

Wax Inhibition by Comb-Like Polymers: Support of the Incorporation-Perturbation Mechanism from Molecular Dynamics Simulations

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Supporting Information

S1. Force field for *n*-alkanes (oil, wax, tail of inhibitor)

The methylene (CH₂) and methyl (CH₃) groups in *n*-alkanes (oil solvent, wax, and long alkyl tails of inhibitors) were treated as united atoms, that is, each CH₂ or CH₃ unit was treated as a single neutral pseudoatom. The force field (FF) for these united atoms (Table S1) was mostly taken from the SKS (Siepmann-Karanorn-Smit) FF.¹⁻³ The bond-stretching force constant, which is not in the original SKS FF, was taken from the AMBER FF.^{4,5} The torsion potential of the SKS FF had been taken from the OPLS (Optimized Potentials for Liquid Simulation) FF.⁶ The nonbonded interactions between 1-2, 1-3, and 1-4 neighbors were excluded, and other nonbonded interactions were truncated by a cubic spline cutoff function defined to be 1 for $r < 13.8$ Å, 0 for $r > 14.3$ Å, and decrease smoothly in between. The off-diagonal vdW parameters between CH₂ and CH₃ are given by the geometric means of diagonal ones:

$$R_0^{AB} = \sqrt{R_0^A R_0^B} \text{ and } D_0^{AB} = \sqrt{D_0^A D_0^B}.$$

Table S1. Force field used for *n*-alkanes.

$E^{vdW}(R) = D_0 \left[\left(\frac{R_0}{R} \right)^{12} - 2 \left(\frac{R_0}{R} \right)^6 \right], \quad E^{bond}(R) = \frac{1}{2} K_R (R - R_0)^2, \quad E^{angle}(\theta) = \frac{1}{2} K_\theta (\theta - \theta_0)^2, \quad E^{torsion}(\phi) = \sum_n \frac{1}{2} V_n [1 - d_n \cos(n\phi)]$							
E^{vdW}	CH ₂	R_0^a	4.4113	D_0^b	0.09339		
	CH ₃	R_0^a	4.4113	D_0^b	0.2265		
E^{bond}	CH ₂ -CH ₂	R_0^a	1.54	K_b^c	520		
	CH ₂ -CH ₃	R_0^a	1.54	K_b^c	520		
E^{angle}	CH ₂ -CH ₂ -CH ₂	θ_0^d	114	K_θ^e	124.19		
	CH ₂ -CH ₂ -CH ₃	θ_0^d	114	K_θ^e	124.19		
$E^{torsion}$	CH ₂ -CH ₂ -CH ₂ -CH ₂	$V_1^b(d_1)$	1.4109 (-1)	$V_2^b(d_2)$	-0.271 (1)	$V_3^b(d_3)$	2.787 (-1)
	CH ₂ -CH ₂ -CH ₂ -CH ₃	$V_1^b(d_1)$	1.4109 (-1)	$V_2^b(d_2)$	-0.271 (1)	$V_3^b(d_3)$	2.787 (-1)

^aIn Å. This R_0 has been converted from the original literature value σ by the relation $R_0 = \sqrt[6]{2}\sigma$, since in the original literature E^{vdW} was defined as

$$E^{vdW}(R) = 4\epsilon \left[\left(\frac{\sigma}{R} \right)^{12} - \left(\frac{\sigma}{R} \right)^6 \right].$$

^bIn kcal/mol. This D_0 (kcal/mol) has been converted from the original literature value $\epsilon (= D_0/k$ where k is the Boltzmann constant; K) by the relation by $\epsilon \times k$. ^cIn kcal/mol/Å². ^dIn degrees. ^eIn kcal/mol/rad².

S2. Force field for the backbone of inhibitor

The FF for the acrylate backbone (Table S2 and Figures S1-S2) was mostly taken from the OPLS FF developed for liquid methyl acetate^{7,8} except several following modifications.

Atomic charges of the acrylate backbone (-CHCOO-) were refined by *ab initio* quantum mechanics (QM) calculation on a model compound representing the acrylate group, (CH₃)₂CHCOOCH₂CH₃ (Figure S1a). The calculation was done at HF/6-31G** level using Jaguar v3.5.^{9,10} The atomic charges were fitted to reproduce the electrostatic potential derived from the wavefunction by the CHELP method (ESP-fitted charges).¹¹⁻¹⁵ Because the charges of the methyl groups, which will be replaced by neutral united atoms in simulations, were close to zero within 0.1 |e| (Figure S1a), they were constrained to zero in a new ESP fitting to determine the charges of the rest of the molecule (Figure S1b). The resulted atomic charges were reasonably similar to those used in the OPLS FF (Figure S1c).

