

**Probing Side Chain Dynamics in Proteins by the Measurement of Nine Deuterium Relaxation
Rates Per Methyl Group**

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Simulations of $R_2(\text{CHD}_2)$ Decay in $(^{13}\text{CH})\text{D}_2$ Methyl Groups. Relaxation of ^2H single-quantum transitions in $(^{13}\text{CH})\text{D}_2$ methyl groups ρ_{i-j} (Figure 2 of the main text) can be conveniently represented in the following form,

$$\frac{d}{dT} \begin{bmatrix} \rho_{1-2} \\ \rho_{2-3} \\ \rho_{3-4} \\ \rho_{4-5} \\ \rho_{5-6} \\ \rho_{6-7} \end{bmatrix} = -\left(\frac{1}{80}\right)(2\pi QCC)^2 \left[\begin{aligned} & \begin{pmatrix} 9 & 0 & 0 & 0 & 0 & 0 \\ 0 & 9 & 0 & 0 & 0 & 0 \\ 0 & 0 & 9 & 0 & 0 & 0 \\ 0 & 0 & 0 & 9 & 0 & 0 \\ 0 & 0 & 0 & 0 & 9 & 0 \\ 0 & 0 & 0 & 0 & 0 & 9 \end{pmatrix} J^A(0) + \begin{pmatrix} 15 & -\sqrt{6} & 0 & 0 & 3\sqrt{2} & 0 \\ -\sqrt{6} & 19 & 1 & 0 & 0 & -3\sqrt{3} \\ 0 & 1 & 19 & -\sqrt{6} & 3\sqrt{3} & 0 \\ 0 & 0 & -\sqrt{6} & 15 & 0 & 3\sqrt{2} \\ 3\sqrt{2} & 0 & 3\sqrt{3} & 0 & 15 & 3 \\ 0 & -3\sqrt{3} & 0 & 3\sqrt{2} & 3 & 15 \end{pmatrix} J^A(\omega_D) \\ & + \begin{pmatrix} 6 & -\sqrt{6} & 0 & 0 & -3\sqrt{2} & 0 \\ -\sqrt{6} & -2 & 1 & 0 & 0 & 3\sqrt{3} \\ 0 & 1 & -2 & -\sqrt{6} & 3\sqrt{3} & 0 \\ 0 & 0 & -\sqrt{6} & 6 & 0 & -3\sqrt{2} \\ -3\sqrt{2} & 0 & 3\sqrt{3} & 0 & -6 & 3 \\ 0 & 3\sqrt{3} & 0 & -3\sqrt{2} & 3 & -6 \end{pmatrix} J^C(\omega_D) \\ & + \begin{pmatrix} 18 & 0 & -2\sqrt{6} & 0 & 0 & -6\sqrt{2} \\ 0 & 10 & 0 & -2\sqrt{6} & 0 & 0 \\ -2\sqrt{6} & 0 & 10 & 0 & 0 & 0 \\ 0 & -2\sqrt{6} & 0 & 18 & -6\sqrt{2} & 0 \\ 0 & 0 & 0 & -6\sqrt{2} & 18 & 0 \\ -6\sqrt{2} & 0 & 0 & 0 & 0 & 18 \end{pmatrix} J^A(2\omega_D) \\ & + \begin{pmatrix} 0 & 0 & -2\sqrt{6} & 0 & 0 & 6\sqrt{2} \\ 0 & 4 & 0 & -2\sqrt{6} & 0 & 0 \\ -2\sqrt{6} & 0 & 4 & 0 & 0 & 0 \\ 0 & -2\sqrt{6} & 0 & 0 & 6\sqrt{2} & 0 \\ 0 & 0 & 0 & 6\sqrt{2} & -12 & 0 \\ 6\sqrt{2} & 0 & 0 & 0 & 0 & -12 \end{pmatrix} J^C(2\omega_D) \end{aligned} \right] \begin{bmatrix} \rho_{1-2} \\ \rho_{2-3} \\ \rho_{3-4} \\ \rho_{4-5} \\ \rho_{5-6} \\ \rho_{6-7} \end{bmatrix}$$

where $QCC=(e^2qQ/h)$ is a quadrupolar coupling constant (167 kHz¹) and $J(\omega)$ is the spectral density function,²⁻⁴

$$J(\omega) = \frac{1}{9} S_{axis}^2 \left(\frac{\tau_c}{1 + (\omega\tau_c)^2} \right) + \left(P_2(\cos\alpha) - \frac{1}{9} S_{axis}^2 \right) \frac{\tau_e}{1 + (\omega\tau_e)^2} \quad (\text{S1})$$

where S_{axis}^2 is the generalized order parameter of the methyl 3-fold symmetry axis; $P_2 = (1/2)[3\cos^2(x)-1]$; $\alpha=0^\circ$ and $P_2(\cos\alpha) = 1$ for the auto-correlation spectral density function, $J^A(\omega)$, while $\alpha = 109.5^\circ$ and $P_2(\cos\alpha) = -1/3$

for the quadrupolar cross-correlation spectral density, $J^C(\omega)$; τ_c is the (isotropic) correlation time of overall molecular re-orientation; $1/\tau_e = 1/\tau_f + 1/\tau_c$ where τ_f is the correlation time of fast local motions.

