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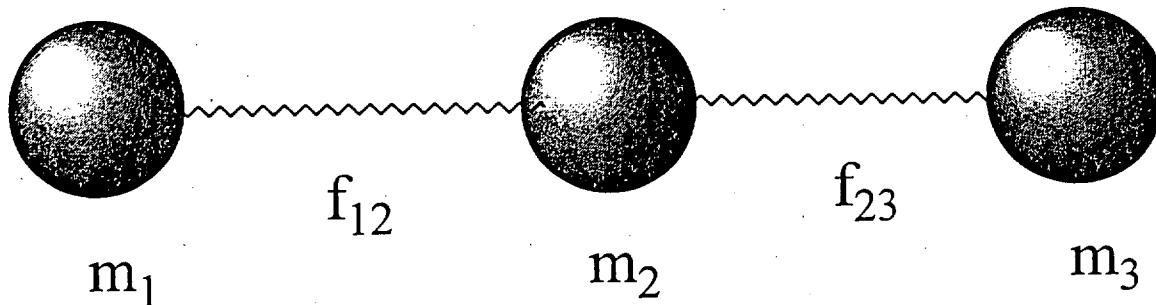
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## Supplementary Material

## Three-Body Model

The three mass, spring model is:



The potential energy of the system is given by:

$$2V = f_{12}(\Delta r_{12})^2 + f_{23}(\Delta r_{23})^2$$

where  $\Delta r_{12} = \Delta x_1 - \Delta x_2$  and  $\Delta r_{23} = \Delta x_2 - \Delta x_3$  so that

$$2V = f_{12}(\Delta x_1 - \Delta x_2)^2 + f_{23}(\Delta x_2 - \Delta x_3)^2$$

$$\text{or } 2V = f_{12}(\Delta x_1)^2 + (f_{12} + f_{23})(\Delta x_2)^2 + f_{23}(\Delta x_3)^2 - 2f_{23}(\Delta x_2 \Delta x_3) - 2f_{12}(\Delta x_1 \Delta x_2)$$

The kinetic energy is

$$2T = m_1(\Delta x_1)^2 + m_2(\Delta x_2)^2 + m_3(\Delta x_3)^2$$

The equations of motion are then:

$$m_1(\Delta x_1) + f_{12}(\Delta x_1) - f_{12}(\Delta x_2) = 0$$

$$m_2(\Delta x_2) + (f_{12} + f_{23})(\Delta x_2) - f_{23}(\Delta x_3) - f_{12}(\Delta x_1) = 0$$

$$m_3(\Delta x_3) + f_{23}(\Delta x_3) - f_{23}(\Delta x_2) = 0$$

The substitution

$$\Delta x_j = A_j \cos(\sqrt{\lambda} t + \epsilon)$$

leads to the equations

$$(f_{12} - m_1\lambda)A_1 - f_{12}A_2 = 0$$

$$-f_{12}A_1 + (f_{12} + f_{23} - m_2\lambda)A_2 - f_{23}A_3 = 0$$

$$(f_{23} - m_3\lambda)A_3 - f_{23}A_2 = 0$$

The equations have non-zero solutions only if the determinant of the coefficients vanishes:

$$\begin{array}{ccc} (f_{12} - m_1\lambda) & -f_{12} & 0 \\ -f_{12} & (f_{12} + f_{23} - m_2\lambda) & -f_{23} \\ 0 & -f_{23} & (f_{23} - m_3\lambda) \end{array}$$

$\lambda = 0$  is one eigenvalue because there are only two force constants.

$$-m_1m_3m_2\lambda^2 + ((m_1 + m_2)m_3f_{12} + (m_2 + m_3)m_1f_{23})\lambda - (m_1 + m_2 + m_3)f_{12}f_{23} = 0$$

which has the solution

$$\lambda = \frac{\left((m_1 + m_2)m_3f_{12} + (m_2 + m_3)m_1f_{23}\right) \pm \sqrt{\left((m_1 + m_2)m_3f_{12} + (m_2 + m_3)m_1f_{23}\right)^2 - 4m_1m_2m_3(m_1 + m_2 + m_3)}}{2m_1m_2m_3}$$

In the special case where the force constants are equal:

$$-m_1m_3m_2\lambda^2 + f(m_2m_3 + m_1m_2 + 2m_1m_3)\lambda - f^2(m_1 + m_2 + m_3) = 0$$

the result is

$$\lambda = f \frac{\left(m_2m_3 + m_1m_2 + 2m_1m_3\right) \pm \sqrt{m_2^2m_3^2 + m_1^2m_2^2 + 4m_1^2m_3^2 - 2m_1m_2m_3}}{2m_1m_2m_3}$$

Our three-body masses are:

$$m_1 = 488 \text{ a.m.u}$$

$$m_2 = 56 \text{ a.m.u}$$

$$m_3 = 68 \text{ a.m.u}$$

It would be possible for the observed vibrational mode to be either (+) predominantly 12 or (-) predominantly 23 combination. For example, assuming  $f = 3.36 \text{ mdyne}/\text{\AA}$  the two normal modes in this analysis have frequencies of 507.3 (+) and 203.7 (-).

A table of isotope effects for three body model with identical force constants of 3.36 mdyne/\AA are given below (Table A1). This result can be compared to the two body model (again with  $f = 3.36 \text{ mdyne}/\text{\AA}$ ). For a two-body model  $1/\mu = (m_1 + m_2)/m_1m_2 = 0.0325$  and the frequency is  $431.0 \text{ cm}^{-1}$  for this force constant. The three-body model presented with 3.36 mdyne/\AA for the axial ligand mode is acceptable

provided that the observed mode is the (-) combination. Notice that the pattern of isotope shifts is that shifts on the iron are smaller and imidazole are larger for the (-) combination. Neither this nor the fact that the (+) combination should also be observed are consistent with the data on the iron-axial-ligand stretching mode in heme proteins.

Table A1. Isotope effects calculated for a three-body model with two identical force constants of 3.36 mdyne/ Å

	$1/\mu$	$1/\mu$	$\text{sq}(\mu'/\mu)-1$	$\Delta\nu$	$\text{sq}(\mu'/\mu)-1$	$\Delta\nu$
Mode	(+)	(-)	(+)	507 cm <sup>-1</sup>	(-)	203 cm <sup>-1</sup>
Na	0.045180	0.0072893	0.0000	-----	0.000	-----
4N <sub>p</sub> on P	0.045174	0.0072782	-6.64 E -5	-0.033	-7.61 E -4	-0.15
d <sub>3</sub> Im	0.044767	0.0070801	-0.0040	-2.02	-0.0144	-2.92
di- <sup>15</sup> N Im	0.044901	0.0071483	-0.0031	-1.57	-0.00971	-1.97
<sup>54</sup> Fe	0.046466	0.0073260	0.0141	-7.16	0.00251	-0.509
<sup>57</sup> Fe	0.044572	0.0072710	-0.0067	3.39	-0.0012	0.24

Since this result is not acceptable the relative frequencies of the two modes for different values of the force constant f are compared in Table A2.

Table A2. Isotope effects calculated for a three-body model with two identical force constant. The calculation is shown as a function of the force constant given in the first column.

f (mdyne/Å)	(+) (cm <sup>-1</sup> )	(-) (cm <sup>-1</sup> )
3.36	507.3	203.7
1.67	358.0	143.8
0.7	231.7	93.0
0.66	224.8	88.6

The calculation with force constants of 0.66 mdyne/Å has the closest agreement with experiment in terms of the frequency of the (+) combination mode. The isotope effect can be readily calculated for the most acceptable values as shown in Table A3.

Table A3. Isotope effects for the three-body model calculation with two identical force constant of 0.66 mdyne/ Å.

Mode	1/μ (+)	1/μ (-)	sq(μ'/μ)-1 (+)	Δν 225 cm⁻¹	sq(μ'/μ)-1 (-)	Δν 88.6 cm⁻¹
Na	0.045180	0.0072893	0.0000	-----	0.000	-----
4N <sub>p</sub> on P	0.045174	0.0072782	-6.64 E -5	-0.022	-7.61 E -4	-0.06
d <sub>3</sub> Im	0.044767	0.0070801	-0.0040	-0.9	-0.0144	-1.3
di- <sup>15</sup> N Im	0.044901	0.0071483	-0.0031	-0.7	-0.00971	-0.9
<sup>54</sup> Fe	0.046466	0.0073260	0.0141	-3.2	0.00251	-0.22
<sup>57</sup> Fe	0.044572	0.0072710	-0.0067	1.5	-0.0012	0.1

Using the approximation of identical force constants for the two internal stretching coordinates of the three-body model shown in Figure 1A, we can calculate the percent potential energy distribution (%PED) as follows.

Beginning with the definition of equations for the linear coefficient  $A_i$  and equating  $f_{12} = f_{23} = f$ , we have

$$\begin{aligned}(f - m_1\lambda)A_1 - fA_2 &= 0 \\ -fA_1 + (2f - m_2\lambda)A_2 - fA_3 &= 0 \\ (f - m_3\lambda)A_3 - fA_2 &= 0\end{aligned}$$

We found eigenvalues of  $\lambda = 0.045$  and  $0.007$  for the (+) and (-) solutions, respectively.

Thus, for a given force constant and  $\lambda$  we can solve for the ratio of  $A_1$  and  $A_3$  to  $A_2$ .

for  $\lambda = 0.007$   
we have for the (-) combination

$$\begin{aligned}A_1 &= [f/(f - m_1\lambda_-)]A_2 = f/(f - 0.007 * 488) A_2 = K_- A_2 \\ A_3 &= [f/(f - m_3\lambda_-)]A_2 = f/(f - 0.007 * 68) A_2 = J_- A_2\end{aligned}$$

for  $\lambda = 0.045$   
we have for the (+) combination

$$\begin{aligned}A_1 &= [f/(f - m_1\lambda_+)]A_2 = f/(f - 0.045 * 488) A_2 = K_+ A_2 \\ A_3 &= [f/(f - m_3\lambda_+)]A_2 = f/(f - 0.045 * 68) A_2 = J_+ A_2\end{aligned}$$

The potential energy is given by  $V = \frac{1}{2}[f_{12}(\Delta x_1 - \Delta x_2)^2 + f_{23}(\Delta x_2 - \Delta x_3)^2]$  so that we can define the total potential energy as  $V = \frac{1}{2}[f(\Delta x_1 - \Delta x_2)^2 + f(\Delta x_2 - \Delta x_3)^2]$  for the equal force constant model discussed here. Substituting in the above information in terms of unnormalized displacements we have:

$$V = \frac{1}{2}[f([K - 1]A_2)^2 + f([J - 1]A_2)^2] \text{ where } A_1 = KA_2 \text{ and } A_3 = JA_2.$$

The %PED for either of the coordinates is given by:

$$\frac{100*(K - 1)^2}{[(K - 1)^2 + (J - 1)^2]}$$

So that we can calculate for the models above that the %PEDs are given in table A4.

Table A4. Potential energy distributions for three-body model calculations.

f	(+)	(+)	(-)	(-)
(mdyne/Å)	Fe-L	Fe-P	Fe-L	Fe-P
3.36	100.0	0.0	1.3	98.7
1.67	80.6	19.4	4.0	96.0
0.7	61.6	38.5	24.7	75.3
0.66	60.5	39.5	18.5	81.5

Several calculations were considered with variable force constants for the purpose of exploring the parameter space.

#### Case 1

$$f_{12} = 3.36$$

$$f_{23} = 2.0$$

Calculate numerically.

$\lambda = 0.01970$  and  $0.11228$  which corresponds to frequencies of  $186.7$  and  $445.74 \text{ cm}^{-1}$

	$\lambda$ (+)	$\lambda$ (-)	$ (\lambda'/\lambda)-1 $ (+)	v $\Delta v$	$(\lambda'/\lambda)-1$ (-)	v $\Delta v$
Na	0.11228	0.019700	-----	$445.7 \text{ cm}^{-1}$	-----	$99.9 \text{ cm}^{-1}$
$^{15}\text{Np}$	0.11225	0.019680	-1.3E-4	-0.06	-5.1 E-04	-0.09
$d_3 \text{ Im}$	0.111645	0.019032	-0.0028	-1.24	-0.0171	-3.2
$^{54}\text{Fe}$	0.11577	0.019752	0.0154	6.9	0.0013	0.24

#### Case 2

$$f_{12} = 1.0$$

$$f_{23} = 0.6$$

Calculate numerically.