The force constants for bond stretching, angle bending, and inversion, which are not in the original OPLS FF, were taken from the AMBER FF (Figures S2b-c).^{4,5}

The torsion FF was refined to reproduce the torsion potential curves obtained from HF/6-31G** QM calculations on small model compounds. The QM torsion potential was calculated by optimizing geometrical parameters at a series of fixed values of a torsion angle of interest.

A model compound $(\text{CH}_3\text{CH}_2)_2\text{CHCOOCH}_3$ (Figures S3a-b) was employed to obtain the potential along the $\text{CH}_2\text{-CH-C=O}$ torsion angle. The conformation with H-C-C=O at 180° (Figure S3b) is 0.3 kcal/mol lower than the conformation with H-C-C=O at 0° (Figure S3a), and these two minima are separated by a barrier of 4.2 kcal/mol. This QM potential for the $\text{CH}_2\text{-CH-C=O}$ torsion angle was reproduced very well (Figure 6a) by choosing appropriate torsion parameters.

A model compound $(\text{CH}_3)_2\text{CHCOOCH}_2\text{CH}_3$ (Figure S3c) was employed for the potential along the CH-C-O-CH_2 torsion angle. Since both conformations with H-C-C=O at 0° and at 180° are shown above to have significant populations and not to convert to each other easily with a barrier of 4.2 kcal/mol, the CH-C-O-CH_2 torsion potential was calculated for both cases: Figures S3c and S4c show the case where H-C-C=O is at 0° and Figures S3d and S4d show the case where H-C-C=O is at 180° . In any case, the *trans* conformation where C-C-O-C is at 180° (Figure S4c) is more stable by 10-16 kcal/mol than the *cis* conformation where C-C-O-C is at 0° (Figures S4c-d), obviously because of the repulsion between the $-\text{CH}_2-$ (or $-\text{CH}_3$) groups located close to each other in the *cis* conformation. This CH-C-O-CH_2 torsion potential was reproduced reasonably well by choosing appropriate torsion parameters (Figures S4c-d and Table S2).

Two model compounds, $(\text{CH}_3)_2\text{CHCOOCH}_2\text{CH}_3$ (Figure S3c) and $\text{CH}_3\text{COOCH}_2\text{CH}_2\text{CH}_3$ (Figures S3d-e) were used to calculate the torsion potential along $\text{C-O-CH}_2\text{-CH}_2$ (or $-\text{CH}_3$). These two compounds gave essentially the same $\text{C-O-CH}_2\text{-CH}_2$ torsion potential quantum mechanically (Figure S4b). This torsion is relatively free (with a barrier less than 1.2 kcal/mol) unless the $-\text{CH}_2-$ or $-\text{CH}_3$ group approaches the carbonyl oxygen ($=\text{O}$) too close in the *cis* conformation with $\text{C-O-CH}_2\text{-CH}_2$ at 0° , where the energy is 8 kcal/mol higher than the *trans* conformation with $\text{C-O-CH}_2\text{-CH}_2$ at 180° . The FF calculation with appropriately chosen parameters reproduces the overall shape of this QM torsion potential, but gives too high energy for the *cis* conformation (11 kcal/mol higher than the *trans* conformation; Figure S4b). This is thought to be because the united atom model in the FF calculation exaggerates the repulsion between the CH_2 pseudoatom and the carbonyl oxygen in the *cis* configuration; there would be no attractive interaction between O and H of CH_2 unit in this united atom model. However, the problematic range showing a discrepancy between QM and FF ($\text{C-O-CH}_2\text{-CH}_2 < 40^\circ$) is accessible only with a high energy (at least 3 kcal/mol higher than the *trans* conformation; Figure S4b), and thus the contribution of this range is expected to be very small.

Table S2. Force field used for the acrylate backbone of the inhibitor.

