

Supporting Information for: Molecular Vibration Induced Plasmon Decay

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Analysis of the transitions in the LR-TDDFT calculations for Ag₄ with 5 quanta of vibrational energy

Time step: 100fs

The transverse mode is composed of two nearly degenerate peaks (the x- and y-components of the Π_u transition of perfectly straight nanowires). At 100 fs, the nanowire has not significantly distorted from its initial linear configuration. The two peaks comprising the transverse mode appear around 5.01 eV with an oscillator strength of $f \sim 0.7$. These excited states are:

Excited State 32: 5.0099 eV 247.48 nm $f=0.7399$ $\langle S^2 \rangle=0.000$

19 -> 39 0.11266

34 -> 40 -0.10597

37 -> 41 0.47821 ($\Sigma_1 \rightarrow \Pi_1$ (y component))

38 -> 44 0.43914 ($\Sigma_2 \rightarrow \Pi_2$ (y component))

Excited State 33: 5.0183 eV 247.06 nm $f=0.6667$ $\langle S^2 \rangle=0.000$

35 -> 40 -0.10308

37 -> 42 0.45719 ($\Sigma_1 \rightarrow \Pi_1$ (x component))

38 -> 43 0.33699 ($\Sigma_2 \rightarrow \Pi_2$ (x component))

38 -> 45 -0.31289

The 37->41 and 37->42 transitions are $\Sigma_1 \rightarrow \Pi_1$ transitions. The 38->43 and 38->44 transitions are $\Sigma_2 \rightarrow \Pi_2$ transitions. Both of these transitions between delocalized orbitals are known to result in large transition dipole moments.¹ These transitions are shown in blue throughout the SI.

Questions we will answer: The two primary transitions are mixed with other excitations/deexcitations. How might these contributions change over time? How does the splitting between the x and y components of the transverse mode change over time?

Excited state 29 also has some oscillator strength (most likely because of its 37->42 and 38->43 contributions). However, it is currently much weaker than states 32 and 33.

Excited State 29: 4.9047 eV 252.79 nm $f=0.0806$ $\langle S^2 \rangle=0.000$

18 -> 39	-0.10159
35 -> 40	-0.16619
37 -> 40	-0.15889
37 -> 42	0.12613
38 -> 43	0.29070
38 -> 45	0.53853

It is important to note that the peaks with the largest weights are *not necessarily* the ones that contribute the most to the oscillator strength. They often do for ‘normal’ molecules, but we see exceptions to this commonly in the case of systems with delocalized orbitals exhibiting collective excitation peaks. (The 38->45 transition becomes important to the analysis at later times.)

States 32 and 33 arise because of constructive interference between the two $\Sigma_n \rightarrow \Pi_n$ transitions. Due to configuration interaction, two (or more, if additional states are mixed) states arising from destructive interference of these transitions must be present.¹ The destructive interference states that are related to the strong transverse modes are 24, 25 and 26:

Excited State 24: 4.3490 eV 285.09 nm $f=0.0057$ $\langle S^2 \rangle=0.000$

19 -> 39	0.61591
37 -> 41	-0.30076
38 -> 44	0.17227

Excited State 25: 4.4416 eV 279.14 nm $f=0.0233$ $\langle S^2 \rangle=0.000$

18 -> 39	0.12320
20 -> 39	-0.24790
37 -> 42	-0.41396
38 -> 43	0.48491
38 -> 45	-0.11114

Excited State 26: 4.4550 eV 278.31 nm $f=0.0084$ $\langle S^2 \rangle=0.000$

19 -> 39	-0.32377
37 -> 41	-0.37927
38 -> 44	0.49747

These peaks are in reasonable agreement with the energy of the “new peak” around 4.3 eV that is forming at later time steps. We get two peaks that involve 37->41 and 38->44 because of the +/- combinations with the 19->39 transition. (The 19->39 transition is probably not very important since the splitting and oscillator strengths of these peaks are not very large regardless of whether we consider the + or – combination.)

Time step: 200fs

At 200fs, the nanowire is more distorted from linear compared to the system at 100 fs. The two excited states responsible for the transverse plasmon still occur around similar energy (5.01-5.05 eV), although the x component (corresponding to the direction in which activation of

the vibrational mode has offset the geometry) of the peak has changed slightly in energy and its oscillator strength is somewhat reduced:

Excited State 32: 5.0089 eV 247.53 nm $f=0.7351$ $\langle S^{**2} \rangle=0.000$

19 -> 39 0.11172

37 -> 41 0.48076

38 -> 44 0.43650

Excited State 34: 5.0474 eV 245.64 nm $f=0.4920$ $\langle S^{**2} \rangle=0.000$

18 -> 39 -0.12490

37 -> 42 -0.40892

38 -> 43 -0.17733

38 -> 45 0.46680

These are still the two modes in the correct range with the strongest oscillator strength. However, the 38->45 mode (shown in green) mixes into state 34 rather strongly, which seems to cause more splitting of the two peak energies, an oscillator strength for 34 that differs significantly from $f \sim 0.7$, and unequal weights of the 37->42 and 38->43 transitions.