$\lambda = 0.005878$  and  $0.033537$  which corresponds to frequencies of  $99.9$  and  $238.7 \text{ cm}^{-1}$

	$\lambda$ (+)	$\lambda$ (-)	$ (\lambda'/\lambda)-1 $ (+)	v $\Delta v$	$(\lambda'/\lambda)-1$ (-)	v $\Delta v$
Na	0.033537	0.005878	-----	$238.7 \text{ cm}^{-1}$	-----	$99.9 \text{ cm}^{-1}$
$^{15}\text{Np}$	0.033529	0.005871	-1.2E-04	-0.03	-6.0 E-04	-0.06
$d_3 \text{ Im}$	0.033367	0.005701	-0.0025	-0.6	-0.015	-1.5
$^{54}\text{Fe}$	0.034588	0.005920	0.015	3.6	0.0035	0.356

## Case 3

$$f_{12} = 0.8$$

$$f_{23} = 0.6$$

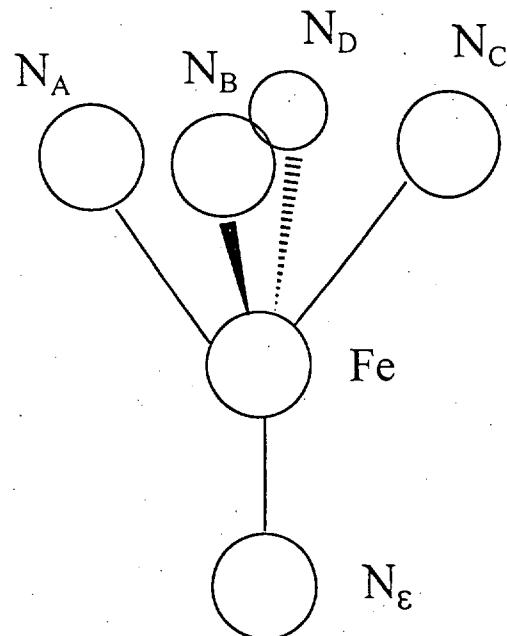
Calculate numerically.

$\lambda = 0.005221$  and  $0.030216$  which corresponds to frequencies of  $96.1$  and  $231.2 \text{ cm}^{-1}$ .

	$\lambda$ (+)	$\lambda$ (-)	$ (\lambda'/\lambda)-1 $ (+)	$\nu$ $\Delta\nu$	$(\lambda'/\lambda)-1$ (-)	$\nu$ $\Delta\nu$
Na	0.030216	0.005221	-----	$231.2 \text{ cm}^{-1}$	-----	$96.1 \text{ cm}^{-1}$
$^{15}\text{Np}$	0.030211	0.005212	-8.3E-05	-0.02	-8.6 E-04	-0.08
$d_3 \text{ Im}$	0.030017	0.005071	-0.0032	-0.74	-0.0145	-1.4
$^{54}\text{Fe}$	0.03109	0.005205	0.0143	3.30	0.0015	0.14

## Six-body model

The six-body model for a heme-axial-ligand complex has  $C_{4v}$  symmetry. We treat the porphyrin ring as consisting of four pyrrole nitrogens ( $N_A, N_B, N_C$ , and  $N_D$ ) of appropriate equivalent mass, one iron, and one axial atom with the equivalent mass of imidazole.



6 atoms x 3 degrees of freedom = 18 degrees of freedom

E	$2C_4$	$C_2$	$2\sigma_v$	$2\sigma_d$
18	2	-2	4	2

$$A_1 = (18 + 4 + -2 + 8 + 4)/8 = 4$$

$$A_2 = (18 + 4 - 2 - 8 - 4)/8 = 1$$

$$B_1 = (18 - 4 - 2 + 8 - 4)/8 = 2$$

$$B_2 = (18 - 4 - 2 - 8 + 4)/8 = 1$$

$$E = (36 + 0 + 4)/8 = 5$$

$$\Gamma = 4 A_1 + 1 A_2 + 2 B_1 + 1 B_2 + 5 E$$

For translation take away 1 E and 1 A<sub>1</sub>

For rotation take away 1 E and 1 A<sub>2</sub>

$$\Gamma_{\text{vib}} = 3 A_1 + 2 B_1 + 1 B_2 + 3 E \text{ or } 12 \text{ normal modes}$$

$$\Gamma_{\text{Raman}} = 3 A_1 + 2 B_1 + 1 B_2 + 3 E$$

Definition of internal coordinates

$\Delta d_A$  is the displacement of nitrogen A with respect to iron (etc.)

There are five total displacement internal coordinates.

$\Delta \theta_{AB}$  is the angle of N<sub>A</sub>-Fe-N<sub>B</sub> (etc.)

There are eight possible angles.

For the equatorial displacements

E	2C <sub>4</sub>	C <sub>2</sub>	2σ <sub>v</sub>	2σ <sub>d</sub>
4	0	0	4	0

$$\Gamma_{\text{EqDis}} = A_1 + B_1 + E$$

For the axial displacement

E	2C <sub>4</sub>	C <sub>2</sub>	2σ <sub>v</sub>	2σ <sub>d</sub>
1	1	1	1	1

$$\Gamma_{\text{AxDis}} = A_1$$

For the equatorial angles

E	2C <sub>4</sub>	C <sub>2</sub>	2σ <sub>v</sub>	2σ <sub>d</sub>
4	0	0	0	4

$$\Gamma_{\text{EqAngle}} = A_1 + B_2 + E$$

For the axial angles

E	2C <sub>4</sub>	C <sub>2</sub>	2σ <sub>v</sub>	2σ <sub>d</sub>

E	2C <sub>4</sub>	C <sub>2</sub>	2σ <sub>v</sub>	2σ <sub>d</sub>
---	-----------------	----------------	-----------------	-----------------

1	1	1	1	1
---	---	---	---	---

$$\Gamma_{\text{AxDis}} = A_1$$

For the equatorial angles

E	2C <sub>4</sub>	C <sub>2</sub>	2σ <sub>v</sub>	2σ <sub>d</sub>
4	0	0	0	4

$$\Gamma_{\text{EqAngle}} = A_1 + B_2 + E$$

For the axial angles

E	2C <sub>4</sub>	C <sub>2</sub>	2σ <sub>v</sub>	2σ <sub>d</sub>
4	0	0	4	0

$$\Gamma_{\text{Ax Angle}} = A_1 + B_1 + E$$

Projection operators to generate symmetry coordinates for normal modes.

$$P^{A_1} \Delta\theta_{AB} = 1/2(\Delta\theta_{AB} + \Delta\theta_{CD} + \Delta\theta_{BC} + \Delta\theta_{DA})$$

$$P^{A_1} \Delta d_A = 1/2(\Delta d_A + \Delta d_B + \Delta d_C + \Delta d_D)$$

$$P^{A_1} \Delta d_X = \Delta d_X$$

$$P^{A_1} \Delta\theta_{AX} = 1/2(\Delta\theta_{AX} + \Delta\theta_{BX} + \Delta\theta_{CX} + \Delta\theta_{DX})$$

The two type of angle deformation have at least one redundant angle so  $\Delta\theta_{DA}$  is removed.

$$P^{B_2} \Delta\theta_{AB} = 1/2(\Delta\theta_{AB} - \Delta\theta_{BC} + \Delta\theta_{CD} - \Delta\theta_{DA})$$

$$P^{B_2} \Delta d_A = \Delta d_A - \Delta d_B + \Delta d_C - \Delta d_D = 0$$

$$P^{B_2} \Delta d_X = 0$$

$$P^{B_2} \Delta\theta_{AX} = \Delta\theta_{AX} - \Delta\theta_{BX} + \Delta\theta_{CX} - \Delta\theta_{DX} = 0$$

$$P^{B_1} \Delta\theta_{AB} = 0$$

$$P^{B_1} \Delta d_A = 1/2(\Delta d_A - \Delta d_B + \Delta d_C - \Delta d_D)$$

$$P^{B_1} \Delta d_X = 0$$

$$P^{B_1} \Delta\theta_{AX} = 1/2(\Delta\theta_{AX} - \Delta\theta_{BX} + \Delta\theta_{CX} - \Delta\theta_{DX})$$

$$P^E \Delta\theta_{AB} = 1/2^{1/2}(\Delta\theta_{AB} - \Delta\theta_{CD})$$

$$P^E \Delta\theta_{BC} = 1/2^{1/2}(\Delta\theta_{BC} - \Delta\theta_{DA})$$

$$P^E \Delta d_A = 1/2^{1/2}(\Delta d_A - \Delta d_C)$$

$$P^E \Delta d_B = 1/2^{1/2}(\Delta d_B - \Delta d_D)$$

$$P^E \Delta d_X = 0$$

$$P^E \Delta\theta_{AX} = 1/2^{1/2}(\Delta\theta_{AX} - \Delta\theta_{CX})$$

$$P^E \Delta\theta_{BX} = 1/2^{1/2}(\Delta\theta_{BX} - \Delta\theta_{DX})$$

Based on these symmetry coordinates the U-matrix is:

$\Delta\theta_{AX}$	$\Delta\theta_{BX}$	$\Delta\theta_{CX}$	$\Delta\theta_{DX}$	$\Delta\theta_{AB}$	$\Delta\theta_{BC}$	$\Delta\theta_{CD}$	$\Delta d_A$	$\Delta d_B$	$\Delta d_C$	$\Delta d_D$	$\Delta d_x$
0	0	0	0	0	0	0	0	0	0	0	1
0	0	0	0	0	0	0	1/2	1/2	1/2	1/2	0
0	0	0	0	0	0	0	1/2	-1/2	1/2	-1/2	0
0	0	0	0	0	0	0	1/r2	0	-1/r2	0	0
0	0	0	0	0	0	0	0	1/r2	0	-1/r2	0
0	0	0	0	1/2	-1/2	1/2	0	0	0	0	0
0	0	0	0	1/r2	0	-1/r2	0	0	0	0	0
0	0	0	0	0	1/sr2	0	0	0	0	0	0
1/2	1/2	1/2	1/2	0	0	0	0	0	0	0	0
1/2	-1/2	1/2	-1/2	0	0	0	0	0	0	0	0
1/r2	0	-1/s2	0	0	0	0	0	0	0	0	0
0	1/r2	0	-1/r2	0	0	0	0	0	0	0	0

The f-matrix is:

$f_{\Delta\theta X}$	0	0	0	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	0	0	0	0
0	$f_{\Delta\theta X}$	0	0	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	0	0	0	0
0	0	$f_{\Delta\theta X}$	0	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	0	0	0	0
0	0	0	$f_{\Delta\theta X}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	0	0	0	0
$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta A}$	0	0	0	0	0	0	0
$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	0	$f_{\Delta\theta A}$	0	0	0	0	0	0
$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	0	0	$f_{\Delta\theta A}$	0	0	0	0	0
$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	$f_{\Delta\theta AX}$	0	0	0	$f_{\Delta\theta A}$	0	0	0	0
0	0	0	0	0	0	0	0	$f_{\Delta d A}$	0	0	$f_{\Delta d AX}$
0	0	0	0	0	0	0	0	0	$f_{\Delta d A}$	0	$f_{\Delta d AX}$
0	0	0	0	0	0	0	0	0	$f_{\Delta d A}$	0	$f_{\Delta d AX}$
0	0	0	0	0	0	0	0	0	0	$f_{\Delta d A}$	$f_{\Delta d AX}$
0	0	0	0	0	0	0	0	$f_{\Delta d AX}$	$f_{\Delta d AX}$	$f_{\Delta d AX}$	$f_{\Delta d X}$

The U' matrix is:

0	0	0	0	0	0	0	0	0	1/2	1/2	1/r2	0
0	0	0	0	0	0	0	0	0	1/2	-1/2	0	1/r2
0	0	0	0	0	0	0	0	0	1/2	1/2	-1/r2	0
0	0	0	0	0	0	0	0	0	1/2	-1/2	0	-1/r2
0	0	0	0	0	0	1/2	1/2	1/r2	0	0	0	0
0	0	0	0	0	0	1/2	-1/2	0	1/r2	0	0	0
0	0	0	0	0	0	1/2	1/2	-1/r2	0	0	0	0
0	0	0	0	0	0	1/2	-1/2	0	-1/r2	0	0	0
0	1/2	1/2	1/r2	0	0	0	0	0	0	0	0	0
0	1/2	-1/2	0	1/r2	0	0	0	0	0	0	0	0
0	1/2	1/2	-1/r2	0	0	0	0	0	0	0	0	0
0	1/2	-1/2	0	-1/r2	0	0	0	0	0	0	0	0
1	0	0	0	0	0	0	0	0	0	0	0	0