$E^{\text{vdW}}(R) = D_0 \left\{ \left(\frac{R_0}{R} \right)^{12} - 2 \left(\frac{R_0}{R} \right)^6 \right\}, \quad E^{\text{bond}}(R) = \frac{1}{2} K_R (R - R_0)^2, \quad E^{\text{angle}}(\theta) = \frac{1}{2} K_\theta (\theta - \theta_0)^2, \quad E^{\text{torsion}}(\phi) = \sum_n \frac{1}{2} V_n [1 - d_n \cos(n\phi)]$						
E^{vdW}	$=\text{O} (sp^2)$	R_0^a	3.3225	D_0^b	0.210	
	$-\text{O}- (sp^3)$	R_0^a	3.3674	D_0^b	0.170	
	$\text{C} (sp^2)$	R_0^a	4.2092	D_0^b	0.105	
	$\text{C} (sp^3)$	R_0	3.8050	D_0	0.06921	
	H	R_0	2.9267	D_0	0.0335	
E^{bond}	$\text{C}(sp^2)=\text{O}$	R_0^a	1.200	K_b^c	1140	
	$\text{C}(sp^2)-\text{O}$	R_0^a	1.344	K_b^c	900	
	$\text{O-CH}_2/\text{CH}_3$	R_0^a	1.437	K_b^c	640	
	$\text{C}(sp^2)-\text{C}(sp^3)$	R_0^a	1.520	K_b^c	634	
	$\text{C}(sp^3)-\text{H}$	R_0	1.09	K_b	662	
	$\text{C}(sp^3)-\text{CH}_2/\text{CH}_3$	R_0	1.54	K_b	520	
	$\text{O}=\text{C}(sp^2)-\text{O}$	θ_0^d	125	K_θ^e	160	
	$\text{O}=\text{C}(sp^2)-\text{C}(sp^3)$	θ_0^d	125	K_θ^e	160	
E^{angle}	$\text{O}-\text{C}(sp^2)-\text{C}(sp^3)$	θ_0^d	110	K_θ^e	140	
	$\text{C}(sp^2)-\text{O-CH}_2/\text{CH}_3$	θ_0^d	115	K_θ^e	200	
	$\text{O-CH}_2\text{-CH}_2/\text{CH}_3$	θ_0^d	109.5	K_θ^e	160	
	$\text{H-C}(sp^3)-\text{C}(sp^2)$	θ_0^d	109.5	K_θ^e	70	
	$\text{H-C}(sp^3)-\text{CH}_2$	θ_0^d	109.5	K_θ^e	70	
	$\text{CH}_2\text{-C}(sp^3)-\text{CH}_2$	θ_0^d	114	K_θ^e	124.19	
	$\text{CH}_2\text{-C}(sp^3)-\text{C}(sp^2)$	θ_0^d	111.1	K_θ^e	126	
	$\text{C-CH}_2\text{-C-CH}_2$	$V_1^b (d_1)$	-0.50 (+1)	$V_2^b (d_2)$	1.70 (-1)	
		$V_3^b (d_3)$	-0.44 (+1)	$V_4^b (d_4)$	1.75 (-1)	
	$\text{O}=\text{C-C-CH}_2$	$V_1^b (d_1)$	0.42 (+1)	$V_2^b (d_2)$	0.38 (+1)	$V_3^b (d_3)$ 0.17 (+1)
E^{torsion}	O-C-C-CH_2	$V_1^b (d_1)$	0.42 (-1)	$V_2^b (d_2)$	0.38 (+1)	$V_3^b (d_3)$ 0.17 (+1)
	$\text{C-C-O-CH}_2/\text{CH}_3$	$V_1^b (d_1)$	0.54 (+1)	$V_2^b (d_2)$	6.18 (+1)	
		$V_3^b (d_3)$	0.38 (-1)	$V_4^b (d_4)$	0.17 (-1)	
	$\text{O}=\text{C-O-CH}_2/\text{CH}_3$	$V_1^b (d_1)$	0.54 (-1)	$V_2^b (d_2)$	6.18 (+1)	
		$V_3^b (d_3)$	0.38 (+1)	$V_4^b (d_4)$	0.17 (-1)	
	$\text{C-O-CH}_2\text{-CH}_2/\text{CH}_3$	$V_1^b (d_1)$	2.14 (+1)	$V_2^b (d_2)$	0.79 (+1)	$V_3^b (d_3)$ 1.29 (-1)

^aIn Å. This R_0 has been converted from the original literature value σ by the relation $R_0 = \sqrt[3]{2}\sigma$, since in the original literature E^{vdW} is defined as

$$E^{\text{vdW}}(R) = 4\epsilon \left\{ \left(\frac{\sigma}{R} \right)^{12} - \left(\frac{\sigma}{R} \right)^6 \right\}.$$

^bIn kcal/mol. This D_0 (kcal/mol) has been converted from the original literature value ϵ ($= D_0/k$ where k is the Boltzmann constant; K) by the relation by $\epsilon \times k$. ^cIn kcal/mol/Å². ^dIn degrees. ^eIn kcal/mol/rad².

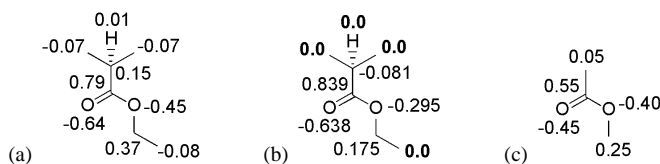


Figure S1. Atomic charges for the acrylate backbone of PAA1. (a) ESP-fitted charges from *ab initio* QM calculation, (b) ESP-fitted charges with constraints (denoted in bold face) on the methyl groups which will be replaced by neutral united atoms in simulations, and (c) atomic charge parameter used in OPLS FF. Charges in (b) were used in the simulations.