Simulations of the transverse magnetization decay in (^{13}CH)D₂ groups, $R_2(\text{CHD}_2)$, have been performed using the matrix exponential of the matrices in eqs S1 expanded to include the triple-quantum (3Q) terms present during relaxation delay T , $(\rho_{1-4}^{3Q} + \rho_{2-5}^{3Q})$ (see expression 5 of the main text). The following initial conditions for $\rho_{i,j}$ have been used:

$$\rho(0) = (1/9) \left[\sqrt{F_1} \rho_{1-2}; \sqrt{F_2} \rho_{2-3}; \sqrt{F_2} \rho_{3-4}; \sqrt{F_1} \rho_{4-5}; \sqrt{F_3} \rho_{5-6}; \sqrt{F_3} \rho_{6-7}; \sqrt{F_{3Q}} (\rho_{1-4}^{3Q} + \rho_{2-5}^{3Q}) \right]^T, \text{ where}$$

F_i is the ‘out-and-back’ transfer function of each individual transition in the experiment of Figure 1 of the main text, and superscript T denotes the transposition operation. From expression (5) of the main text it follows that:

$$F_1 = 0.25 \left[2 \sin(4\pi J_{CD} T_C) + \sin(8\pi J_{CD} T_C) \right]^2 \quad (\text{S2a})$$

$$F_2 = 1.5 \sin^2(8\pi J_{CD} T_C) \quad (\text{S2b})$$

$$F_3 = 2 \sin^2(4\pi J_{CD} T_C) \quad (\text{S2c})$$

$$F_{3Q} = \left[0.5 \sin(8\pi J_{CD} T_C) - \sin(4\pi J_{CD} T_C) \right]^2 \quad (\text{S2d})$$

The sum of all transfer functions in eqs S2 above gives $4 \sin^2(4\pi J_{CD} T_C) + 2 \sin^2(8\pi J_{CD} T_C)$ - the transfer function used for the sum of *all* transitions relaxing during the delay T in the experiment for $R_2(\text{CHD}_2)$ measurement (Figure 1, inset A). To account for the differential transfer efficiency of $\rho_{i,j}$ in the reverse pathway of the scheme, after the relaxation period T each of the transitions has been pre-multiplied by its corresponding $\sqrt{F_i}$ before summation. Note the scaling of $\rho(0)$ as above ensures that the sum of all transitions at $T = 0$ is normalized to unity.

Thus simulated $R_2(\text{CHD}_2)$ decays have been fitted to a single-exponential function $I = I_0 \exp(-RT)$, where I is the simulated intensity and R is the relaxation rate. The results of these simulations expressed as relative errors in $R_2(\text{CHD}_2)$ arising from non-single-exponential decay at higher frequencies, $100 * R_2(\text{CHD}_2) - R(D_+, \text{CH}_2\text{D}) / R(D_+, \text{CH}_2\text{D})$, are illustrated in Figure 4 of the main text for several global molecular correlation times.