The F-matrix is:

$$\begin{matrix}
 f_{\Delta dX} & 2f_{\Delta dAX} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 2f_{\Delta dAX} & f_{\Delta dA} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & f_{\Delta dA} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & f_{\Delta dA} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & f_{\Delta dA} & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & f_{\Delta \theta A} & 0 & 0 & 0 & 4f_{\Delta \theta AX} & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & f_{\Delta \theta A} & 0 & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & f_{\Delta \theta A} & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & f_{\Delta \theta A} & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & f_{\Delta \theta A} & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 4f_{\Delta \theta AX} & 0 & 0 & 0 & 0 & f_{\Delta \theta X} & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & f_{\Delta \theta X} & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & f_{\Delta \theta X} & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & f_{\Delta \theta X}
 \end{matrix}$$

We wish to calculate  $\mathbf{G} = \mathbf{UgU}'$  where  $\mathbf{g}$  is:

$$\begin{matrix}
 \Delta\theta_{AX} & \Delta\theta_{BX} & \Delta\theta_{CX} & \Delta\theta_{DX} & \Delta\theta_{AB} & \Delta\theta_{BC} & \Delta\theta_{CD} & \Delta\theta_{DA} & \Delta d_A & \Delta d_B & \Delta d_C & \Delta d_D & \Delta d_X \\
 \Delta\theta_{AX} & g_{\phi\phi}^3 & g_{\phi\phi}^2 & g_{\phi\phi}^2 & g_{\phi\phi}^2 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^1 \\
 \Delta\theta_{BX} & g_{\phi\phi}^2 & g_{\phi\phi}^3 & g_{\phi\phi}^2 & g_{\phi\phi}^2 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{r\phi}^1 & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^2 \\
 \Delta\theta_{CX} & g_{\phi\phi}^2 & g_{\phi\phi}^2 & g_{\phi\phi}^3 & g_{\phi\phi}^2 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^2 \\
 \Delta\theta_{DX} & g_{\phi\phi}^2 & g_{\phi\phi}^2 & g_{\phi\phi}^2 & g_{\phi\phi}^3 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^2 & g_{r\phi}^2 \\
 \Delta\theta_{AB} & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^3 & g_{\phi\phi}^2 & g_{\phi\phi}^2 & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^1 \\
 \Delta\theta_{BC} & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^2 & g_{\phi\phi}^3 & g_{\phi\phi}^2 & g_{r\phi}^1 & g_{r\phi}^2 & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^1 \\
 \Delta\theta_{CD} & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^2 & g_{\phi\phi}^2 & g_{\phi\phi}^3 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^2 & g_{r\phi}^2 & g_{r\phi}^1 \\
 \Delta\theta_{DA} & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^1 & g_{\phi\phi}^2 & g_{\phi\phi}^2 & g_{\phi\phi}^2 & g_{r\phi}^3 & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^2 \\
 \Delta d_A & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^1 & g_{\pi}^2 & g_{\pi}^1 & g_{\pi}^1 & g_{\pi}^1 \\
 \Delta d_B & g_{r\phi}^1 & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^2 & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^1 & g_{\pi}^1 & g_{\pi}^2 & g_{\pi}^1 & g_{\pi}^1 \\
 \Delta d_C & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^2 & g_{r\phi}^2 & g_{r\phi}^1 & g_{\pi}^1 & g_{\pi}^1 & g_{\pi}^2 & g_{\pi}^1 \\
 \Delta d_D & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^2 & g_{r\phi}^2 & g_{\pi}^1 & g_{\pi}^1 & g_{\pi}^1 & g_{\pi}^2 \\
 \Delta d_X & g_{r\phi}^2 & g_{r\phi}^2 & g_{r\phi}^2 & g_{r\phi}^2 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^1 & g_{r\phi}^1 & g_{\pi}^1 & g_{\pi}^1 & g_{\pi}^1 & g_{\pi}^1
 \end{matrix}$$

The A<sub>1</sub> symmetry block is:

$$\begin{array}{cccc}
 f_{\Delta dX} & 2f_{\Delta AX} & 0 & 0 \\
 2f_{\Delta dAX} & f_{\Delta dA} & 0 & 0 \\
 0 & 0 & f_{\Delta \theta A} & 4f_{\Delta \theta AX} \\
 0 & 0 & 4f_{\Delta \theta AX} & f_{\Delta \theta X} \\
 \\ 
 g_{rr}^2 & 2g_{rr}^{-1} & 2g_{r\phi}^{-1} & 2g_{r\phi}^2 \\
 2g_{rr}^{-1} & g_{rr}^{-2} + 3g_{rr}^{-1} & 2(g_{r\phi}^{-1} + g_{r\phi}^2) & g_{r\phi}^{-2} + 3g_{r\phi}^{-1} \\
 2g_{r\phi}^{-1} & 2(g_{r\phi}^{-2} + g_{r\phi}^{-1}) & g_{\phi\phi}^{-3} + 3g_{\phi\phi}^{-2} & 4g_{\phi\phi}^{-1} \\
 g_{r\phi}^2 & g_{r\phi}^{-2} + 3g_{r\phi}^{-1} & 4g_{\phi\phi}^{-1} & g_{\phi\phi}^{-3} + 3g_{\phi\phi}^{-2}
 \end{array}$$

We can eliminate the redundant coordinate at this point. It is ΔθA.

$$\begin{array}{cccc}
 f_{\Delta dX} & 2f_{\Delta AX} & 0 & \\
 2f_{\Delta dAX} & f_{\Delta dA} & 0 & \\
 0 & 0 & f_{\Delta \theta X} & \\
 \\ 
 1.67 & 0 & 0 & \\
 0 & 1.67 & 0 & \\
 0 & 0 & 0.83 & \\
 \\ 
 g_{rr}^2 & 2g_{rr}^{-1} & 2g_{r\phi}^{-2} & \\
 2g_{rr}^{-1} & g_{rr}^{-2A} + 3g_{rr}^{-1} & g_{r\phi}^{-2A} + 3g_{r\phi}^{-1} & \\
 2g_{r\phi}^{-2} & g_{r\phi}^{-2A} + 3g_{r\phi}^{-1} & g_{\phi\phi}^{-3} + 3g_{\phi\phi}^{-2} & \\
 \\ 
 0.03255 & -0.0068 & -0.01668 & \\
 -0.0068 & 0.01584 & -0.008389 & \\
 -0.01668 & -0.008389 & 0.0149 &
 \end{array}$$

for  $f_{\Delta dX}$

$$g_{rr}^{-2} = \mu_1 + \mu_2 = 1/68 + 1/56 = (0.01470) + (0.01785) = 0.03255$$

for  $f_{\Delta dA}$

$$g_{rr}^{-2A} = \mu_3 + \mu_2 = 1/122 + 1/56 = (0.008196) + (0.01785) = 0.02604$$

for  $f_{\Delta dA}$  and  $f_{\Delta dX}$

$$g_{rr}^{-1} = \mu_2 \cos\phi = 1/56 \cos(101.04) = (0.01785)(-0.191) = -0.0034$$

for  $f_{\Delta dA}$  and  $f_{\Delta \theta X}$

$$g_{r\phi}^{-1} = \rho_{13} \mu_3 \sin\phi \cos\tau = 0$$

for  $f_{\Delta dX}$  and  $f_{\Delta \theta X}$

$$g\tau\phi^2 = -\rho_{23} \mu_2 \sin\phi = -(1/2.1)(1/56) \sin(101.04) = -(0.4761)(0.01785)(0.981) = -0.008340$$

for  $f_{\Delta dA}$  and  $f_{\Delta \theta X}$

$$g\tau\phi^2 A = -\rho_{12} \mu_2 \sin\phi = -(1/2.088)(1/56) \sin(101.04) = -(0.4789)(0.01785)(0.981) = -0.008389$$

for  $f_{\Delta \theta X}$

$$g\phi\tau^3 = \rho_{12}^2 \mu_1 + \rho_{23}^2 \mu_3 + (\rho_{12}^2 + \rho_{23}^2 - 2\rho_{12}\rho_{23}\cos\phi)\mu_2 = (1/122)(1/2.088)^2 + (1/68)(1/2.1)^2 - (0.008196)(0.2293) + (0.01470)(0.2267) + (0.2293 + 0.2267 - 0.4561(-0.191))(0.01785) = 0.01490$$

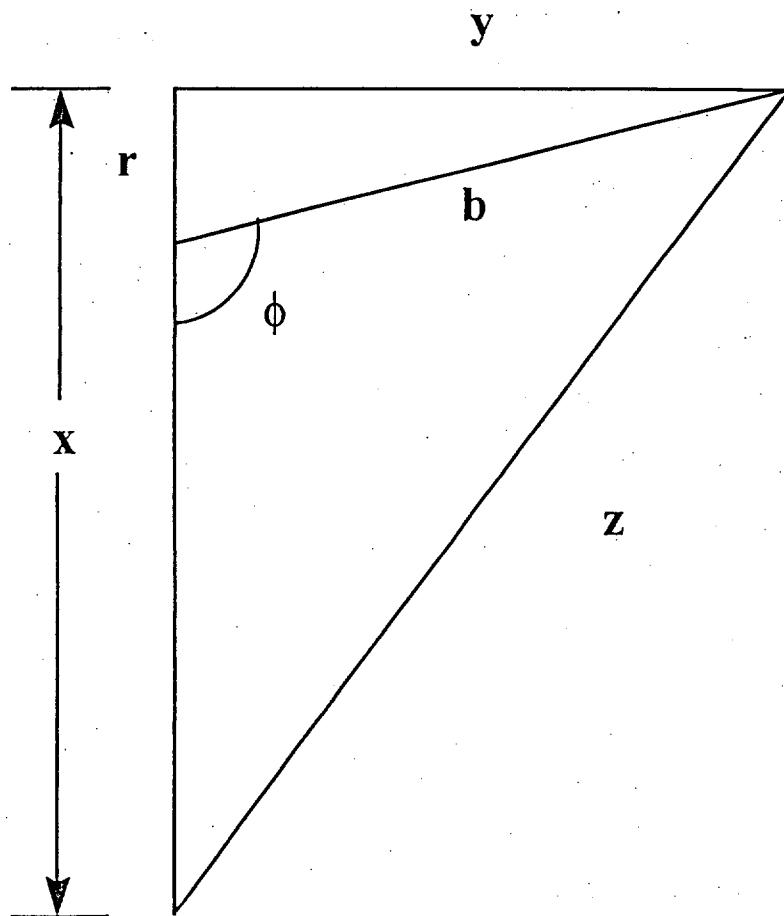
$g\phi\tau^2 = 0$  because the dihedral angle between the planes is 90 degrees and therefore  $\cos \tau$  is zero.

Make the following identifications.

1 = pyrrole nitrogens

2 = iron

3 = axial ligand



$r$  is the iron doming coordinate (i.e. relative to the plane of the pyrrole nitrogens)

Note that  $\phi = \pi - \eta$  where  $\eta = \arccos r/b$  or  $\arctan y/r$

In fact the length  $x$  and  $z$  are irrelevant for determination of the angle.

So if  $y = 2.05 \text{ \AA}$  and  $r = 0.4 \text{ \AA}$  then  $y/r = 5.125$

$$\eta = 78.959$$

$$\phi = 101.04$$

$$\cos \phi = -0.191$$

$$\sin \phi = 0.981$$

Assuming that  $x = 2.5 \text{ \AA}$  then  $b = \sqrt{(r^2 + y^2)} = 2.088 \text{ \AA}$   
 $x = 2.5 \text{ \AA}$  means that the iron axial ligand bond is  $2.1 \text{ \AA}$ .

$$\begin{array}{ccc} f_{\Delta dX} & 2f_{\Delta dAX} & 0 \\ 2f_{\Delta dAX} & f_{\Delta dA} & 0 \\ 0 & 0 & f_{\Delta \theta X} \end{array}$$

$$\begin{array}{ccc} g_{rr}^2 & 2g_{rr'}^{-1} & 2g_{r\phi}^2 \\ 2g_{rr'}^{-1} & g_{rr}^{2A} + 3g_{rr'}^{-1} & g_{r\phi}^{2A} \\ g_{r\phi}^2 & g_{r\phi}^{2A} & g_{\phi\phi}^3 \end{array}$$

The FG Matrix is the product of these two.

