# A "two-step" chronoamperometric method for studying the anaerobic inactivation of an oxygen tolerant NiFe hydrogenase Supplementary material 

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## S1 Influence of the high-potential slope of $i(E)$ on the value of $E_{\text {switch }}$ in the limit of zero scan rate

We have derived the value of $E_{\mathrm{switch}}^{v \rightarrow 0}$ in the stationnary limit, by assuming that $i(E)$ is independent on potential in the potential range of potential around $E_{\text {switch }}^{v \rightarrow 0}$ (eq. 18). However, the data in figure 5 suggest that $E_{\mathrm{switch}}^{v \rightarrow 0}$ may fall in the region where $i(E)$ depends linearly on potential, as predicted by the model in reference 34. Let us now estimate how this shifts the position of the inflection point of $i(E)=i_{A}(E) \times A_{\infty}(E)$.

We are aiming at calculating the position of the inflection point of $i(E)$ defined by:

$$
\begin{equation*}
i(E)=\left[1+\epsilon\left(E-E_{\mathrm{switch}}^{v \rightarrow 0}\right)\right] A_{\infty}(E) \tag{S1}
\end{equation*}
$$

where $A_{\infty}(E)$ is a function that displays an inflection point at $E_{\mathrm{switch}}^{v \rightarrow 0}$. The second derivative of $i(E)$ is:

$$
\begin{equation*}
i^{\prime \prime}(E)=\left[1+\epsilon\left(E-E_{\mathrm{switch}}^{v \rightarrow 0}\right)\right] A_{\infty}^{\prime \prime}(E)+2 \epsilon A_{\infty}^{\prime}(E) \tag{S2}
\end{equation*}
$$

The inflection point of $i(E)$ is shifted with respect to that of $A_{\infty}(E)$ by a quantity $\delta E$, defined by:

$$
\begin{equation*}
i^{\prime \prime}\left(E_{\mathrm{switch}}^{v \rightarrow 0}+\delta E\right)=0 \tag{S3}
\end{equation*}
$$

This translates into:

$$
\begin{equation*}
[1+\epsilon \delta E] A_{\infty}^{\prime \prime}\left(E_{\mathrm{switch}}^{v \rightarrow 0}+\delta E\right)+2 \epsilon A_{\infty}^{\prime}\left(E_{\mathrm{switch}}^{v \rightarrow 0}+\delta E\right)=0 \tag{S4}
\end{equation*}
$$

We estimate the values of $A_{\infty}^{\prime \prime}\left(E_{\text {switch }}^{v \rightarrow 0}+\delta E\right)$ and $A_{\infty}^{\prime}\left(E_{\text {switch }}^{v \rightarrow 0}+\delta E\right)$ to the first order in $\delta E$ :

$$
\begin{align*}
& A_{\infty}^{\prime \prime}\left(E_{\mathrm{switch}}^{v \rightarrow 0}+\delta E\right)=A_{\infty}^{\prime \prime}\left(E_{\mathrm{switch}}^{v \rightarrow 0}\right)+\delta E A_{\infty}^{\prime \prime \prime}\left(E_{\mathrm{switch}}^{v \rightarrow 0}\right)=\delta E A_{\infty}^{\prime \prime \prime}\left(E_{\mathrm{switch}}^{v \rightarrow 0}\right)  \tag{S5}\\
& A_{\infty}^{\prime}\left(E_{\mathrm{switch}}^{v \rightarrow 0}+\delta E\right)=A_{\infty}^{\prime}\left(E_{\mathrm{switch}}^{v \rightarrow 0}\right)+\delta E A_{\infty}^{\prime \prime}\left(E_{\mathrm{switch}}^{v \rightarrow 0}\right)=A_{\infty}^{\prime}\left(E_{\mathrm{switch}}^{v \rightarrow 0}\right) \tag{S6}
\end{align*}
$$

Keeping only the first-order terms in $\delta E$ gives:

$$
\begin{equation*}
\delta E=-\frac{2 \epsilon A_{\infty}^{\prime}\left(E_{\mathrm{switch}}^{v \rightarrow 0}\right)}{A_{\infty}^{\prime \prime \prime}\left(E_{\mathrm{switch}}^{v \rightarrow 0}\right)} \tag{S7}
\end{equation*}
$$