Figure S1

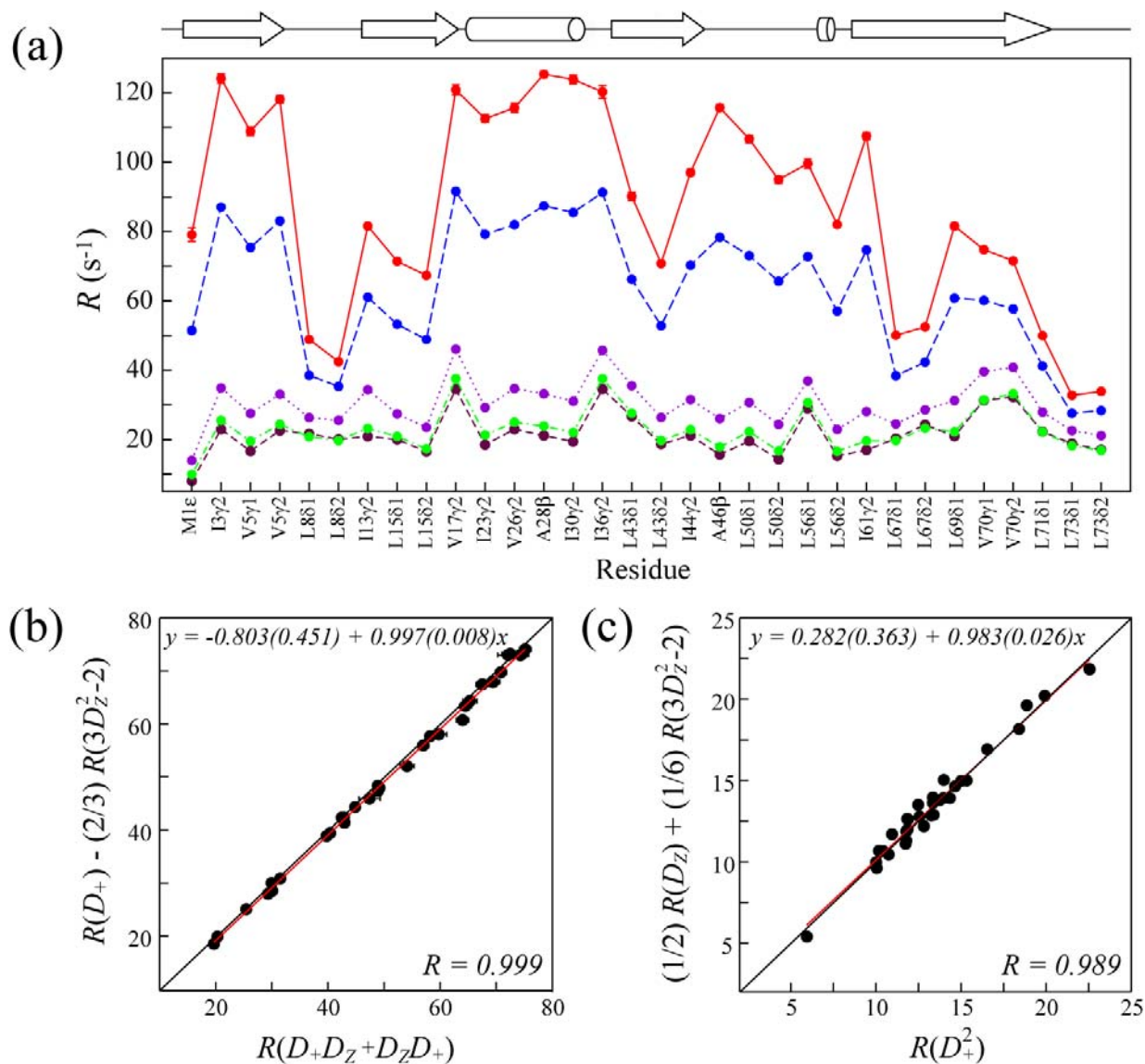


Figure S1. (a) All ²H relaxation rates measured in ¹³CH₂D methyl groups of ubiquitin at 27 °C satisfy the consistency relations. (a) $(5/3)R(D_+D_z + D_zD_+) \geq R(D_+) \geq (5/3)R(3D_z^2-2) \geq (5/3)R(D_+^2) \geq R(D_z)$. The rates $(5/3)R(D_+D_z + D_zD_+)$, $(5/3)R(D_+)$, $R(3D_z^2-2)$, $(5/3)R(D_+^2)$ and $R(D_z)$ are shown with red, blue, purple, green and brown circles/lines, respectively, vs. methyl positions; (b-c) Linear correlation plots of (b) $R(D_+D_z + D_zD_+)$ rates (x-axis) vs. $R(D_+) - (2/3)R(3D_z^2-2)$ (y-axis), and (c) $R(D_+^2)$ (x-axis) vs. $(1/2)R(D_z) + (1/6)R(3D_z^2-2)$ (y-axis) for 32 methyl groups in ubiquitin.