### The 38-body model

Although the addition of interaction force constants and torsional force constants could be used to refine the force field we find that this is not warranted. We present no quantitative conclusions for such modeling because there is no justification for introducing a greater number of parameters based on the present experimental data. Qualitatively, the effect of introducing out-of-plane force constants (3) is to introduce significant mixing of porphyrin doming and ruffling coordinates into the normal modes that involve the iron-axial-ligand stretch. If the motion of the axial ligand is purely axial then the ruffling coordinate does not contribute by symmetry, however, ruffling can cause mixing of nominally in-plane internal coordinates with off-axis tilting of the ligand. By including these internal coordinates the pyrrole  $^{15}\text{N}$  isotope effect begins to become much larger than the values of  $0.2 - 0.4 \text{ cm}^{-1}$  for modes A and B shown in Table 6. The experimental data rule out larger isotope effects. The recent observation that proteins enforce distortions along the low frequency saddling, ruffling, and doming vibrational modes suggests a significant constraint exists on these degrees of freedom. (2) It is possible that torsional degrees of freedom in crystal or solution may be significantly different than in proteins. Our results suggest that porphyrin torsions would almost necessarily lead to an isotope effect for  $^{15}\text{N}$  substituted porphyrins in myoglobin, contrary to experimental observation. (1) The view of the iron-axial-ligand mode that emerges from our analysis is

of a structurally rigid porphyrin with significant ligand bending in a mode that is approximately 50% an iron-ligand stretch.

The experimental evidence for a protein role in enforcing distortions of metalloporphyrins has been presented. The globins (including H93G myoglobin studied here) have a known domed structure that is enforced by the protein and the interplay between protein structure in the R and T state of hemoglobins. The domed heme of the T state lowers the iron-histidine mode frequency. The lowering of the frequency may occur due to off-axis tilting that increases anharmonic coupling to low frequency bending modes. It is remarkable that the frequency lowering in the T state of two different types of hemoglobins are similar in spite of different mechanisms for enforcing the T state. (4) This observations suggests a key role for the heme geometry. Peroxidases have an enforced saddled conformation of the heme and this may affect the mixing of the porphyrin out-of-plane modes and axial-ligand stretching. (2) These examples serve to indicate the importance of structural constraints on out-of-plane torsional and wagging internal coordinates.

The 38-body model applied here uses the SVIB program available from the Quantum Chemistry Program Exchange. The parameters used in a representative input file are given below. No hydrogens are included. The carbons in the imidazole ring were pseudo-atoms with a mass of 13 a.m.u. (carbon plus hydrogen). The eight peripheral substituents in the  $\beta$ -position are pseudo-atoms with a mass 24 a.m.u. This number was chosen for consistency with the three-body and six-body models and represents an average number for the methyl, vinyl, and propionate substituents of heme. The parameters that were varied systematically in the model calculations are shown in bold in the input file. The masses were altered to model the isotope effect. For example, the pyrrole nitrogen isotope effect was modeled by replacing:

XMATRX	3	2.027970	-0.077730	-0.115320	14.000000	NA	_____
XMATRX	4	-0.043490	-2.101650	-0.245570	14.000000	NB	_____
XMATRX	5	-2.058600	-0.028090	-0.430350	14.000000	NC	_____
XMATRX	6	0.015760	1.998760	-0.300210	14.000000	ND	_____

with

XMATRX	3	2.027970	-0.077730	-0.115320	15.000000	NA	_____
XMATRX	4	-0.043490	-2.101650	-0.245570	15.000000	NB	_____
XMATRX	5	-2.058600	-0.028090	-0.430350	15.000000	NC	_____
XMATRX	6	0.015760	1.998760	-0.300210	15.000000	ND	_____

#### Sample input file for SVIB

```
*****
SPV VERSION 100
HEADER
```

COMPND im\_heme\_model

ISOTOP Na

REMARK File created by franzen on Tue Apr 29 15:14:36 PDT 1997

COMMENT # Xcoord Ycoord Zcoord AtomMass Atomtype

XMATRX 1 -0.031260 -0.049220 -0.056900 56.000000 FE  
 XMATRX 2 -0.191350 -0.019700 2.027460 14.000000 NX  
 XMATRX 3 2.027970 -0.077730 -0.115320 14.000000 NA  
 XMATRX 4 -0.043490 -2.101650 -0.245570 14.000000 NB  
 XMATRX 5 -2.058600 -0.028090 -0.430350 14.000000 NC  
 XMATRX 6 0.015760 1.998760 -0.300210 14.000000 ND  
 XMATRX 7 2.819680 -1.194340 -0.082010 12.000000 C1  
 XMATRX 8 4.196500 -0.793220 0.032530 12.000000 C2  
 XMATRX 9 4.220420 0.556410 0.070200 12.000000 C3  
 XMATRX 10 2.858520 1.012740 -0.063270 12.000000 C4  
 XMATRX 11 2.372150 -2.507660 -0.137740 12.000000 C5  
 XMATRX 12 -1.153830 -2.902510 -0.324430 12.000000 C6  
 XMATRX 13 -0.738120 -4.282420 -0.267610 12.000000 C7  
 XMATRX 14 0.612230 -4.295360 -0.228850 12.000000 C8  
 XMATRX 15 1.050500 -2.925130 -0.202400 12.000000 C9  
 XMATRX 16 -2.462830 -2.451130 -0.456250 12.000000 C1\_0  
 XMATRX 17 -2.843870 1.086340 -0.558280 12.000000 C1\_1  
 XMATRX 18 -4.219420 0.675860 -0.682070 12.000000 C1\_2  
 XMATRX 19 -4.245320 -0.679980 -0.707970 12.000000 C1\_3  
 XMATRX 20 -2.882560 -1.123260 -0.518070 12.000000 C1\_4  
 XMATRX 21 -2.387760 2.400640 -0.575130 12.000000 C1\_5  
 XMATRX 22 1.134190 2.791280 -0.256680 12.000000 C1\_6  
 XMATRX 23 0.730870 4.171100 -0.348380 12.000000 C1\_7  
 XMATRX 24 -0.613730 4.197390 -0.524750 12.000000 C1\_8  
 XMATRX 25 -1.066260 2.825680 -0.461960 12.000000 C1\_9  
 XMATRX 26 2.445540 2.338450 -0.138220 24.000000 C2\_0  
 XMATRX 27 0.526410 0.750420 2.860610 13.000000 C2\_9  
 XMATRX 28 0.148630 0.492430 4.122550 15.000000 NY  
 XMATRX 29 -0.811120 -0.442810 4.093050 13.000000 C3\_0  
 XMATRX 30 -1.027190 -0.768080 2.761230 13.000000 C3\_1  
 XMATRX 31 -1.630010 -5.472330 -0.311000 24.000000 C2\_1  
 XMATRX 32 5.431970 1.413520 0.169420 24.000000 C2\_2  
 XMATRX 33 1.643290 5.345680 -0.255670 24.000000 C2\_3  
 XMATRX 34 -5.391020 1.593650 -0.759170 24.000000 C2\_4  
 XMATRX 35 -1.439490 5.283660 -0.715290 24.000000 C2\_5  
 XMATRX 36 -5.324950 -1.521150 -0.865910 24.000000 C2\_6  
 XMATRX 37 1.501520 -5.496590 -0.175850 24.000000 C2\_7  
 XMATRX 38 5.357390 -1.726490 0.154370 24.000000 C2\_8

COMMENT # Atom Numbers of connected atoms

CONECT 1 4 3 6 5 2 0  
 CONECT 2 1 27 30 0 0 0  
 CONECT 3 1 7 10 0 0 0  
 CONECT 4 1 12 15 0 0 0  
 CONECT 5 1 17 20 0 0 0  
 CONECT 6 1 22 25 0 0 0  
 CONECT 7 3 8 11 0 0 0  
 CONECT 8 7 9 38 0 0 0  
 CONECT 9 8 10 32 0 0 0  
 CONECT 10 3 9 26 0 0 0  
 CONECT 11 7 15 0 0 0 0  
 CONECT 12 4 13 16 0 0 0  
 CONECT 13 12 14 31 0 0 0  
 CONECT 14 13 15 37 0 0 0  
 CONECT 15 4 11 14 0 0 0  
 CONECT 16 12 20 0 0 0 0  
 CONECT 17 5 18 21 0 0 0

CONECT	18	17	19	34	0	0	0
CONECT	19	18	20	36	0	0	0
CONECT	20	5	16	19	0	0	0
CONECT	21	17	25	0	0	0	0
CONECT	22	6	23	26	0	0	0
CONECT	23	22	24	33	0	0	0
CONECT	24	23	25	35	0	0	0
CONECT	25	6	21	24	0	0	0
CONECT	26	10	22	0	0	0	0
CONECT	27	2	28	0	0	0	0
CONECT	28	27	29	0	0	0	0
CONECT	29	28	30	0	0	0	0
CONECT	30	2	29	0	0	0	0
CONECT	31	13	0	0	0	0	0
CONECT	32	9	0	0	0	0	0
CONECT	33	23	0	0	0	0	0
CONECT	34	18	0	0	0	0	0
CONECT	35	24	0	0	0	0	0
CONECT	36	19	0	0	0	0	0
CONECT	37	14	0	0	0	0	0
CONECT	38	8	0	0	0	0	0
COMMENT	#	ForceConstant	LowerBound	Upperbound	DescriptiveLabel		
DIF_FC	1	6.980	0.000	0.000	K(C10C14)		
DIF_FC	2	6.980	0.000	0.000	K(C10C6_)		
DIF_FC	3	5.430	0.000	0.000	K(C11C12)		
DIF_FC	4	6.980	0.000	0.000	K(C11C15)		
DIF_FC	5	5.640	0.000	0.000	K(C11NC_)		
DIF_FC	6	7.120	0.000	0.000	K(C12C13)		
DIF_FC	7	4.200	0.000	0.000	K(C12C24)		
DIF_FC	8	5.430	0.000	0.000	K(C13C14)		
DIF_FC	9	4.200	0.000	0.000	K(C13C26)		
DIF_FC	10	5.640	0.000	0.000	K(C14NC_)		
DIF_FC	11	6.980	0.000	0.000	K(C15C19)		
DIF_FC	12	5.430	0.000	0.000	K(C16C17)		
DIF_FC	13	6.980	0.000	0.000	K(C16C20)		
DIF_FC	14	5.640	0.000	0.000	K(C16ND_)		
DIF_FC	15	7.120	0.000	0.000	K(C17C18)		
DIF_FC	16	4.200	0.000	0.000	K(C17C23)		
DIF_FC	17	5.430	0.000	0.000	K(C18C19)		
DIF_FC	18	4.200	0.000	0.000	K(C18C25)		
DIF_FC	19	5.640	0.000	0.000	K(C19ND_)		
DIF_FC	20	5.430	0.000	0.000	K(C1_C2_)		
DIF_FC	21	6.980	0.000	0.000	K(C1_C5_)		
DIF_FC	22	5.640	0.000	0.000	K(C1_NA_)		
DIF_FC	23	6.980	0.000	0.000	K(C20C4_)		
DIF_FC	24	4.200	0.000	0.000	K(C21C7_)		
DIF_FC	25	4.200	0.000	0.000	K(C22C3_)		
DIF_FC	26	4.200	0.000	0.000	K(C27C8_)		
DIF_FC	27	4.200	0.000	0.000	K(C28C2_)		
DIF_FC	28	7.560	0.000	0.000	K(C29NX_)		
DIF_FC	29	6.060	0.000	0.000	K(C29NY_)		
DIF_FC	30	7.120	0.000	0.000	K(C2_C3_)		
DIF_FC	31	8.130	0.000	0.000	K(C30C31)		
DIF_FC	32	5.300	0.000	0.000	K(C30NY_)		
DIF_FC	33	5.070	0.000	0.000	K(C31NX_)		
DIF_FC	34	5.430	0.000	0.000	K(C3_C4_)		
DIF_FC	35	5.640	0.000	0.000	K(C4_NA_)		
DIF_FC	36	6.980	0.000	0.000	K(C5_C9_)		
DIF_FC	37	5.430	0.000	0.000	K(C6_C7_)		
DIF_FC	38	5.640	0.000	0.000	K(C6_NB_)		

