**, 17953 (1994)], Eq. (1) reads

$$i\mathcal{T}^\dagger \mathcal{T} \frac{\partial}{\partial t} \tilde{\psi}_n(\mathbf{r}, t) = \mathcal{T}^\dagger H_{\text{KS}}(t) \mathcal{T} \tilde{\psi}_n(\mathbf{r}, t), \quad (\text{S1})$$

where $\tilde{\psi}_n(\mathbf{r}, t)$ is a pseudo wave function and \mathcal{T} denotes the PAW transformation operator.

In the LCAO method, the pseudo wave function $\tilde{\psi}_n(\mathbf{r}, t)$ is expanded in localized basis functions $\tilde{\phi}_\mu(\mathbf{r})$ centered at atomic coordinates

$$\tilde{\psi}_n(\mathbf{r}, t) = \sum_\mu \tilde{\phi}_\mu(\mathbf{r}) C_{\mu n}(t), \quad (\text{S2})$$

with expansion coefficients $C_{\mu n}(t)$. The corresponding all-electron wave function is given by [compare to Eq. (7)]

$$\psi_n(\mathbf{r}, t) = \mathcal{T} \tilde{\psi}_n(\mathbf{r}, t) = \sum_\mu \phi_\mu(\mathbf{r}) C_{\mu n}(t), \quad (\text{S3})$$

where the all-electron basis functions have been defined as $\phi_\mu = \mathcal{T} \tilde{\phi}_\mu$.

The time-dependent all-electron real-space KS density matrix can be obtained as

$$\rho(\mathbf{r}, \mathbf{r}', t) = \sum_{\mu\nu} \phi_\mu(\mathbf{r}) \rho_{\mu\nu}(t) \phi_\nu^*(\mathbf{r}'), \quad (\text{S4})$$

where the KS density matrix in the LCAO basis $\rho_{\mu\nu}(t)$ is given by Eq. (8).

The transformation of the real-space KS density matrix to the basis defined by the ground-state KS orbitals $\psi_n^{(0)}(\mathbf{r})$, see Eq. (3), is given by

$$\rho_{nn'}(t) = \int d\mathbf{r} \int d\mathbf{r}' \psi_n^{(0)*}(\mathbf{r}) \rho(\mathbf{r}, \mathbf{r}', t) \psi_{n'}^{(0)}(\mathbf{r}'). \quad (\text{S5})$$

By expanding $\psi_n^{(0)}(\mathbf{r})$ in the LCAO basis as in Eq. (S3) and inserting Eq. (S4) into Eq. (S5), we obtain after re-ordering the integrals

$$\begin{aligned} \rho_{nn'}(t) = & \sum_\mu C_{\mu n}^{(0)*} \underbrace{\sum_{\mu'} \int d\mathbf{r} \phi_\mu^*(\mathbf{r}) \phi_{\mu'}(\mathbf{r})}_{S_{\mu\mu'}} \\ & \cdot \sum_{\nu'} \rho_{\mu'\nu'}(t) \underbrace{\sum_\nu \int d\mathbf{r}' \phi_{\nu'}^*(\mathbf{r}') \phi_\nu(\mathbf{r}')}_{S_{\nu\nu'}^*} C_{\nu n'}^{(0)}. \end{aligned} \quad (\text{S6})$$

Here, we have isolated the overlap integrals $S_{\mu\mu'}$ used regularly in LCAO calculations, *i.e.*,

$$S_{\mu\mu'} = \int d\mathbf{r} \phi_\mu^*(\mathbf{r}) \phi_{\mu'}(\mathbf{r}) = \int d\mathbf{r} \tilde{\phi}_\mu^*(\mathbf{r}) \mathcal{T}^\dagger \mathcal{T} \tilde{\phi}_{\mu'}(\mathbf{r}). \quad (\text{S7})$$

After simplifying the overlap integrals in Eq. (S6), we obtain Eq. (9). We note that the PAW transformation affects only the evaluation of the overlap integrals $S_{\mu\mu'}$, see Eq. (S7).

S2 Additional transition contribution maps for Ag₅₅

On the following pages, we present the transition contribution maps (TCMs) for the Ag₅₅ nanoparticle from $\omega = 3.66$ eV to $\omega = 4.25$ eV in steps of 0.01 eV. All TCMs use the same color scale (see the main text for definitions).