Figure S2

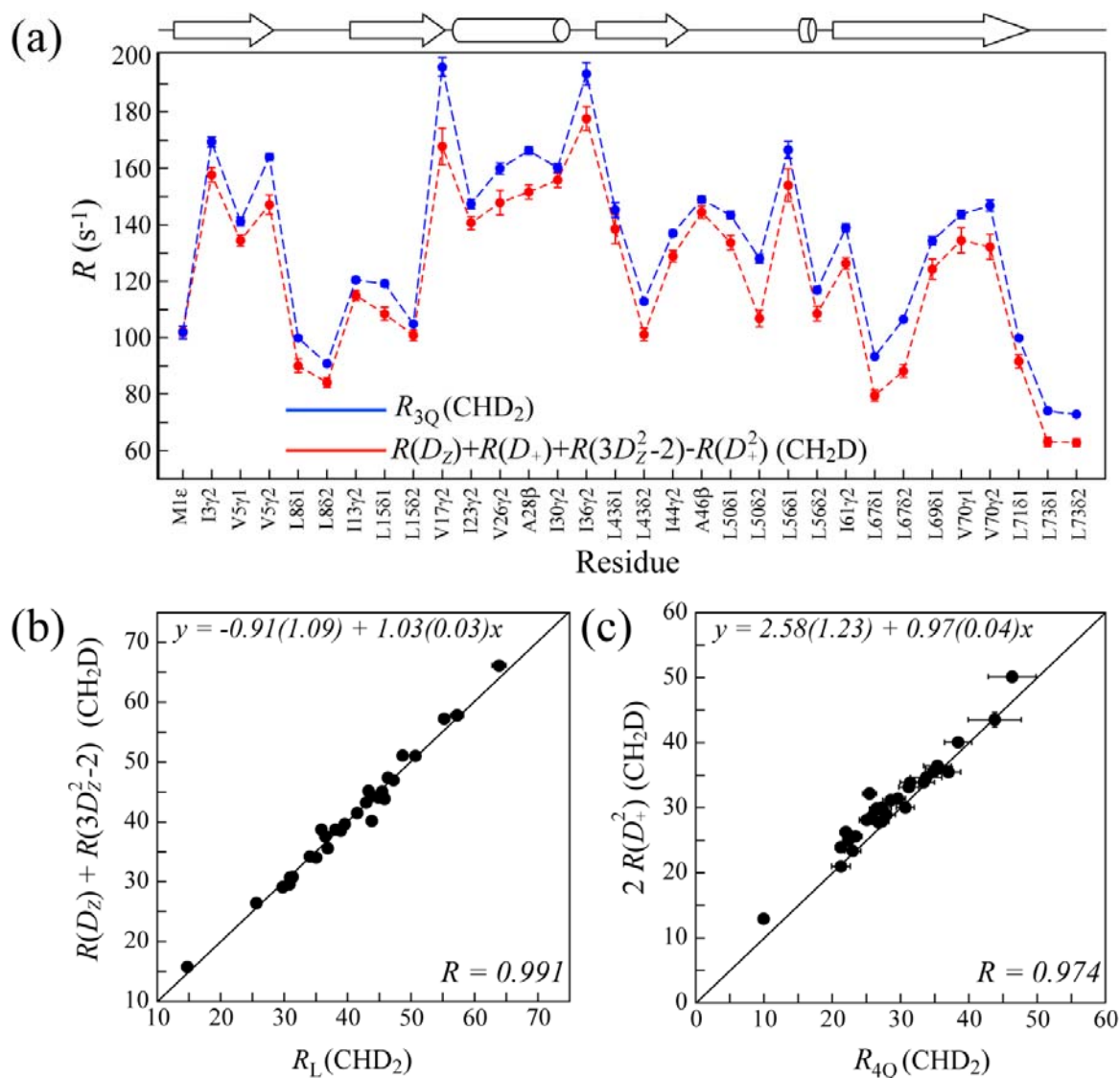


Figure S2. (a) All ²H relaxation rates measured in ¹³CHD₂ methyl groups of ubiquitin at 10 °C satisfy the consistency relations. (a) $R_{3Q}(\text{CHD}_2) \leq R(D_Z) + R(D_+) + R(3D_Z^2 - 2) - R(D_+^2)$ (CH₂D) (shown with red and blue circles/lines, respectively, vs. methyl position). (b-c) Linear correlation plots of (b) $R_L(\text{CHD}_2)$ rates (x-axis) vs. $R(D_Z) + R(3D_Z^2 - 2)$ (CH₂D) (y-axis) and (c) $R_{4Q}(\text{CHD}_2)$ (x-axis) vs. $2R(D_+^2)$ (CH₂D) (y-axis) for 29 methyl groups of ubiquitin. Three residues with large errors in the measurements (> 6 %) have been eliminated from the plots in (b-c).