DIF_FC	39	7.120	0.000	0.000	K(C7_C8_)
DIF_FC	40	5.430	0.000	0.000	K(C8_C9_)
DIF_FC	41	5.640	0.000	0.000	K(C9_NB_)
DIF_FC	42	1.000	0.000	0.000	K(FE_NA_)
DIF_FC	43	1.000	0.000	0.000	K(FE_NB_)
DIF_FC	44	1.000	0.000	0.000	K(FE_NC_)
DIF_FC	45	1.000	0.000	0.000	K(FE_ND_)
DIF_FC	46	1.000	0.000	0.000	K(FE_NX_)
DIF_FC	47	0.830	0.000	0.000	H(C10C14C13)
DIF_FC	48	0.830	0.000	0.000	H(C10C14NC_)
DIF_FC	49	0.830	0.000	0.000	H(C10C6_C7_)
DIF_FC	50	0.830	0.000	0.000	H(C10C6_NB_)
DIF_FC	51	1.370	0.000	0.000	H(C11C12C13)
DIF_FC	52	1.200	0.000	0.000	H(C11C12C24)
DIF_FC	53	1.100	0.000	0.000	H(C11C15C19)
DIF_FC	54	1.620	0.000	0.000	H(C11NC_C14)
DIF_FC	55	0.300	0.000	0.000	H(C11NC_FE_)
DIF_FC	56	0.830	0.000	0.000	H(C12C11C15)
DIF_FC	57	1.370	0.000	0.000	H(C12C11NC_)
DIF_FC	58	1.370	0.000	0.000	H(C12C13C14)
DIF_FC	59	1.200	0.000	0.000	H(C12C13C26)
DIF_FC	60	1.200	0.000	0.000	H(C13C12C24)
DIF_FC	61	1.370	0.000	0.000	H(C13C14NC_)
DIF_FC	62	1.100	0.000	0.000	H(C14C10C6_)
DIF_FC	63	1.200	0.000	0.000	H(C14C13C26)
DIF_FC	64	0.300	0.000	0.000	H(C14NC_FE_)
DIF_FC	65	0.830	0.000	0.000	H(C15C11NC_)
DIF_FC	66	0.830	0.000	0.000	H(C15C19C18)
DIF_FC	67	0.830	0.000	0.000	H(C15C19ND_)
DIF_FC	68	1.370	0.000	0.000	H(C16C17C18)
DIF_FC	69	1.200	0.000	0.000	H(C16C17C23)
DIF_FC	70	1.100	0.000	0.000	H(C16C20C4_)
DIF_FC	71	1.620	0.000	0.000	H(C16ND_C19)
DIF_FC	72	0.300	0.000	0.000	H(C16ND_FE_)
DIF_FC	73	0.830	0.000	0.000	H(C17C16C20)
DIF_FC	74	1.370	0.000	0.000	H(C17C16ND_)
DIF_FC	75	1.370	0.000	0.000	H(C17C18C19)
DIF_FC	76	1.200	0.000	0.000	H(C17C18C25)
DIF_FC	77	1.200	0.000	0.000	H(C18C17C23)
DIF_FC	78	1.370	0.000	0.000	H(C18C19ND_)
DIF_FC	79	1.200	0.000	0.000	H(C19C18C25)
DIF_FC	80	0.300	0.000	0.000	H(C19ND_FE_)
DIF_FC	81	1.200	0.000	0.000	H(C1_C2_C28)
DIF_FC	82	1.370	0.000	0.000	H(C1_C2_C3_)
DIF_FC	83	1.100	0.000	0.000	H(C1_C5_C9_)
DIF_FC	84	1.620	0.000	0.000	H(C1_NA_C4_)
DIF_FC	85	0.300	0.000	0.000	H(C1_NA_FE_)
DIF_FC	86	0.830	0.000	0.000	H(C20C16ND_)
DIF_FC	87	0.830	0.000	0.000	H(C20C4_C3_)
DIF_FC	88	0.830	0.000	0.000	H(C20C4_NA_)
DIF_FC	89	1.200	0.000	0.000	H(C21C7_C6_)
DIF_FC	90	1.200	0.000	0.000	H(C21C7_C8_)
DIF_FC	91	1.200	0.000	0.000	H(C22C3_C2_)
DIF_FC	92	1.200	0.000	0.000	H(C22C3_C4_)
DIF_FC	93	1.200	0.000	0.000	H(C27C8_C7_)
DIF_FC	94	1.200	0.000	0.000	H(C27C8_C9_)
DIF_FC	95	1.200	0.000	0.000	H(C28C2_C3_)
DIF_FC	96	1.200	0.000	0.000	H(C29NX_C31)
DIF_FC	97	1.200	0.000	0.000	H(C29NX_FE_)

DIF\_FC 98 1.200 0.000 0.000 H(C29NY\_C30)  
 DIF\_FC 99 0.830 0.000 0.000 H(C2\_C1\_C5\_ )  
 DIF\_FC 100 1.370 0.000 0.000 H(C2\_C1\_NA\_ )  
 DIF\_FC 101 1.370 0.000 0.000 H(C2\_C3\_C4\_ )  
 DIF\_FC 102 1.620 0.000 0.000 H(C30C31NX\_ )  
 DIF\_FC 103 1.370 0.000 0.000 H(C31C30NY\_ )  
 DIF\_FC 104 0.300 0.000 0.000 H(C31NX\_FE\_ )  
 DIF\_FC 105 1.370 0.000 0.000 H(C3\_C4\_NA\_ )  
 DIF\_FC 106 0.300 0.000 0.000 H(C4\_NA\_FE\_ )  
 DIF\_FC 107 1.370 0.000 0.000 H(C5\_C1\_NA\_ )  
 DIF\_FC 108 0.830 0.000 0.000 H(C5\_C9\_C8\_ )  
 DIF\_FC 109 0.830 0.000 0.000 H(C5\_C9\_NB\_ )  
 DIF\_FC 110 1.370 0.000 0.000 H(C6\_C7\_C8\_ )  
 DIF\_FC 111 1.620 0.000 0.000 H(C6\_NB\_C9\_ )  
 DIF\_FC 112 0.300 0.000 0.000 H(C6\_NB\_FE\_ )  
 DIF\_FC 113 1.370 0.000 0.000 H(C7\_C6\_NB\_ )  
 DIF\_FC 114 1.370 0.000 0.000 H(C7\_C8\_C9\_ )  
 DIF\_FC 115 1.370 0.000 0.000 H(C8\_C9\_NB\_ )  
 DIF\_FC 116 0.300 0.000 0.000 H(C9\_NB\_FE\_ )  
 DIF\_FC 117 0.250 0.000 0.000 H(NA\_FE\_NB\_ )  
 DIF\_FC 118 0.250 0.000 0.000 H(NA\_FE\_ND\_ )  
 DIF\_FC 119 0.300 0.000 0.000 H(NA\_FE\_NX\_ )  
 DIF\_FC 120 0.250 0.000 0.000 H(NB\_FE\_NC\_ )  
 DIF\_FC 121 0.300 0.000 0.000 H(NB\_FE\_NX\_ )  
 DIF\_FC 122 0.250 0.000 0.000 H(NC\_FE\_ND\_ )  
 DIF\_FC 123 0.300 0.000 0.000 H(NC\_FE\_NX\_ )  
 DIF\_FC 124 0.300 0.000 0.000 H(ND\_FE\_NX\_ )  
 DIF\_FC 125 1.400 0.000 0.000 H(NX\_C29NY\_ )

COMMENT	#	Typ	IC1	IC2	At1	At2	At3	At4	At5	At6	At7	At8	ScaleF	FC#
Not.used.														
FMATRX	1	1	1	1	16-	20	0	0	0	0	0	0	1.0000	1
K(C10C14)														
FMATRX	2	1	2	2	12-	16	0	0	0	0	0	0	1.0000	2
K(C10C6_ )														
FMATRX	3	1	3	3	17-	18	0	0	0	0	0	0	1.0000	3
K(C11C12)														
FMATRX	4	1	4	4	17-	21	0	0	0	0	0	0	1.0000	4
K(C11C15)														
FMATRX	5	1	5	5	5-	17	0	0	0	0	0	0	1.0000	5
K(C11NC_ )														
FMATRX	6	1	6	6	18-	19	0	0	0	0	0	0	1.0000	6
K(C12C13)														
FMATRX	7	1	7	7	18-	34	0	0	0	0	0	0	1.0000	7
K(C12C24)														
FMATRX	8	1	8	8	19-	20	0	0	0	0	0	0	1.0000	8
K(C13C14)														
FMATRX	9	1	9	9	19-	36	0	0	0	0	0	0	1.0000	9
K(C13C26)														
FMATRX	10	1	10	10	5-	20	0	0	0	0	0	0	1.0000	10
K(C14NC_ )														
FMATRX	11	1	11	11	21-	25	0	0	0	0	0	0	1.0000	11
K(C15C19)														
FMATRX	12	1	12	12	22-	23	0	0	0	0	0	0	1.0000	12
K(C16C17)														
FMATRX	13	1	13	13	22-	26	0	0	0	0	0	0	1.0000	13
K(C16C20)														
FMATRX	14	1	14	14	6-	22	0	0	0	0	0	0	1.0000	14
K(C16ND_ )														

FMATRX	15	1	15	15	23-	24	0	0	0	0	0	0	1.0000	15
K(C17C18)														
FMATRX	16	1	16	16	23-	33	0	0	0	0	0	0	1.0000	16
K(C17C23)														
FMATRX	17	1	17	17	24-	25	0	0	0	0	0	0	1.0000	17
K(C18C19)														
FMATRX	18	1	18	18	24-	35	0	0	0	0	0	0	1.0000	18
K(C18C25)														
FMATRX	19	1	19	19	6-	25	0	0	0	0	0	0	1.0000	19
K(C19ND_)														
FMATRX	20	1	20	20	7-	8	0	0	0	0	0	0	1.0000	20
K(C1_C2_)														
FMATRX	21	1	21	21	7-	11	0	0	0	0	0	0	1.0000	21
K(C1_C5_)														
FMATRX	22	1	22	22	3-	7	0	0	0	0	0	0	1.0000	22
K(C1_NA_)														
FMATRX	23	1	23	23	10-	26	0	0	0	0	0	0	1.0000	23
K(C20C4_)														
FMATRX	24	1	24	24	13-	31	0	0	0	0	0	0	1.0000	24
K(C21C7_)														
FMATRX	25	1	25	25	9-	32	0	0	0	0	0	0	1.0000	25
K(C22C3_)														
FMATRX	26	1	26	26	14-	37	0	0	0	0	0	0	1.0000	26
K(C27C8_)														
FMATRX	27	1	27	27	8-	38	0	0	0	0	0	0	1.0000	27
K(C28C2_)														
FMATRX	28	1	28	28	2-	27	0	0	0	0	0	0	1.0000	28
K(C29NX_)														
FMATRX	29	1	29	29	27-	28	0	0	0	0	0	0	1.0000	29
K(C29NY_)														
FMATRX	30	1	30	30	8-	9	0	0	0	0	0	0	1.0000	30
K(C2_C3_)														
FMATRX	31	1	31	31	29-	30	0	0	0	0	0	0	1.0000	31
K(C30C31)														
FMATRX	32	1	32	32	28-	29	0	0	0	0	0	0	1.0000	32
K(C30NY_)														
FMATRX	33	1	33	33	2-	30	0	0	0	0	0	0	1.0000	33
K(C31NX_)														
FMATRX	34	1	34	34	9-	10	0	0	0	0	0	0	1.0000	34
K(C3_C4_)														
FMATRX	35	1	35	35	3-	10	0	0	0	0	0	0	1.0000	35
K(C4_NA_)														
FMATRX	36	1	36	36	11-	15	0	0	0	0	0	0	1.0000	36
K(C5_C9_)														
FMATRX	37	1	37	37	12-	13	0	0	0	0	0	0	1.0000	37
K(C6_C7_)														
FMATRX	38	1	38	38	4-	12	0	0	0	0	0	0	1.0000	38
K(C6_NB_)														
FMATRX	39	1	39	39	13-	14	0	0	0	0	0	0	1.0000	39
K(C7_C8_)														
FMATRX	40	1	40	40	14-	15	0	0	0	0	0	0	1.0000	40
K(C8_C9_)														
FMATRX	41	1	41	41	4-	15	0	0	0	0	0	0	1.0000	41
K(C9_NB_)														
FMATRX	42	1	42	42	1-	3	0	0	0	0	0	0	1.0000	42
K(FE_NA_)														
FMATRX	43	1	43	43	1-	4	0	0	0	0	0	0	1.0000	43
K(FE_NB_)														
FMATRX	44	1	44	44	1-	5	0	0	0	0	0	0	1.0000	44
K(FE_NC_)														