States 27, 29, and 30 (energy range: 4.65-4.90 eV) have moderately strong oscillator strengths ($f \sim 0.12$ -0.14). This is likely because of transitions out of orbitals 37 and 38 into orbitals 42-45.

Excited State 27: 4.6465 eV 266.83 nm $f=0.1429$ $\langle S^{**2} \rangle=0.000$

18 -> 39 0.67394

38 -> 43 -0.10175

38 -> 45 0.15865

Excited State 29: 4.8732 eV 254.42 nm $f=0.1252$ $\langle S^{**2} \rangle=0.000$

35 -> 40 0.44781

37 -> 40 0.12331

37 -> 42 -0.17974

38 -> 43 -0.31171

38 -> 45 -0.34831

Excited State 30: 4.9035 eV 252.85 nm $f=0.1270$ $\langle S^{**2} \rangle=0.000$

35 -> 40 0.53885

37 -> 40 -0.11480

37 -> 42 0.19641

38 -> 43 0.29834

38 -> 45 0.20813

Furthermore, states around 4.35-4.45 eV are also present and arise from destructive interference of these transitions:

Excited State 24: 4.3538 eV 284.77 nm $f=0.0050$ $\langle S^{**2} \rangle=0.000$

19 -> 39 0.60967

37 -> 41 -0.30630

38 -> 44 0.18429

Excited State 25: 4.4369 eV 279.44 nm $f=0.0140$ $\langle S^{**2} \rangle=0.000$

18 -> 39	0.12171
37 -> 42	-0.46297
38 -> 43	0.46533
38 -> 45	-0.20573

Excited State 26: 4.4549 eV 278.31 nm f=0.0094 <S**2>=0.000

19 -> 39	-0.33580
37 -> 41	-0.37117
38 -> 44	0.49558

Note: Again, 24 and 26 have very similar contributions, but mix in with the 19->39 transition in +/- ways. The oscillator strength contribution from the 19->39 transition is probably relatively small since both of these states have small, relatively similar oscillator strengths.

Time step: 300fs

At 300fs, the nanowire is more distorted from linear compared to the system at 100 or 200 fs. The two excited states responsible for the transverse plasmon still occur around similar energy (5.01-5.10 eV), although the x component of the peak has increased further in energy and its oscillator strength is further reduced:

Excited State 33: 5.0083 eV 247.56 nm f=0.7034 <S**2>=0.000

19 -> 39	0.11153
30 -> 40	0.12186
37 -> 41	0.48511
38 -> 44	0.42160

Excited State 36: 5.1023 eV 243.00 nm f=0.2650 <S**2>=0.000

18 -> 39	-0.13285
31 -> 40	-0.23456
37 -> 42	-0.34087
38 -> 45	0.50390

The oscillator strength of the other strong mode in this region has decreased (now only 0.2650). The energy of this part of the transverse mode has increased. This correlates with a slight increase in the energy of the transverse mode that is observed in the Ehrenfest dynamics with real-time TDDFT. In addition, the transitions responsible for this mode have changed. Orbital 43 is no longer significantly involved, and 45 has a strong role. The 38->43 transition now appears to be involved in lower energy excited states, which makes sense because the $\Sigma_2 \rightarrow \Pi_2(x)$ (38->43) orbital energy difference is decreasing (Figure 7).

Lower energy excited states 28 and 31 include contributions from 38->43 (note: these still have a constructive interaction between 37->42 and 38->43). These states are:

Excited State 28: 4.7805 eV 259.35 nm f=0.1769 <S**2>=0.000

18 -> 39	0.50707
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35 -> 40	0.32929
37 -> 42	-0.21830
38 -> 43	-0.24288

Excited State 31: 4.8825 eV 253.94 nm f=0.2891 <S**2>=0.000

18 -> 39	0.26620
35 -> 40	0.19893
37 -> 40	-0.16402
37 -> 42	0.29462
38 -> 43	0.34700
38 -> 45	0.32792

In addition, the “destructive interference” excited states are still present and include 38->43:

Excited State 26: 4.4621 eV 277.86 nm f=0.0369 <S**2>=0.000

37 -> 40	-0.11949
37 -> 42	-0.44384
38 -> 43	0.46807
38 -> 45	-0.22729

There are also the two destructive ones that involve 37->41 and 38->44; these are related to the y transverse mode that has not changed in energy/oscillator strength:

Excited State 24: 4.3771 eV 283.26 nm f=0.0022 <S**2>=0.000

19 -> 39	0.55821
37 -> 41	-0.34866
38 -> 44	0.25743

Excited State 25: 4.4613 eV 277.91 nm f=0.0149 <S**2>=0.000

19 -> 39	-0.41577
37 -> 41	-0.31671
38 -> 44	0.47209

Time step: 400fs

At 400fs, the nanowire is significantly more distorted from linear compared to the system at earlier times. The y component of the transverse plasmon still arises around 4.99 eV:

Excited State 34: 4.9947 eV 248.23 nm f=0.4790 <S**2>=0.000

20 -> 39	0.10593
26 -> 40	0.11258
30 -> 40	0.35455
37 -> 41	0.44174
38 -> 44	0.34371

At 400 fs the oscillator strength of this peak is actually starting to decrease. It is also interesting to note that more transitions now have strong weights. These observations are related, since the

interaction with the 26->40 and 30->40 transitions splits this peak into two peaks, so the oscillator strength is divided. The other related peak is:

Excited State 37: 5.0937 eV 243.41 nm $f=0.1691$ $\langle S^{**2} \rangle=0.000$

26 -> 40	-0.22892
30 -> 40	0.59297
37 -> 41	-0.22451
38 -> 44	-0.17289

(Again, is important to note that the peaks with the largest weights are *not necessarily* the ones that contribute the most to the oscillator strength.)

The other transverse peak in this energy range is:

Excited State 36: 5.0434 eV 245.84 nm $f=0.1076$ $\langle S^{**2} \rangle=0.000$

18 -> 39	-0.36067
31 -> 40	0.48547
38 -> 45	-0.31693

Again, the 38->45 transition has become important. Moreover, the 38->43 transition does not exceed Gaussian's default cutoff to show the weight of this transition. Instead, the 38->43 transition contributes primarily to a variety of excited states at lower energy (which again makes sense since the $\Sigma_2 \rightarrow \Pi_2(x)$ energy difference is lowered in energy):

Excited State 30: 4.8289 eV 256.75 nm $f=0.3330$ $\langle S^{**2} \rangle=0.000$

18 -> 39	0.25769
37 -> 40	0.12332
37 -> 42	0.47757
38 -> 43	-0.29703
38 -> 45	-0.24747

Excited State 27: 4.6593 eV 266.10 nm $f=0.0178$ $\langle S^{**2} \rangle=0.000$

36 -> 40	0.66165
37 -> 42	-0.14612
38 -> 43	-0.11056
38 -> 45	0.13281

Excited State 24: 4.4781 eV 276.87 nm $f=0.1774$ $\langle S^{**2} \rangle=0.000$

22 -> 39	-0.13008
36 -> 40	0.21873
37 -> 40	-0.17953
37 -> 42	0.30654
38 -> 43	0.51282
38 -> 45	-0.16173

Excited State 21: 4.2695 eV 290.39 nm $f=0.0130$ $\langle S^{**2} \rangle=0.000$

22 -> 39	0.68350
38 -> 43	0.17025

We also see the destructive peak that correlates with the ‘normal’ transverse peak that is essentially unaffected by the vibration:

Excited State 23: 4.4097 eV 281.16 nm $f=0.0027$ $\langle S^{**2} \rangle=0.000$
 20 -> 39 0.30381
 37 -> 41 -0.43419
 38 -> 44 0.46659

Time step: 500fs

At 500fs, the nanowire is very distorted from linear. The y component of the transverse plasmon still arises around 4.98-5.13 eV. It again arises from excitations involving 37->41 and 38->44, which are again slightly split due to interactions with other transitions:

Excited State 38: 4.9815 eV 248.89 nm $f=0.1654$ $\langle S^{**2} \rangle=0.000$
 20 -> 39 0.11391
 26 -> 40 0.57706
 37 -> 41 0.31641
 38 -> 44 0.19359
 Excited State 41: 5.1346 eV 241.47 nm $f=0.5304$ $\langle S^{**2} \rangle=0.000$
 26 -> 40 -0.40067
 30 -> 40 -0.14831
 37 -> 41 0.42285
 38 -> 44 0.29172

Very few other strong peaks can be observed in the ~ 5 eV region. Because the 37->41 and 38->44 transitions are related to the y component of the excitation that is orthogonal to the vibration in the x direction (and, more importantly in this case, to the applied electric field in the x direction), we should not observe many peaks in this region of the Ehrenfest spectrum, which agrees with the observations. There is a small 38->45 contribution to several excited states as follows, but the oscillator strengths are low:

Excited State 37: 4.9500 eV 250.47 nm $f=0.0315$ $\langle S^{**2} \rangle=0.000$
 27 -> 40 0.64788
 37 -> 42 0.18633
 38 -> 45 -0.17542
 Excited State 43: 5.2312 eV 237.01 nm $f=0.0997$ $\langle S^{**2} \rangle=0.000$
 17 -> 39 0.59485
 38 -> 45 0.32645
 Excited State 48: 5.3769 eV 230.59 nm $f=0.0510$ $\langle S^{**2} \rangle=0.000$
 17 -> 39 -0.20541
 24 -> 40 -0.18393
 34 -> 41 0.11639
 37 -> 42 0.24844
 38 -> 45 0.37221
 38 -> 47 0.41836

We note the destructive peak related to 37->41 and 38->44 (again, still essentially unchanged):

Excited State 22: 4.4189 eV 280.58 nm f=0.0259 <S**2>=0.000

37 -> 41 -0.40579

38 -> 44 0.56923

So, a question arises: where has all the oscillator strength gone?

This can be answered by considering the energy region from 4.40-4.93 eV, in which a number of new states with reasonable oscillator strength appear:

Excited State 21: 4.4011 eV 281.71 nm f=0.2080 <S**2>=0.000

36 -> 40 0.47056

37 -> 40 0.16786

37 -> 42 -0.12132

38 -> 43 -0.45153

Excited State 24: 4.5213 eV 274.22 nm f=0.1568 <S**2>=0.000

22 -> 39 0.50751

36 -> 40 -0.36036

37 -> 40 0.12598

38 -> 43 -0.27318

Excited State 26: 4.6089 eV 269.01 nm f=0.1468 <S**2>=0.000

22 -> 39 0.48655

36 -> 40 0.37184

37 -> 40 -0.15293

38 -> 43 0.25205

Excited State 34: 4.9273 eV 251.63 nm f=0.1176 <S**2>=0.000

17 -> 39 0.13742

27 -> 40 -0.27336

37 -> 42 0.51730

38 -> 43 -0.12842

38 -> 45 -0.32017

These states have contributions from 38->43 and occasionally 37->42, which is likely why they have relatively large oscillator strengths. However, there is significant mixing from other transitions (notably 36->40 and 37->40, as well as excitations from occupied orbitals much lower in energy), so the picture is not extremely clean.

Conclusions (LR-TDDFT calculations, 100-500 fs)

Overall, the change in the geometry as time progresses during the simulation is responsible for decreasing the energy of the 38->43 transition. In consequence, how this transition mixes with other transitions is affected. The oscillator strengths of the excited states change, and a significant portion of the oscillator strength is shifted into lower energy excited states.

Orbitals involved in the transverse plasmon excitation from the Ehrenfest dynamics

The transverse plasmon excitation is characterized by two electronic transitions: $\Sigma_1 \rightarrow \Pi_1$ and $\Sigma_2 \rightarrow \Pi_2$. The orbitals involved in both transitions are shown for some representative snapshots extracted from the Ehrenfest dynamics having as initial molecular structure and momenta the ones generated by imparting zero-point vibrational energy into the out-of-chain bending mode. In particular they are shown for the configurations extracted at 100, 200, 300, 400 and 500 fs in Fig. S1, Fig. S2, Fig.S3, Fig.S4 and Fig.S5, respectively.

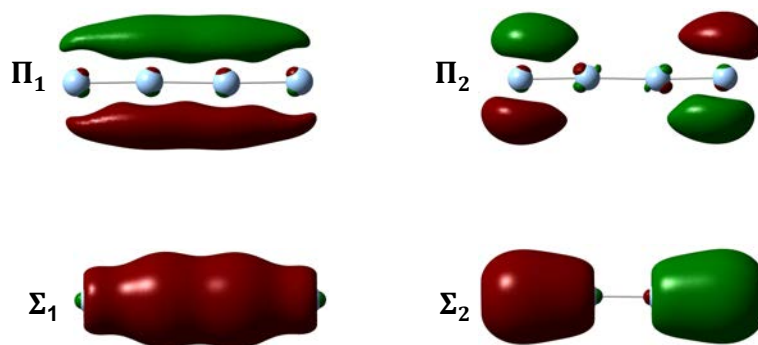


Figure S1. Orbitals involved in the transverse plasmon excitation for the configuration extracted from the Ehrenfest dynamics at 100 fs.