Figure S3

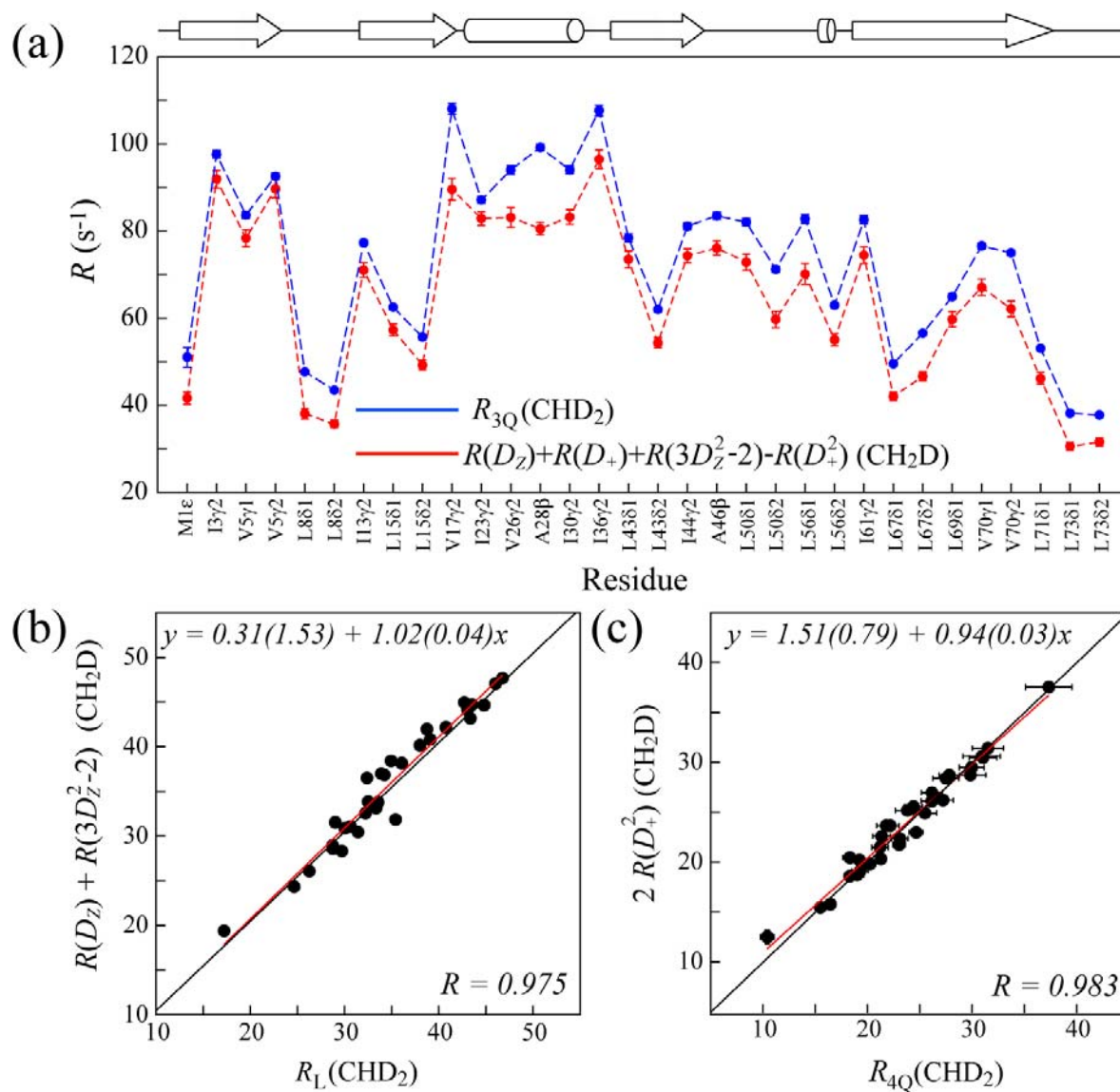


Figure S3. (a) All ²H relaxation rates measured ¹³CHD₂ methyl groups of ubiquitin at 40°C satisfy the consistency relations. (a) $R_{3Q}(\text{CHD}_2) < R(D_Z) + R(D_+) + R(3D_Z^2 - 2) - R(D_+^2)$ (CH₂D) (shown with red and blue circles/lines, respectively vs. methyl position). (b-c) Linear correlation plots of (b) $R_L(\text{CHD}_2)$ rates (x-axis) vs. $R(D_Z) + R(3D_Z^2 - 2)$ (CH₂D) (y-axis) and (c) $R_{4Q}(\text{CHD}_2)$ (x-axis) vs. $2R(D_+^2)$ (CH₂D) (y-axis) for 31 methyl groups in ubiquitin. 1 residue has been eliminated from the plots in (b-c) due to very high errors in the extracted relaxation rates.