FMATRX	45	1	45	45	1-	6	0	0	0	0	0	1.0000	45
K(FE_ND_)													
FMATRX	46	1	46	46	1-	2	0	0	0	0	0	1.0000	46
K(FE_NX_)													
FMATRX	47	2	47	47	16-	20-	19	0	0	0	0	1.0000	47
H(C10C14C13)													
FMATRX	48	2	48	48	5-	20-	16	0	0	0	0	1.0000	48
H(C10C14NC_)													
FMATRX	49	2	49	49	13-	12-	16	0	0	0	0	1.0000	49
H(C10C6_C7_)													
FMATRX	50	2	50	50	4-	12-	16	0	0	0	0	1.0000	50
H(C10C6_NB_)													
FMATRX	51	2	51	51	17-	18-	19	0	0	0	0	1.0000	51
H(C11C12C13)													
FMATRX	52	2	52	52	17-	18-	34	0	0	0	0	1.0000	52
H(C11C12C24)													
FMATRX	53	2	53	53	17-	21-	25	0	0	0	0	1.0000	53
H(C11C15C19)													
FMATRX	54	2	54	54	17-	5-	20	0	0	0	0	1.0000	54
H(C11NC_C14)													
FMATRX	55	2	55	55	1-	5-	17	0	0	0	0	1.0000	55
H(C11NC_FE_)													
FMATRX	56	2	56	56	18-	17-	21	0	0	0	0	1.0000	56
H(C12C11C15)													
FMATRX	57	2	57	57	5-	17-	18	0	0	0	0	1.0000	57
H(C12C11NC_)													
FMATRX	58	2	58	58	18-	19-	20	0	0	0	0	1.0000	58
H(C12C13C14)													
FMATRX	59	2	59	59	18-	19-	36	0	0	0	0	1.0000	59
H(C12C13C26)													
FMATRX	60	2	60	60	19-	18-	34	0	0	0	0	1.0000	60
H(C13C12C24)													
FMATRX	61	2	61	61	5-	20-	19	0	0	0	0	1.0000	61
H(C13C14NC_)													
FMATRX	62	2	62	62	12-	16-	20	0	0	0	0	1.0000	62
H(C14C10C6_)													
FMATRX	63	2	63	63	20-	19-	36	0	0	0	0	1.0000	63
H(C14C13C26)													
FMATRX	64	2	64	64	1-	5-	20	0	0	0	0	1.0000	64
H(C14NC_FE_)													
FMATRX	65	2	65	65	5-	17-	21	0	0	0	0	1.0000	65
H(C15C11NC_)													
FMATRX	66	2	66	66	21-	25-	24	0	0	0	0	1.0000	66
H(C15C19C18)													
FMATRX	67	2	67	67	6-	25-	21	0	0	0	0	1.0000	67
H(C15C19ND_)													
FMATRX	68	2	68	68	22-	23-	24	0	0	0	0	1.0000	68
H(C16C17C18)													
FMATRX	69	2	69	69	22-	23-	33	0	0	0	0	1.0000	69
H(C16C17C23)													
FMATRX	70	2	70	70	10-	26-	22	0	0	0	0	1.0000	70
H(C16C20C4_)													
FMATRX	71	2	71	71	22-	6-	25	0	0	0	0	1.0000	71
H(C16ND_C19)													
FMATRX	72	2	72	72	1-	6-	22	0	0	-0	0	1.0000	72
H(C16ND_FE_)													
FMATRX	73	2	73	73	23-	22-	26	0	0	0	0	1.0000	73
H(C17C16C20)													
FMATRX	74	2	74	74	6-	22-	23	0	0	0	0	1.0000	74
H(C17C16ND_)													

FMATRX	75	2	75	75	23-	24-	25	0	0	0	0	0	1.0000	75
H(C17C18C19)														
FMATRX	76	2	76	76	23-	24-	35	0	0	0	0	0	1.0000	76
H(C17C18C25)														
FMATRX	77	2	77	77	24-	23-	33	0	0	0	0	0	1.0000	77
H(C18C17C23)														
FMATRX	78	2	78	78	6-	25-	24	0	0	0	0	0	1.0000	78
H(C18C19ND_)														
FMATRX	79	2	79	79	25-	24-	35	0	0	0	0	0	1.0000	79
H(C19C18C25)														
FMATRX	80	2	80	80	1-	6-	25	0	0	0	0	0	1.0000	80
H(C19ND_FE_)														
FMATRX	81	2	81	81	7-	8-	38	0	0	0	0	0	1.0000	81
H(C1_C2_C28)														
FMATRX	82	2	82	82	7-	8-	9	0	0	0	0	0	1.0000	82
H(C1_C2_C3_)														
FMATRX	83	2	83	83	7-	11-	15	0	0	0	0	0	1.0000	83
H(C1_C5_C9_)														
FMATRX	84	2	84	84	7-	3-	10	0	0	0	0	0	1.0000	84
H(C1_NA_C4_)														
FMATRX	85	2	85	85	1-	3-	7	0	0	0	0	0	1.0000	85
H(C1_NA_FE_)														
FMATRX	86	2	86	86	6-	22-	26	0	0	0	0	0	1.0000	86
H(C20C16ND_)														
FMATRX	87	2	87	87	9-	10-	26	0	0	0	0	0	1.0000	87
H(C20C4_C3_)														
FMATRX	88	2	88	88	3-	10-	26	0	0	0	0	0	1.0000	88
H(C20C4_NA_)														
FMATRX	89	2	89	89	12-	13-	31	0	0	0	0	0	1.0000	89
H(C21C7_C6_)														
FMATRX	90	2	90	90	14-	13-	31	0	0	0	0	0	1.0000	90
H(C21C7_C8_)														
FMATRX	91	2	91	91	8-	9-	32	0	0	0	0	0	1.0000	91
H(C22C3_C2_)														
FMATRX	92	2	92	92	10-	9-	32	0	0	0	0	0	1.0000	92
H(C22C3_C4_)														
FMATRX	93	2	93	93	13-	14-	37	0	0	0	0	0	1.0000	93
H(C27C8_C7_)														
FMATRX	94	2	94	94	15-	14-	37	0	0	0	0	0	1.0000	94
H(C27C8_C9_)														
FMATRX	95	2	95	95	9-	8-	38	0	0	0	0	0	1.0000	95
H(C28C2_C3_)														
FMATRX	96	2	96	96	27-	2-	30	0	0	0	0	0	1.0000	96
H(C29NX_C31)														
FMATRX	97	2	97	97	1-	2-	27	0	0	0	0	0	1.0000	97
H(C29NX_FE_)														
FMATRX	98	2	98	98	27-	28-	29	0	0	0	0	0	1.0000	98
H(C29NY_C30)														
FMATRX	99	2	99	99	8-	7-	11	0	0	0	0	0	1.0000	99
H(C2_C1_C5_)														
FMATRX	100	2	100	100	3-	7-	8	0	0	0	0	0	1.0000	100
H(C2_C1_NA_)														
FMATRX	101	2	101	101	8-	9-	10	0	0	0	0	0	1.0000	101
H(C2_C3_C4_)														
FMATRX	102	2	102	102	2-	30-	29	0	0	0	0	0	1.0000	102
H(C30C31NX_)														
FMATRX	103	2	103	103	28-	29-	30	0	0	0	0	0	1.0000	103
H(C31C30NY_)														
FMATRX	104	2	104	104	1-	2-	30	0	0	0	0	0	1.0000	104
H(C31NX_FE_)														

FMATRX	105	2	105	105	3-	10-	9	0	0	0	0	0	1.0000	105
H(C3_C4_NA_)														
FMATRX	106	2	106	106	1-	3-	10	0	0	0	0	0	1.0000	106
H(C4_NA_FE_)														
FMATRX	107	2	107	107	3-	7-	11	0	0	0	0	0	1.0000	107
H(C5_C1_NA_)														
FMATRX	108	2	108	108	11-	15-	14	0	0	0	0	0	1.0000	108
H(C5_C9_C8_)														
FMATRX	109	2	109	109	4-	15-	11	0	0	0	0	0	1.0000	109
H(C5_C9_NB_)														
FMATRX	110	2	110	110	12-	13-	14	0	0	0	0	0	1.0000	110
H(C6_C7_C8_)														
FMATRX	111	2	111	111	12-	4-	15	0	0	0	0	0	1.0000	111
H(C6_NB_C9_)														
FMATRX	112	2	112	112	1-	4-	12	0	0	0	0	0	1.0000	112
H(C6_NB_FE_)														
FMATRX	113	2	113	113	4-	12-	13	0	0	0	0	0	1.0000	113
H(C7_C6_NB_)														
FMATRX	114	2	114	114	13-	14-	15	0	0	0	0	0	1.0000	114
H(C7_C8_C9_)														
FMATRX	115	2	115	115	4-	15-	14	0	0	0	0	0	1.0000	115
H(C8_C9_NB_)														
FMATRX	116	2	116	116	1-	4-	15	0	0	0	0	0	1.0000	116
H(C9_NB_FE_)														
FMATRX	117	2	117	117	3-	1-	4	0	0	0	0	0	1.0000	117
H(NA_FE_NB_)														
FMATRX	118	2	118	118	3-	1-	6	0	0	0	0	0	1.0000	118
H(NA_FE_ND_)														
FMATRX	119	2	119	119	2-	1-	3	0	0	0	0	0	1.0000	119
H(NA_FE_NX_)														
FMATRX	120	2	120	120	4-	1-	5	0	0	0	0	0	1.0000	120
H(NB_FE_NC_)														
FMATRX	121	2	121	121	2-	1-	4	0	0	0	0	0	1.0000	121
H(NB_FE_NX_)														
FMATRX	122	2	122	122	5-	1-	6	0	0	0	0	0	1.0000	122
H(NC_FE_ND_)														
FMATRX	123	2	123	123	2-	1-	5	0	0	0	0	0	1.0000	123
H(NC_FE_NX_)														
FMATRX	124	2	124	124	2-	1-	6	0	0	0	0	0	1.0000	124
H(ND_FE_NX_)														
FMATRX	125	2	125	125	2-	27-	28	0	0	0	0	0	1.0000	125
H(NX_C29NY_)														

BEGINOBSFREQS

1 Isotope 125 labels

## Na

1 Mode001.....	1	0
2 Mode002.....	2	0
3 Mode003.....	3	0
4 Mode004.....	4	0
5 Mode005.....	5	0
6 Mode006.....	6	0
7 Mode007.....	7	0
8 Mode008.....	8	0
9 Mode009.....	9	0
10 Mode010.....	10	0
11 Mode011.....	11	0
12 Mode012.....	12	0
13 Mode013.....	13	0
14 Mode014.....	14	0

15 Mode015.....	15	0
16 Mode016.....	16	0
17 Mode017.....	17	0
18 Mode018.....	18	0
19 Mode019.....	19	0
20 Mode020.....	20	0
21 Mode021.....	21	0
22 Mode022.....	22	0
23 Mode023.....	23	0
24 Mode024.....	24	0
25 Mode025.....	25	0
26 Mode026.....	26	0
27 Mode027.....	27	0
28 Mode028.....	28	0
29 Mode029.....	29	0
30 Mode030.....	30	0
31 Mode031.....	31	0
32 Mode032.....	32	0
33 Mode033.....	33	0
34 Mode034.....	34	0
35 Mode035.....	35	0
36 Mode036.....	36	0
37 Mode037.....	37	0
38 Mode038.....	38	0
39 Mode039.....	39	0
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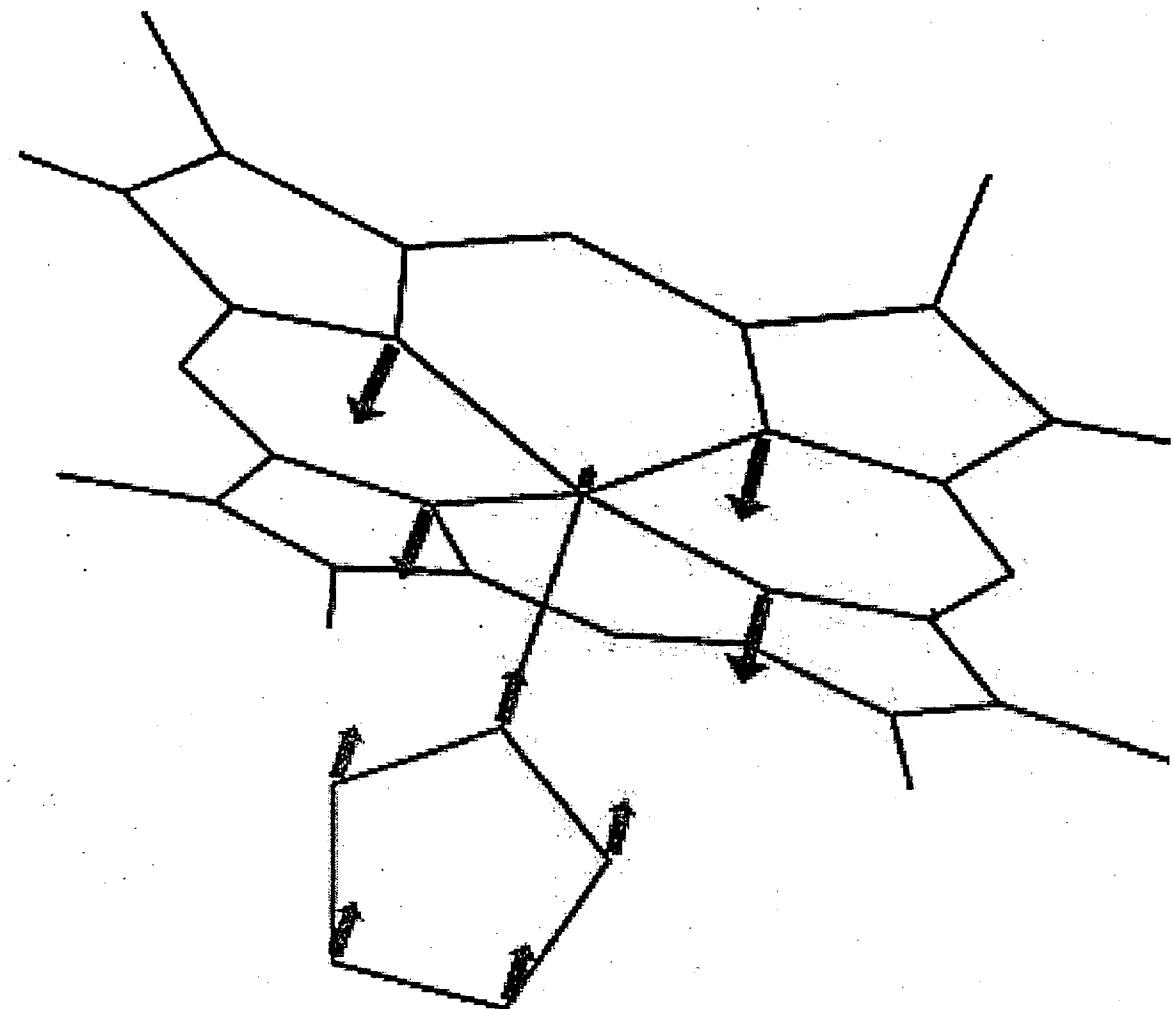
ENDOBSFREQS