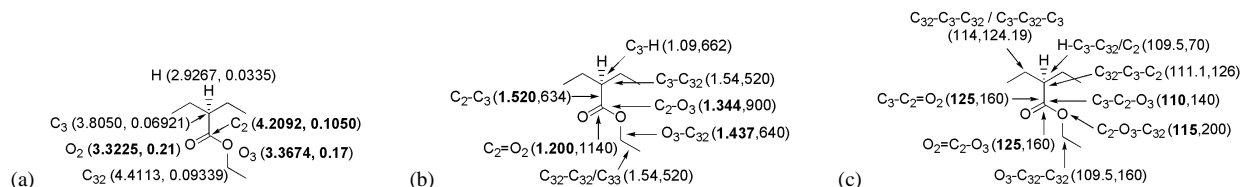


Figure S2. Force field for the acrylate backbone of the PAA1 inhibitor. Numbers in bold faces are taken from the OPLS FF. C_{32} and C_{33} represent united atoms for CH_2 and CH_3 groups, respectively. (a) vdW FF, R_0 (Å) and D_0 (kcal/mol); (b) bond stretching FF, R_0 (Å) and K_b (kcal/mol/Å²); (c) angle bending FF, θ_0 (Å) and K_θ (kcal/mol/rad²). Functionals forms are given in Table S1.

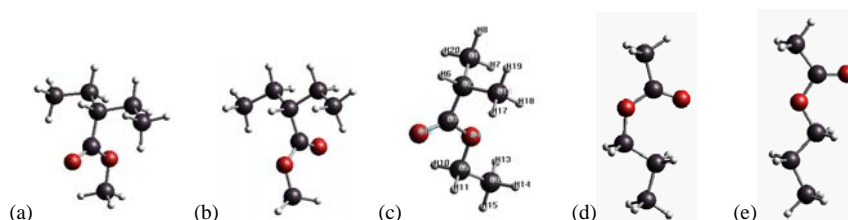


Figure S3. Model compounds used to obtain torsion potentials. (a-b) $(\text{CH}_3\text{CH}_2)_2\text{CHCOOCH}_3$ for torsion $\text{CH}_2\text{-CH-C=O}$ (a: 0° and b: 180°) and (c-e) $(\text{CH}_3)_2\text{CHCOOCH}_2\text{CH}_3$ (c) and $\text{CH}_3\text{COOCH}_2\text{CH}_2\text{CH}_3$ (d-e) for torsion $\text{C-O-CH}_2\text{-CH}_2$ (c: 0° and d-e: 180°). These two different models give essentially the same potential energy curves (dots in Figure S4b).

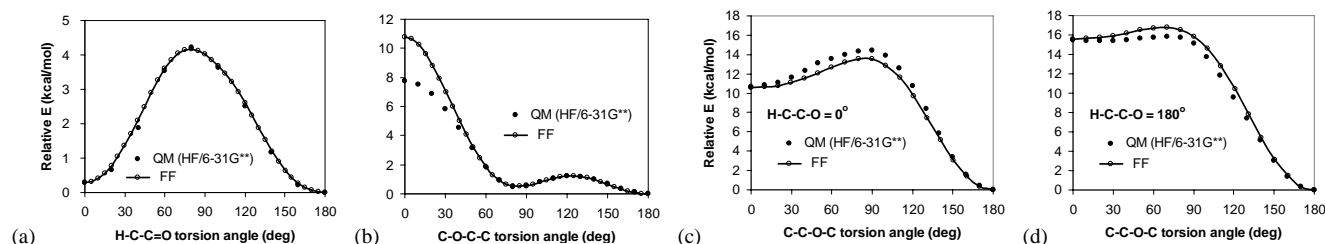


Figure 6. Torsion potential curves calculated on model compounds using QM (filled circle) and FF (solid line with open circles). (a) Torsion potential along $\text{CH}_2\text{-CH-C=O}$ of $(\text{CH}_3\text{CH}_2)_2\text{CHCOOCH}_3$ (Figures S3a-b), (b) torsion potential along $\text{C-O-CH}_2\text{-CH}_2$ of $(\text{CH}_3)_2\text{CHCOOCH}_2\text{CH}_3$ (Figure S3c) and $\text{CH}_3\text{COOCH}_2\text{CH}_2\text{CH}_3$ (Figures S3d-e), and (c-d) torsion potential along CH-C-O-CH_2 of $(\text{CH}_3)_2\text{CHCOOCH}_2\text{CH}_3$ (Figure S3c) when H-C-C=O is 0° (c) or 180° (d).

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