Figure S4

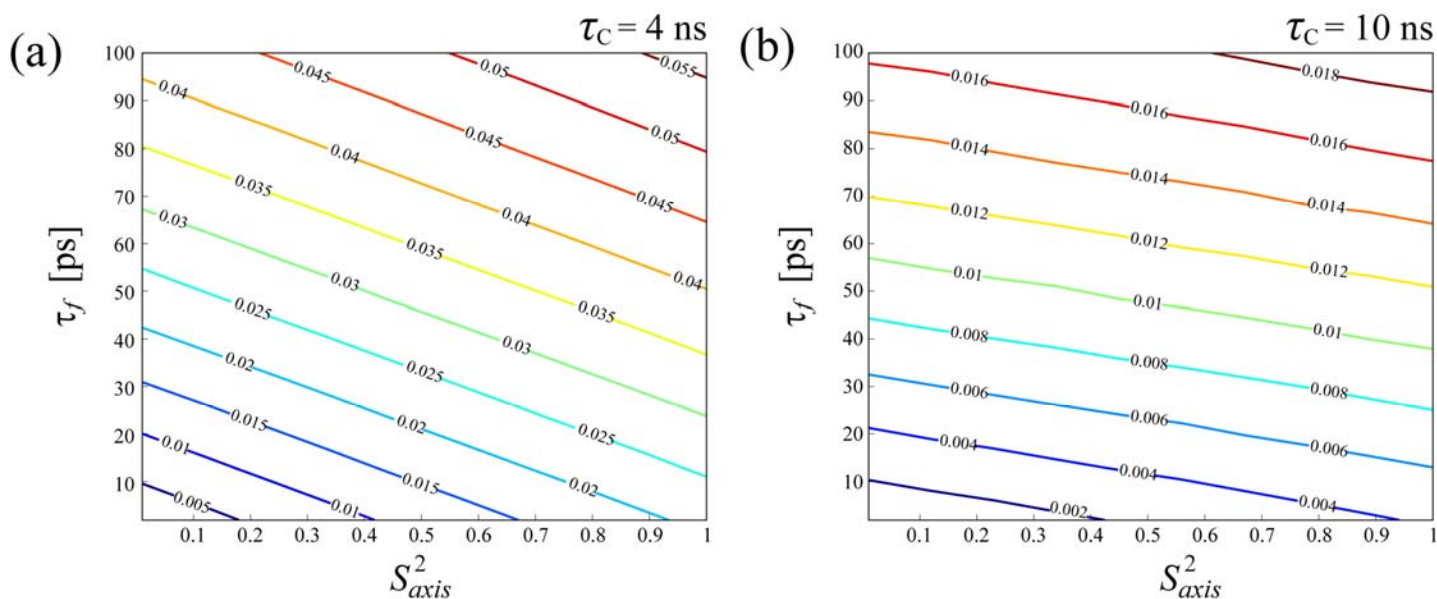


Figure S4. Contour plots showing simulated absolute differences between synthetic- $[R_L; R_2(\text{CHD}_2)]$ -derived and input S^2_{axis} values ($S^2_{axis}(\text{CHD}_2) - S^2_{axis}$) as a function of $(S^2_{axis}; \tau_f)$ for the global isotropic molecular reorientation correlation times (τ_C) of (a) 4 ns; (b) 10 ns. The $R_2(\text{CHD}_2)$ rates have been simulated as described in the Supporting Information and best-fit to $R(D_+; \text{CH}_2\text{D}) = (1/80)(2\pi QCC)^2[9J(0) + 15J(\omega_D) + 6J(2\omega_D)]$, while $R_L(\text{CHD}_2)$ have been fit to $(3/40)(2\pi QCC)^2[4J(\omega_D) + 4J(2\omega_D)] = R(3D_Z^2 - 2; \text{CH}_2\text{D}) + R(D_Z; \text{CH}_2\text{D})$.

Figure S5

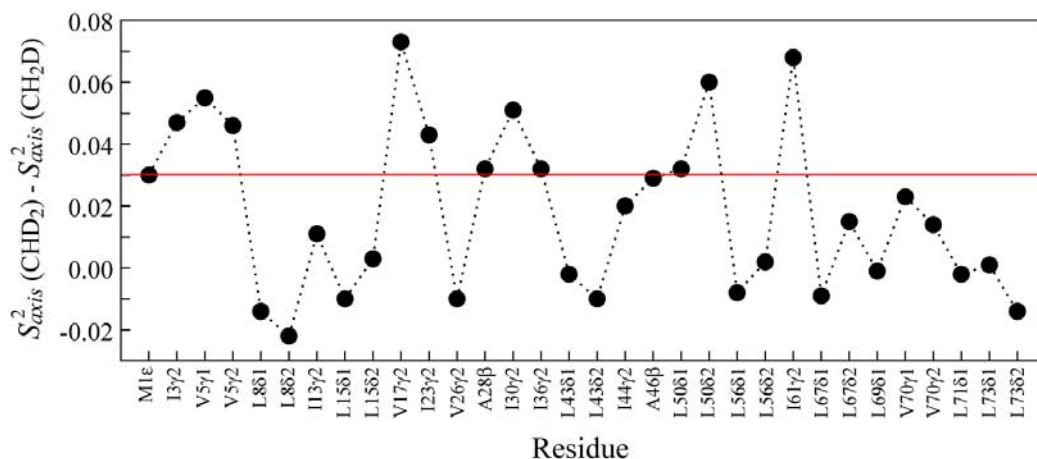


Figure S5. The differences between the $[R_2, R_L(\text{CHD}_2)]$ -derived and $[D_+, D_Z(\text{CH}_2\text{D})]$ -derived S^2_{axis} values in ubiquitin at 40°C $[S^2_{axis}(\text{CHD}_2) - S^2_{axis}(\text{CH}_2\text{D})]$ plotted versus methyl-containing residue numbers. The approximate position of the average absolute deviation $[S^2_{axis}(\text{CHD}_2) - S^2_{axis}(\text{CH}_2\text{D})]$ predicted from simulations for $\tau_C = 4$ ns is shown with horizontal red line.