Now, we equate $A_{\infty}(E)$ and $A_{\infty}(E)$, hence

$$
\begin{equation*}
\frac{A_{\infty}^{\prime}\left(E_{\text {switch }}^{v \rightarrow 0}\right)}{A_{\infty}^{\prime \prime \prime}\left(E_{\text {switch }}^{v \rightarrow 0}\right)}=-\frac{2}{(\alpha f)^{2}}=2.10^{-3} \mathrm{~V}^{2} \tag{S8}
\end{equation*}
$$

with $\alpha=0.66$. We identify $i(E)$ and $1+\epsilon\left(E-E_{\text {switch }}^{v \rightarrow 0}\right)$, and a linear fit of the green trace of figure 5 returns a value of $\epsilon$ of about $2 \mathrm{~V}^{-1}$. This gives:

$$
\begin{equation*}
\delta E=+13 \mathrm{mV} \tag{S9}
\end{equation*}
$$

Under steady-state conditions, $E_{\text {switch }}$ is therefore 10 mV to 15 mV greater than the position of the inflection point of $A_{\infty}(E)$, given by eq. 18 in the main text.

## S2 The EC model

In the framework of scheme 1, the equations giving the evolution of the fractions of enzyme in the NiB , I and A forms read:

$$
\begin{align*}
\frac{\mathrm{d}[\mathrm{NiB}]}{\mathrm{d} t} & =-k_{r}[\mathrm{NiB}]+k_{o}[\mathrm{I}]  \tag{S10}\\
\frac{\mathrm{d}[\mathrm{~A}]}{\mathrm{d} t} & =-k_{-1}[\mathrm{~A}]+k_{1}[\mathrm{I}]  \tag{S11}\\
1 & =[\mathrm{A}]+[\mathrm{NiB}]+[\mathrm{I}] \tag{S12}
\end{align*}
$$

We are interested in the evolution of the active fraction $A(t)$. The quantity $A^{\prime}=[\mathrm{A}]-[\mathrm{A}]_{\infty}$, where $[\mathrm{A}]_{\infty}$ is the asymtotic value for $[\mathrm{A}]_{\infty}$, is governed by:

$$
\begin{equation*}
\frac{\mathrm{d}^{2} A^{\prime}}{\mathrm{d} t^{2}}+\left(k_{r}+k_{o}+k_{1}+k_{-1}\right) \frac{\mathrm{d} A^{\prime}}{\mathrm{d} t}+\left(k_{r} k_{1}+k_{r} k_{-1}+k_{o} k_{-1}\right) A^{\prime}=0 \tag{S13}
\end{equation*}
$$

The general solution is:

$$
\begin{equation*}
A^{\prime}(t)=A_{+} \exp \left(-k_{+} t\right)+A_{-} \exp \left(-k_{-} t\right) \tag{S14}
\end{equation*}
$$

with

$$
\begin{equation*}
k_{ \pm}=\frac{k_{r}+k_{o}+k_{1}+k_{-1} \pm \sqrt{\left(k_{r}+k_{o}+k_{1}+k_{-1}\right)^{2}-4\left(k_{r} k_{1}+k_{r} k_{-1}+k_{o} k_{-1}\right)}}{2} \tag{S15}
\end{equation*}
$$

on condition that $k_{+} \neq k_{-} . A_{+}$and $A_{-}$are determined by the initial conditions. In general, the reaction scheme 1, leads to a bi-exponential relaxation. However, a mono-exponential relaxation can be observed in particular cases: (1) after a rapid equilibration between NiB and I (2) when either $\mathrm{I}, \mathrm{NiB}$ or A is not accumulated (3) after a rapid equilibration between $I$ and $A$. We rule out the latter because it would lead to a potential dependent $k_{I}$. In the following, we consider the two other cases.

Case 1 A rapid equilibrium between NiB and I occurs when $k_{r}+k_{o} \gg k_{1}+k_{-1}$. Eq. S15 gives:

$$
\begin{equation*}
k_{+} \approx k_{o}+k_{r} \quad k_{-} \approx \frac{k_{1} k_{r}}{k_{o}+k_{r}}+k_{-1} \tag{S16}
\end{equation*}
$$

The equilibration between NiB and I , which is determined by the high rate $k_{+}$, is too fast to be observed experimentally. The relaxation is governed by the slow rate $k_{-}$, which matches the value of $\left(k_{A}+k_{I}\right)$ given by eq. 21.

