

# Spontaneous Uphill Movement and Self-Removal of Condensates on Hierarchical Tower-Like Arrays

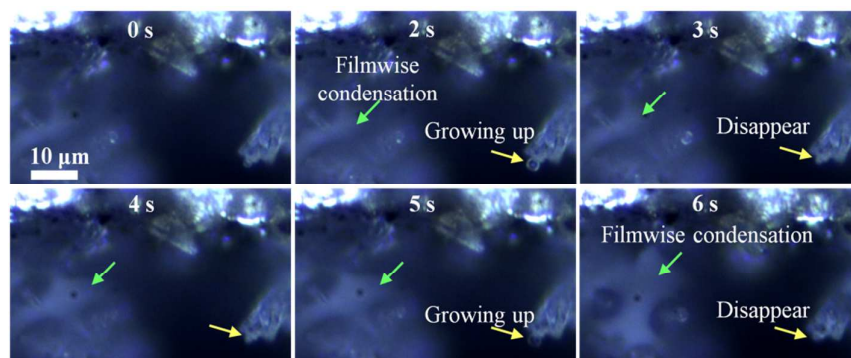
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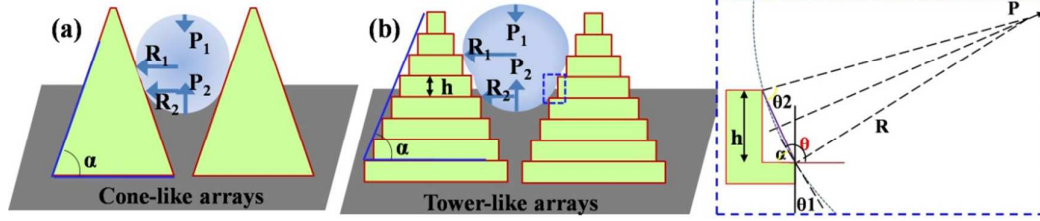


**Figure S1.** Condensation on bare arrays without being modified by PDMS, green arrows show the filmwise condensation, yellow arrows show the formation, growth and disappearance of tiny condensed droplets.



**Figure S2.** A water drop on the bare tower-like structure, the average contact angle of droplets on such surface is 11.3°.

**The Laplace pressure driving condensed droplets moving uphill to the top of tower-like structures**



**Figure S3.** Scheme of a micro-droplet suspending on neighboring cone-like structures (a) and hierarchical towers (b).

The Laplace pressure ( $P$ ) from the ratchet-like sidewall is

$$P = \frac{-2\gamma}{R} \quad (1)$$

$R$  is the radius of curvature,  $\gamma$  is the water surface tension.

In Figure S1 (a), the total Laplace pressure for a droplet entrapped in two neighboring cone-like structures is

$$P^{\text{total}} = -2\gamma \left( \frac{1}{R_2} - \frac{1}{R_1} \right) \cos \theta \quad (2)$$

In Figure S1 (b), The upward component of  $P$  for a droplet entrapped in two neighboring hierarchical towers is

$$P_{\text{up}} = P \cos \alpha \quad (3)$$

$\alpha$  is the tilted angle of a tower-like structure,

$$L = 2R \cos \theta_2 \quad (4)$$

and

$$L = \frac{h}{2 \sin \alpha} \quad (5)$$

$L$  is the length of the line connecting neighboring nano-plates' edge,  $h$  is the thickness of a single plate.

With simple geometric operation, we obtain:

$$\theta_2 = \pi - \alpha - \theta_1 \quad (6)$$

$$\theta_1 = \theta - \frac{\pi}{2} \quad (7)$$

$\theta_1$  is the angle between the tangent of curved water surface and the perpendicular direction,  $\theta_2$  is the angle between the normal line and the line connecting neighboring nano-plates' edge,  $\theta$  is the apparent CA of water droplets on the surface.