Figure S6

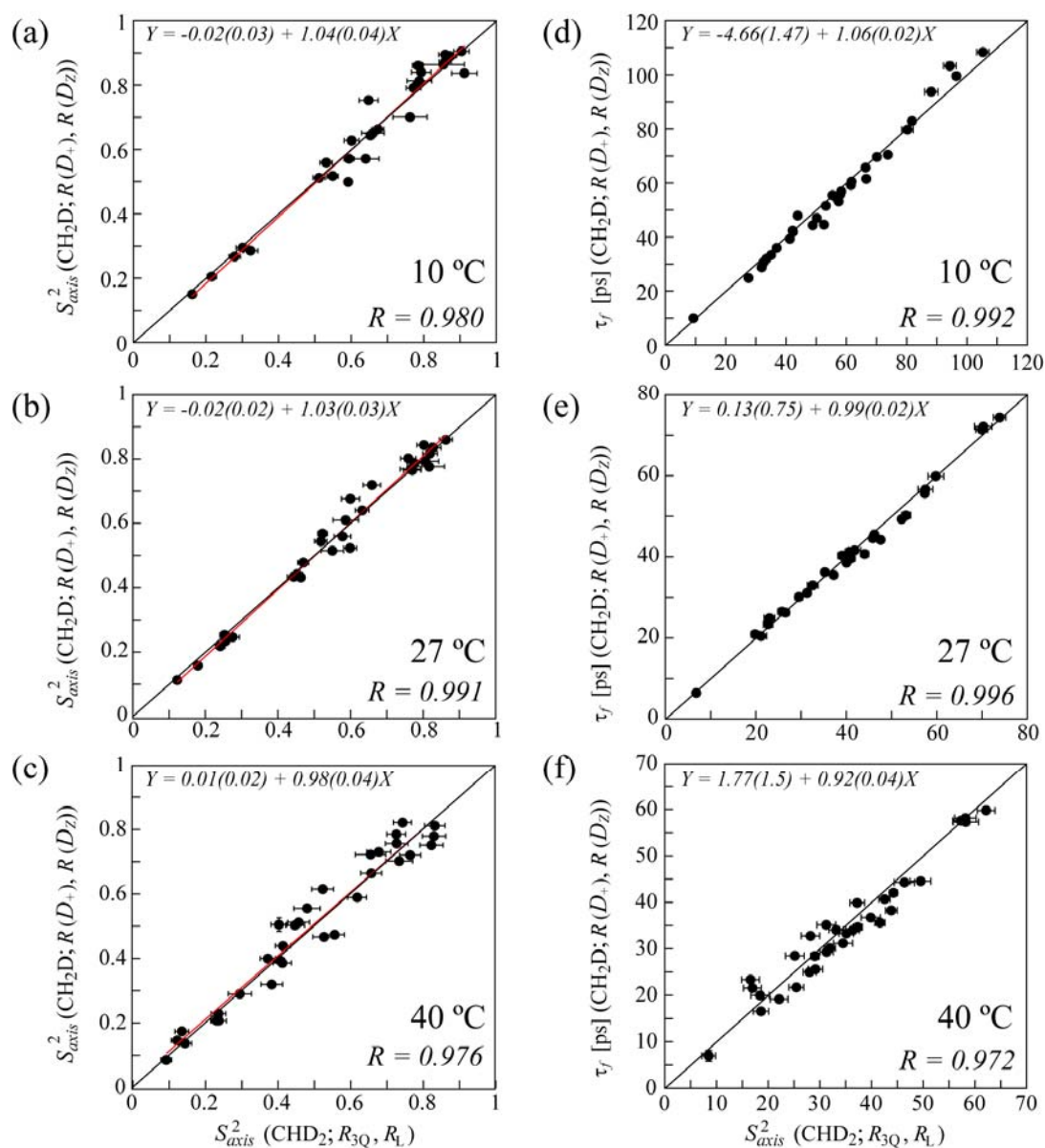


Figure S6. Linear correlation plots of **(a-c)** $[R_{3Q}, R_L(\text{CHD}_2)]$ -derived S^2_{axis} (x-axis) vs. $[R(D_+, D_Z; \text{CH}_2\text{D})]$ -derived S^2_{axis} (y-axis) at 10, 27 and 40 °C; **(d-f)** $[R_{3Q}, R_L(\text{CHD}_2)]$ -derived τ_f values (x-axis) vs. $[R(D_+, D_Z; \text{CH}_2\text{D})]$ -derived τ_f (y-axis) at 10, 27 and 40 °C. Linear regression parameters are shown for each plot along with Pearson correlation coefficients, R . Solid diagonal lines correspond to $y = x$.