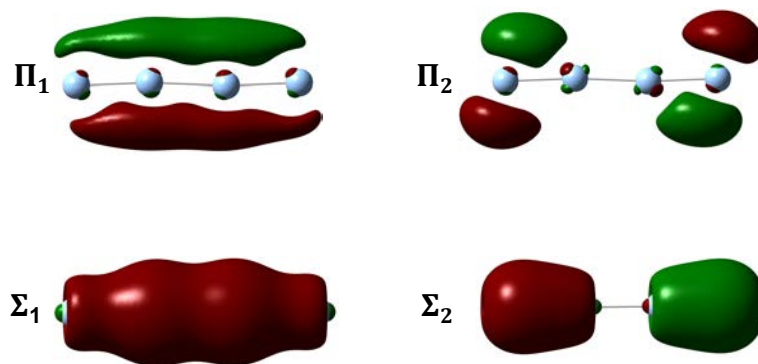


Figure S2. Orbitals involved in the transverse plasmon excitation for the configuration extracted from the Ehrenfest dynamics at 200 fs.

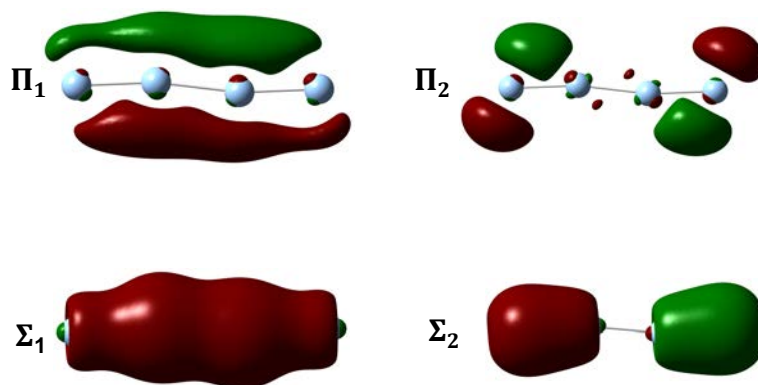


Figure S3. Orbitals involved in the transverse plasmon excitation for the configuration extracted from the Ehrenfest dynamics at 300 fs.

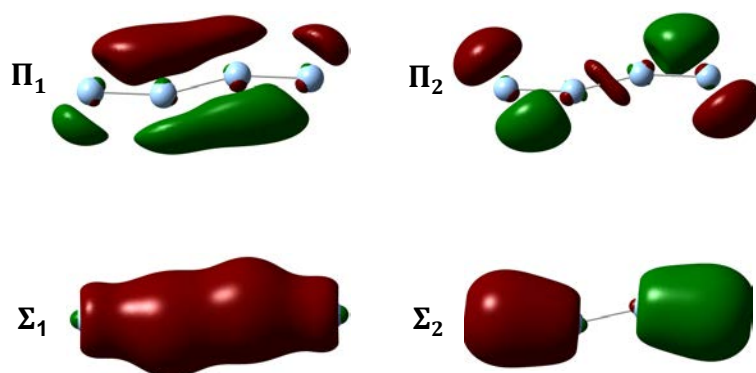


Figure S4. Orbitals involved in the transverse plasmon excitation for the configuration extracted from the Ehrenfest dynamics at 400 fs.

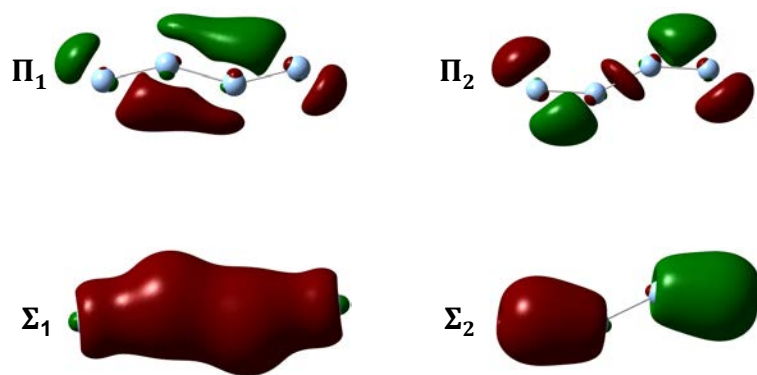


Figure S5. Orbitals involved in the transverse plasmon excitation for the configuration extracted from the Ehrenfest dynamics at 500 fs.

References

1. Guidez, E. B.; Aikens, C. M., Theoretical Analysis of the Optical Excitation Spectra of Silver and Gold Nanowires. *Nanoscale* **2012**, *4*, 4190-4198.