Combining equations (1)-(5) yields an expression that relates the Laplace pressure and the surface properties

$$P_{up} = \frac{-4\gamma \sin(2\alpha) \cos(\alpha + \theta_1)}{h} \quad (8)$$

Thus, the total Laplace pressure driving condensates moving uphill to the top of microtowers is

$$P_{up}^{total} = 2\gamma \left( \frac{1}{R_1} - \frac{1}{R_2} \right) \cos \theta + n \sum_{i=1}^j P_{up} \quad (9)$$

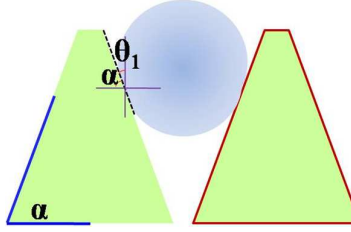
with

$$j = \frac{(R_1 - R_2) \tan \alpha}{h} \quad (10)$$

here  $0 \leq \theta \leq \alpha$

$R_1$  is the half pitch of nano-plates on neighboring tower-like structures contacting with the up curved water surface,  $R_2$  is the half pitch of nano-plates on neighboring tower-like structures contacting with the bottom curved water surface,  $n$  is the number

of tower-like structures contacting with a microdroplet,  $j$  is the number of nano-plates on a single tower-like structure contacting with a microdroplet.



**Figure S4.** Condensed droplet suspends between two cone-like structures.

On the cone-like structures, the  $(\alpha + \theta_1)$  is equal to  $\pi/2$  and consequently the cosine of  $(\alpha + \theta_1)$  is equal to 0 (Figure S4). As a result, the equation (9) can be reduced to the equation (2) which is applied to cone-like structures as mentioned above.

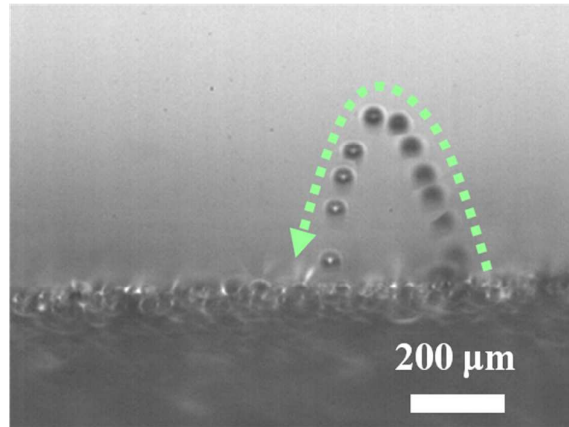
By comparison the Laplace pressure on the two surfaces, we obtain

$$\frac{(\text{tower})P_{\text{up}}^{\text{total}}}{(\text{cone})P_{\text{up}}^{\text{total}}} = 1 + \frac{n \sum_{i=1}^j P_{\text{up}}}{2\gamma \left( \frac{1}{R_1} - \frac{1}{R_2} \right) \cos\theta} \quad (11)$$

Combining equation (6), (7) and (10), the following equation is obtained

$$\frac{(\text{tower})P_{\text{up}}^{\text{total}}}{(\text{cone})P_{\text{up}}^{\text{total}}} = 1 + \frac{n \sum_{i=1}^j R_1 R_2 \sin(2\alpha) \cos(\alpha + \theta)}{(R_1 - R_2) h \cos\theta} \quad (12)$$

The ratio of Laplace pressure on tower-like structures and cone-like structures suggests the advantage of tower-like structures in uphill movement of condensates. Thus, maximum values of the  $\theta_i$  can be determined by the first and second derivative of eq. 11 with respect to  $R_2$ . The result show when  $R_1$  is 2 times of  $R_2$  in the case of  $h$  is  $\sim 100$  nm, the maximum  $R_1$  is  $\sim 2.4$   $\mu\text{m}$  (the pitch of two neighboring microtowers),  $\theta$  is  $108^\circ$ .  $\alpha$  is calculated to be  $\sim 72.6^\circ$ , the Laplace pressure on the tower-like structures is  $\sim 30$  times larger than that on the cone-like structures.



**Figure S5.** An overlapped optical side-view showing the trajectory of a merged micro-droplet jumping from the horizontal  $S_t$ . Dashed green arrow indicates the jump and fall of condensed droplets.