Figure S7

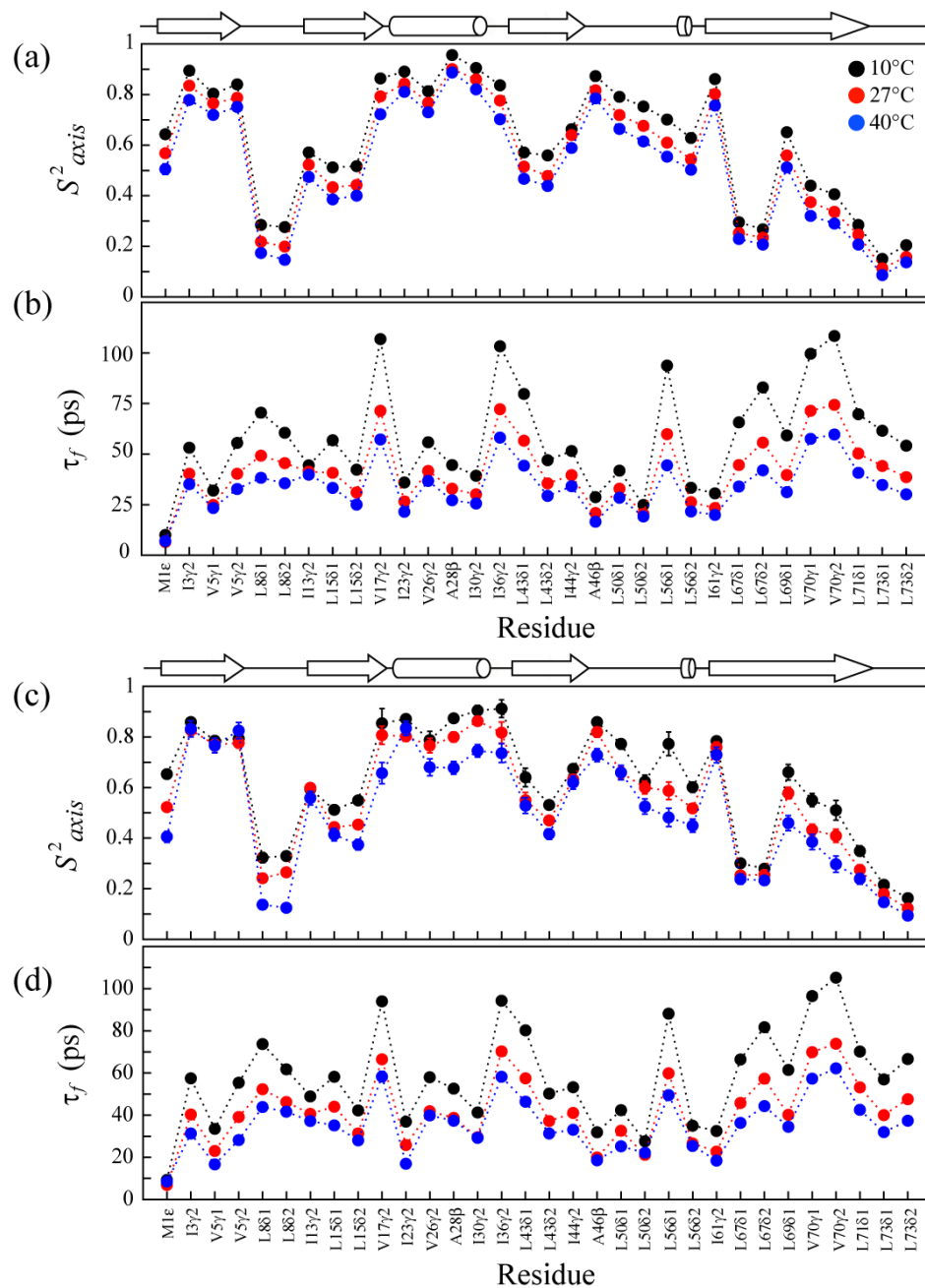


Figure S7. Comparisons of S^2_{axis} and τ_f values obtained for 32 methyl groups of ubiquitin at 10°C (black circles), 27°C (red circles) and 40°C (blue circles): (a) $[R(D_+, D_Z; CH_2D)]$ -derived S^2_{axis} , (b) $[R(D_+, D_Z; CH_2D)]$ -derived τ_f values (ps); (c) $[(R_{3Q}, R_L; CHD_2)]$ -derived S^2_{axis} and (d) $[(R_{3Q}, R_L; CHD_2)]$ -derived τ_f values.

Figure S8

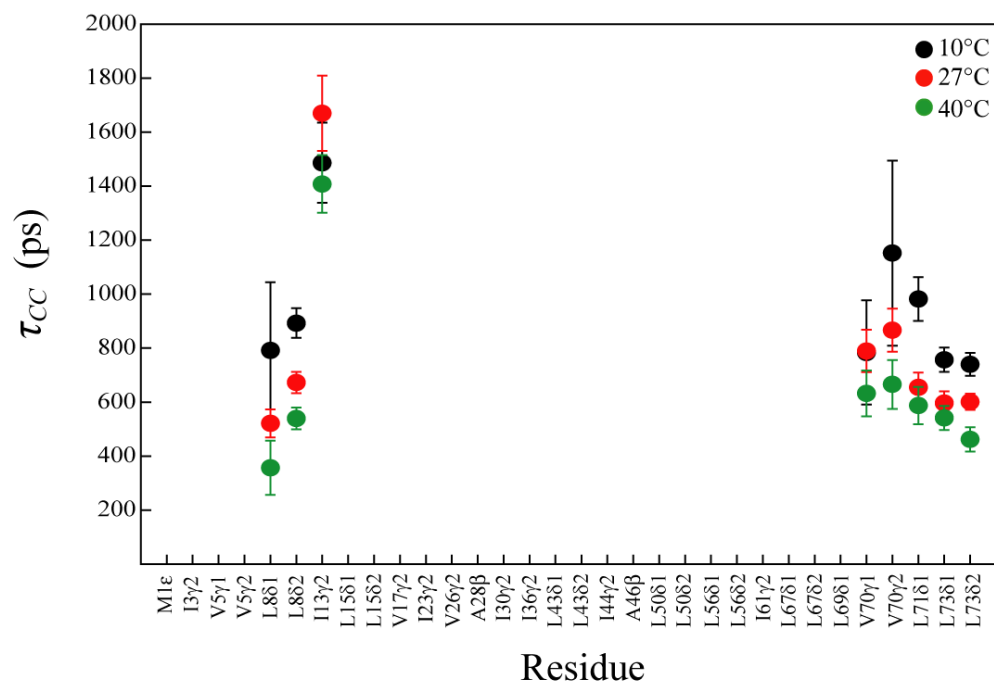


Figure S8. Temperature dependence of the time-scales of slow methyl axis motions/rotameric jumps (τ_{CC} ; ps) for the subset of methyl groups of ubiquitin that have statistically significant reduction in reduced χ^2 when five ^2H relaxation rates (see main text) are best-fit to the spectral density function in eq. 2b. The values of τ_{CC} are shown at 10°C (black circles), 27°C (red circles) and 40°C (green circles).

Supplementary References

- (1) Mittermaier, A.; Kay, L. E. *J. Am. Chem. Soc.* **1999**, *121*, 10608.
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