END

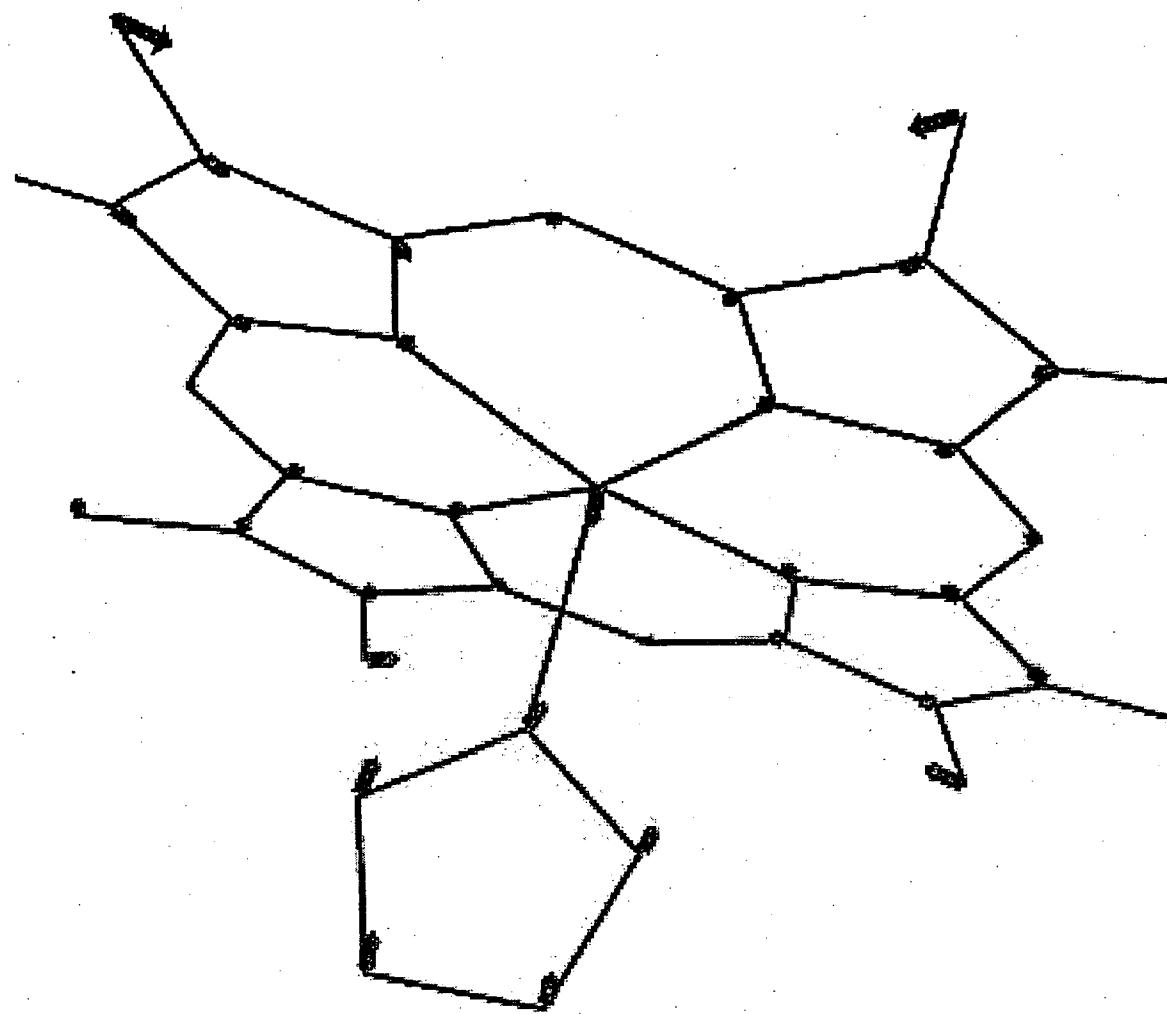
\*\*\*\*\*

Mode A'' Sombrero 102.4 cm<sup>-1</sup>

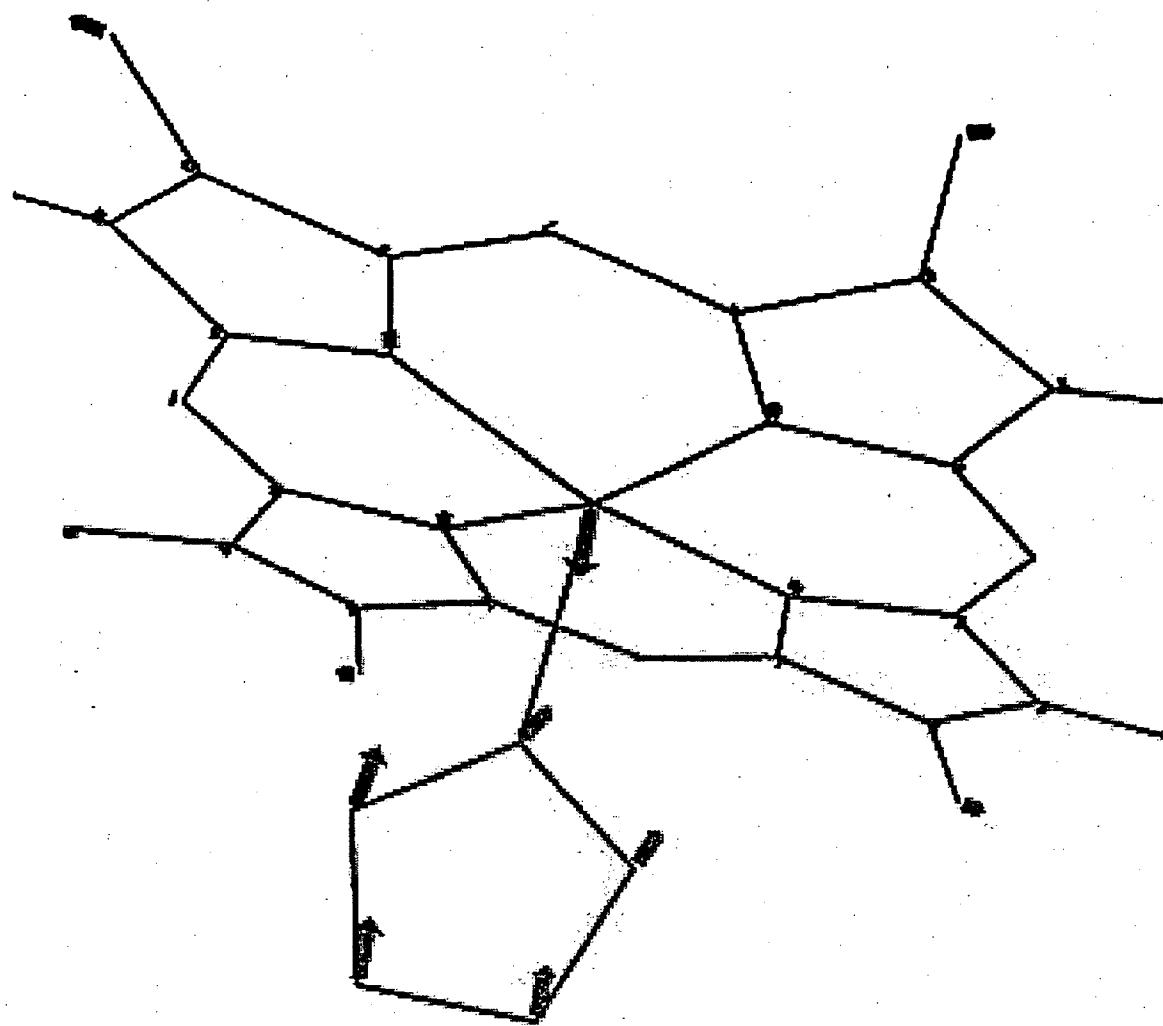
19 % H(ND\_FE\_NX)  
19 % H(NA\_FE\_NX)  
18 % H(NC\_FE\_NX)  
13 % H(NB\_FE\_NX)  
14 % K(FE\_NX)



Mode A' Doming/Breathing  $205.7\text{ cm}^{-1}$   
20 % K(FE\_NX\_)

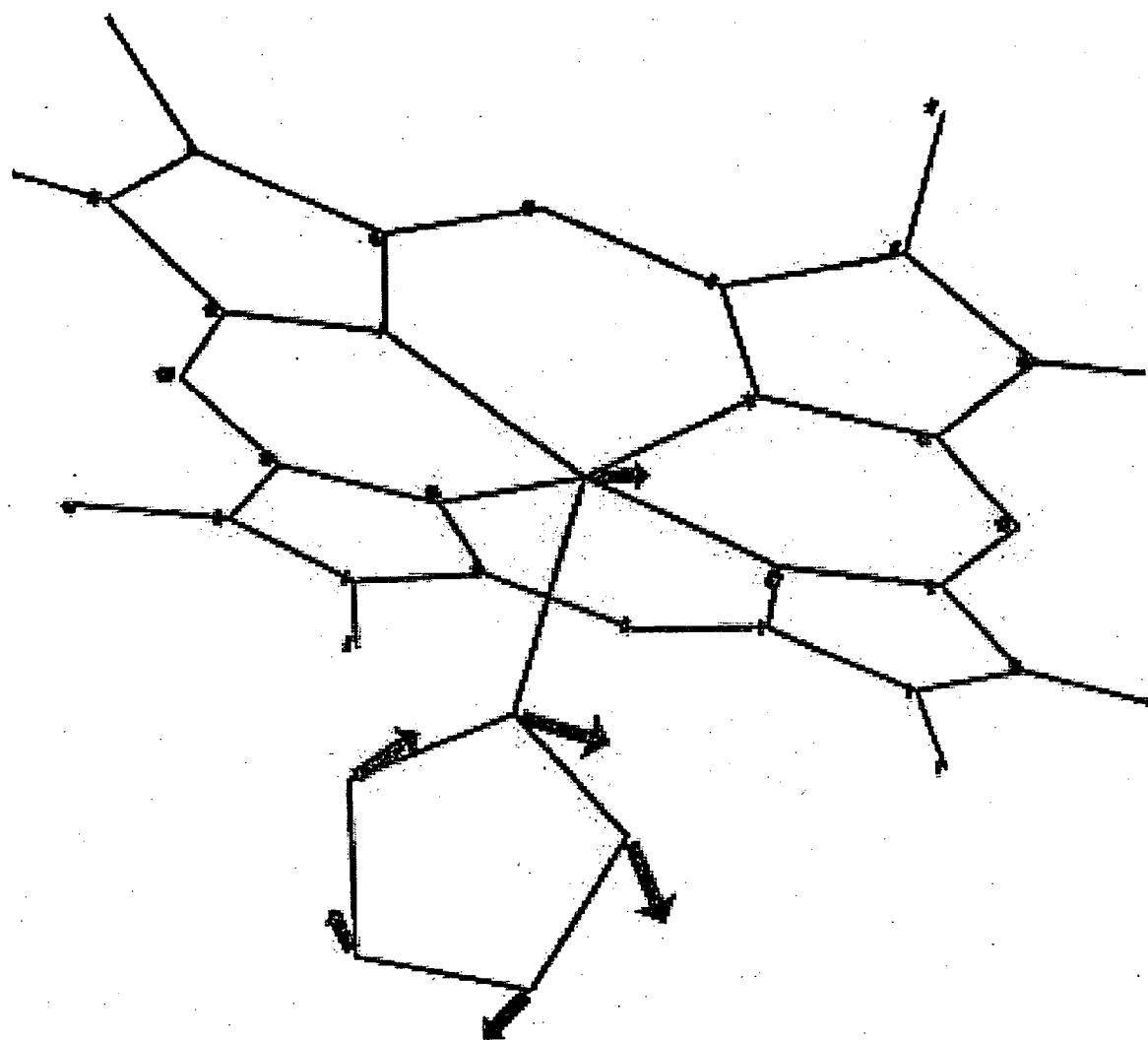


Mode A Doming/Axial Stretching  $219.7 \text{ cm}^{-1}$   
53 % K(FE\_NX\_)