Case 2. The only possibility is that I is not accumulated. This occurs when $k_{o} \gg k_{r}+k_{-1}$ and
$k_{1} \gg k_{r}+k_{-1}$, and leads to:

$$
\begin{equation*}
k_{+} \approx k_{o}+k_{1} \quad k_{-} \approx \frac{k_{r} k_{1}+k_{o} k_{-1}}{k_{o}+k_{1}} \tag{S17}
\end{equation*}
$$

The depletion of I, which is determined by the high rate $k_{+}$, is not observed. The slow rate $k_{-}$determines the relaxation monitored during the experiment, and matches the value of $\left(k_{A}+k_{I}\right)$ given by eq. 21 .

## S3 The simplification of equations 22 and 23

Let us now derive the conditions of validity of the simplified equations 24 and 25 . Equation 24 only holds when $k_{3} \gg k_{o}$, i.e. for

$$
\begin{equation*}
E<E_{\max A}=E^{0}+\frac{1}{f(1-\alpha)} \ln \frac{k_{3}}{k_{0}} \tag{S18}
\end{equation*}
$$

The fact that the experimental value of $k_{A}$ does not deviate from that predicted by eq 24, (fig. 7) implies that $E_{\max A}$ is larger than the potential at which the term $k_{\text {Alim }}$ starts to dominate $k_{A}$ (about +0.2 V for pH 8 ). Due to the relatively low value of $1-\alpha=0.4$, even relatively small values of the ratio $k_{3} / k_{0}$ lead to relatively large values of $E_{\max A}$ with respect to $E^{0}$; for instance, for $k_{3} / k_{0}=10, E_{\max A}$ is 150 mV greater than $E^{0}$.

Similarly, the condition for the validity of equation 25 is:

$$
\begin{equation*}
E>E_{\lim I} \tag{S19}
\end{equation*}
$$

where $E_{\text {lim } I}$ is the maximum of

$$
\begin{equation*}
E^{0^{\prime}}-\frac{1}{(1-\alpha) f} \ln \frac{k_{0}^{\prime}}{k_{1}} \quad \text { and } \quad E^{0^{\prime}}-\frac{1}{\alpha f} \ln \frac{k_{2}}{k_{0}^{\prime}} \tag{S20}
\end{equation*}
$$

where the prime parameters characterize the redox couple $\mathrm{I}_{2} / \mathrm{I}_{1}$, by analogy with eqs 19 . $E_{\text {lim }}$ only needs to be superior to the lowest potential where the determination of $k_{I}$ is unambiguous (ie, $k_{I}>0.1 \times k_{A}$ ).

## S4 Supplementary figures



Figure S1: Illustration of the effect of subtracting the capacitive current recorded in a control experiment with no adsorbed enzyme.

Figure S2: Cyclic voltammograms of adsorbed



A. aeolicus hydrogenase at different scan rates (solid lines), and the corresponding simulations (dotted lines). Conditions: $\mathrm{pH} 6.1,40^{\circ} \mathrm{C}, 1 \mathrm{~atm}$. of $\mathrm{H}_{2}, \omega=4 \mathrm{krpm}$. Parameters: $k_{I}=3.5 \times$ $10^{-3} \mathrm{~s}^{-1}, \alpha=0.80, k_{0} \exp \alpha f E^{0}=0.28 \mathrm{~s}^{-1}$ and $k_{\text {Alim }}=5 \times 10^{-4} \mathrm{~s}^{-1}$, and the parameters describing the wave shape: $E_{O H}=-0.212 \mathrm{~V}$, $E_{H R}=-0.340 \mathrm{~V}, k_{2} / k_{0}^{\max }=4 \times 10^{-2}$ and $i_{\lim } / \beta d_{0} \approx 0.2 \mu \mathrm{~A}$ (the exact value was adjusted to account for film loss). Film loss during the course of each voltammogram was neglected. The small differences with the parameters used in figure 5 in the main text results from the pH being slightly greater and, possibly, uncertainties in the concentration of dissolved $\mathrm{H}_{2}$.

Figure S3: Influence on the value of the vertex (the potential where the direction of the scan is returned) on the value of $E_{\text {switch }}$. Panel A shows the voltammograms recorded with $A$. aeolicus NiFe hydrogenase adsorbed at a rotating electrode, at $\mathrm{pH} 6,40^{\circ} \mathrm{C}$, under 1 atm. $\mathrm{H}_{2}$. Electrode rotation rate 3 krpm , scan rate $1.5 \mathrm{mV} . \mathrm{s}^{-1}$. Panel B shows the $1^{\text {st }}$ derivative of the sweeps towards negative potentials, from which the values of $E_{\text {switch }}$ are determined. Panel C shows the values of $E_{\text {switch }}$ against vertex potential. The color code is the same in all panels.




Figure S4: Influence of scan rate on the value of $E_{\text {switch }}$ in experiments carried out with $D$. fructosovorans NiFe hydrogenase. Panel A shows the voltammograms recorded at $\mathrm{pH} 5.5,40^{\circ} \mathrm{C}$, under $0.1 \mathrm{~atm} . \mathrm{H}_{2}$. Electrode rotation rate 2.2 krpm , scan rate as indicated. Panel B shows the values of $E_{\text {switch }}$ against scan rate. The color code is the same in panel A and B.

Figure S5: Effect of the variation of $\mathrm{H}_{2}$ pressure on activation and inactivation rate constants of A. aeolicus NiFe hydrogenase. Open circles show the values of $k_{I}$, filled squares are the corresponding values of $k_{A}$. The red and blue symbol correspond to $\mathrm{H}_{2}$ pressure of 0.1 and 1 atm , respectively. Other conditions: $40^{\circ} \mathrm{C}$, $\mathrm{pH} 6, \omega=4 \mathrm{krpm}$.

## S5 Supplementary tables

|  | A. vinosum <br> pH 8.8 | A. aeolicus <br> pH 8 |
| :--- | :---: | :---: |
| $E_{0}(\mathrm{mV})$ | -284 | 0 |
| $k_{1}\left(\mathrm{~s}^{-1}\right)$ | 1 | $4.10^{-2}$ |
| $k_{\text {Alim }}\left(\mathrm{s}^{-1}\right)$ | $7.10^{-3 \mathrm{a}}$ | $3.10^{-4}$ |
| $k_{-1}\left(\mathrm{~s}^{-1}\right)$ |  | $7.10^{-3}$ |

[^0]Table S1: Values used for the fits of the data for model 1.

|  | A. vinosum <br> pH 8.8 | A. aeolicus <br> pH 8 |
| :--- | :---: | :---: |
| $\alpha$ | 0.8 | 0.64 |
| $k_{0} \exp \left(-f \alpha E^{0}\right)\left(\mathrm{s}^{-1}\right)$ | $9.10^{-5}$ | $1.10^{-2}$ |
| $k_{\text {Alim }}\left(\mathrm{s}^{-1}\right)$ | $7.10^{-3 \mathrm{a}}$ | $3.10^{-4}$ |
| $k_{-1}\left(\mathrm{~s}^{-1}\right)$ |  | $7.10^{-3}$ |

${ }^{a}$ As for other fits, only the sum of $k_{-1}$ and $k_{\text {Alim }}$ can be determined.

Table S3: Values used for the fits of the data for model 3.

|  | A. vinosum <br> pH 8.8 | A. aeolicus <br> pH 8 |
| :--- | :---: | :---: |
| $E_{0}(\mathrm{mV})$ | -290 | -160 |
| $k_{1}\left(\mathrm{~s}^{-1}\right)$ | 1.5 | 8 |
| $k_{\text {Alim }}\left(\mathrm{s}^{-1}\right)$ |  | $5.10^{-3 \mathrm{a}}$ |
| $k_{-1}\left(\mathrm{~s}^{-1}\right)$ |  | $5.10^{-4}$ |
| $\alpha$ | 0.55 | $8.10^{-3}$ |
| $k_{0}\left(\mathrm{~s}^{-1}\right)$ | 0.9 | 2.30 |

${ }^{a}$ As for other fits, only the sum of $k_{-1}$ and $k_{\text {Alim }}$ can be determined.

Table S2: Values used for the fits of the data for model 2.


[^0]:    ${ }^{a}$ For A. vinosum, as only $k_{I}+k_{A}$ is known $k_{\text {Alim }}$ and $k_{-1}$ cannot be independently determined, only their sum is known.