Mode B Tilting 229.8 cm<sup>-1</sup>

24 & H(C29NX\_FE\_ )  
7 & H(C31NX\_FE\_ )  
8 & K(FE\_ND\_ )  
8 & K(FE\_NA\_ )  
7 & K(FE\_NC\_ )  
7 & K(FE\_NB\_ )



## Fitting Procedure

The fits of the experimental line shape to two Gaussians is illustrated below for both H93G(Im) and H93G(2-Me Im). The fitting function used is a sum of two normalized Gaussians:

$$F(\omega) = \frac{A_1}{\sqrt{2\pi}\sigma_1} \exp\left\{-\frac{(\omega - \omega_1)^2}{2\sigma_1^2}\right\} + \frac{A_2}{\sqrt{2\pi}\sigma_2} \exp\left\{-\frac{(\omega - \omega_2)^2}{2\sigma_2^2}\right\}$$

Table A5. Fits to experimental data in the wavenumber range from 180 - 250 cm<sup>-1</sup>. A two Gaussian model was used to fit the data.

Ligand	A <sub>1</sub>	$\omega_1$	$\sigma_1$	A <sub>2</sub>	$\omega_2$	$\sigma_2$
Im	7560.01	223.533	6.18485	28702.7	227.662	14.4517
d3-Im	7016.74	222.562	6.01776	29509.4	226.495	14.4745
2-Me Im	145366	215.856	11.5496	36045.3	238.416	12.7852
d5 2-Me Im	117985	210.207	10.9328	60419	233.362	13.2352

The fits to two gaussians demonstrate the problem identified in the manuscript that the shift in frequency for the H93G(Im) substitution is too small compared to that expected for a two body model, while that for H93G(2-Me Im) is too large compared to the same model. The fits are illustrated below.

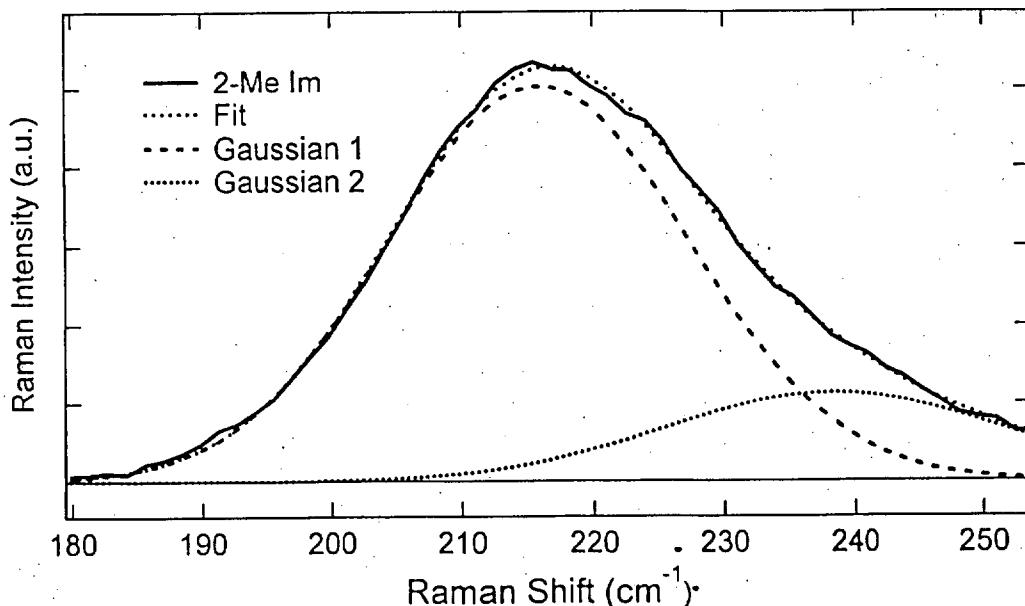


Figure A1. Fit of the data for H93G(2-Me Im) to two Gaussians. Both the fit to the data and individual gaussian components are shown.

The fits do not change qualitatively for the perdeutero isotopomer of 2-methyl imidazole as shown below in Figure A2.

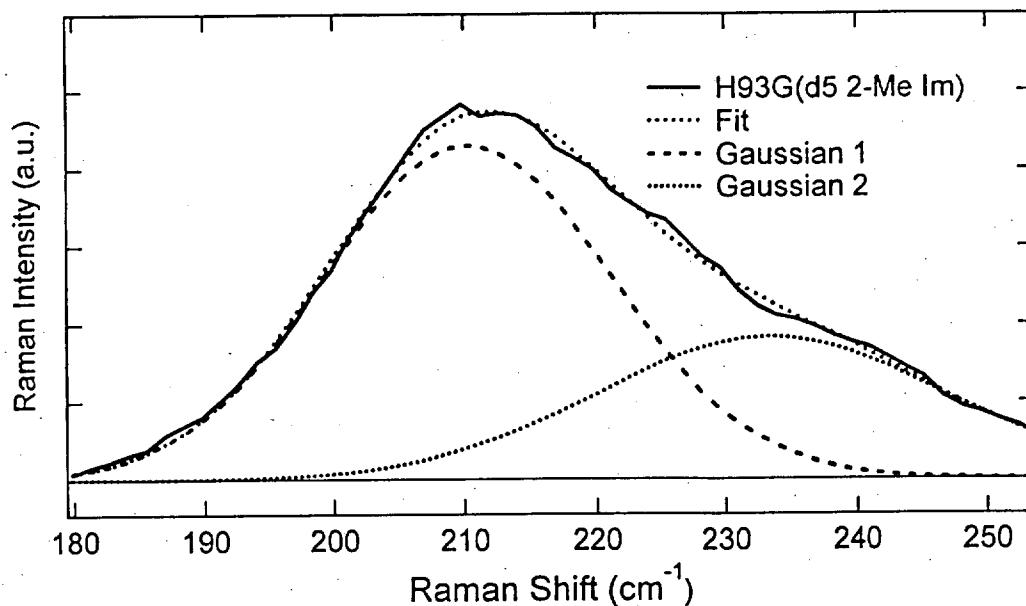


Figure A2. Fit of the axial-ligand out-of-plane band of H93G(d5 2-Me Im) to a two Gaussian model.

A similar situation applies for the fits to H93G(Im) and the d3 isotopomer as shown in Figures A3 and A4 below.

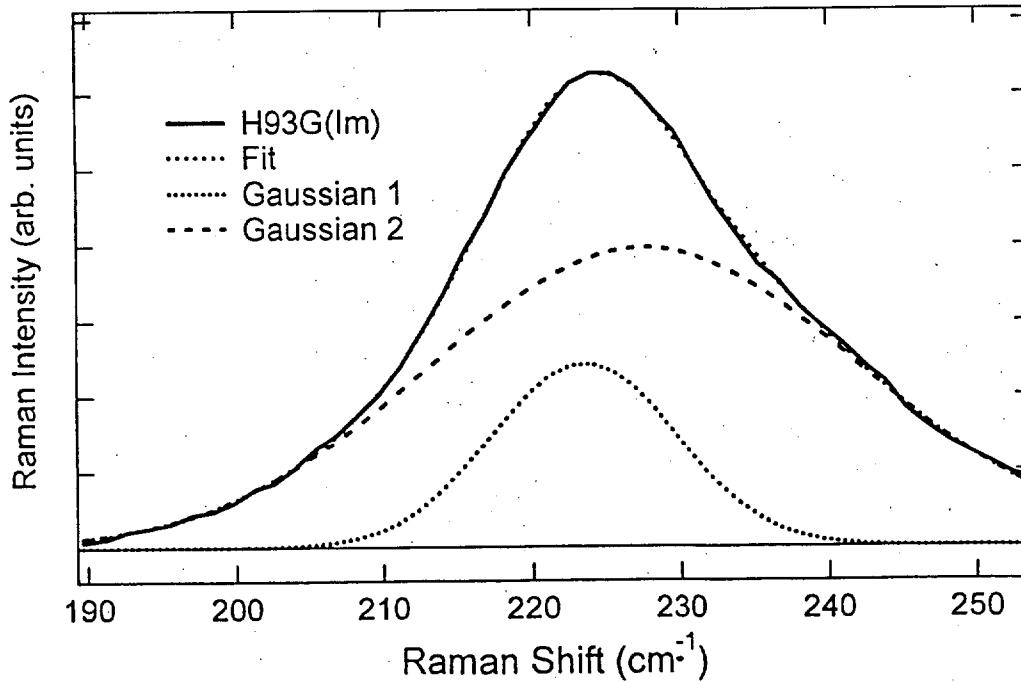


Figure A3. The fit to the axial-ligand out-of-plane band for the imidazole adduct of H93G to a two Gaussian model. Note that the extremely broad peak centered at ca.

$227\text{ cm}^{-1}$  does not have a physical interpretation, however, the shift in that peak is within  $0.2\text{ cm}^{-1}$  of the magnitude of the smaller peak in the data for the d3-Im isotopomer (Figure A4).

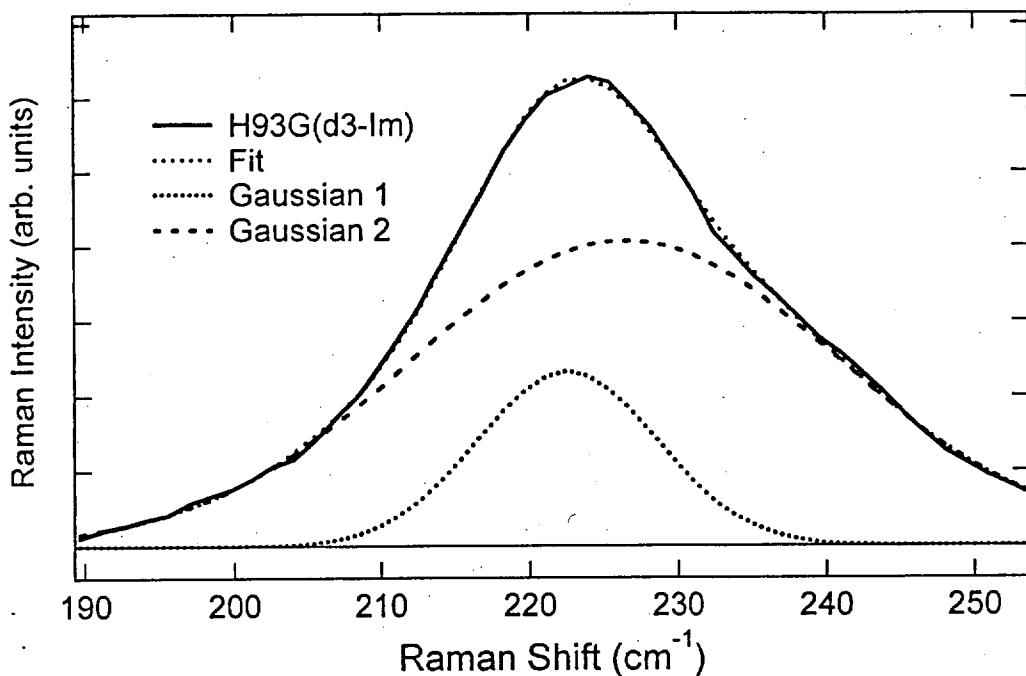


Figure A4. The fit of the axial-ligand out-of-plane band for H93G(d3-Im) to a two Gaussian model.

#### References

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