# SUPPORTING INFORMATION 

for

## By Design: A Macrocyclic 3d-4f Single-Molecule Magnet with Quantifiable Zero-Field Slow Relaxation of Magnetization

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## Supplementary X-ray data



Figure S1: The crystal structure of $\left[\mathrm{Cu}_{3} \mathrm{La}\left(\mathbf{L}^{\mathrm{Bu}}\right)\left(\mathrm{NO}_{3}\right)_{2}(\mathrm{MeOH})_{3}\right]\left(\mathrm{NO}_{3}\right)$. For clarity, non-acidic hydrogen atoms and the non-coordinated nitrate anion have been omitted.

Structure description for $\left[\mathrm{Cu}_{3} \mathrm{La}\left(\mathrm{L}^{\mathrm{Bu}}\right)\left(\mathrm{NO}_{3}\right)_{2}(\mathbf{M e O H})_{3}\right]\left(\mathrm{NO}_{3}\right)$ : Overall, the structure of $\left[\mathrm{Cu}_{3} \mathrm{La}\left(\mathbf{L}^{\mathbf{B u}}\right)\left(\mathrm{NO}_{3}\right)_{2}(\mathrm{MeOH})_{3}\right]\left(\mathrm{NO}_{3}\right)$ is very similar to the terbium(III) analogue, with the large lanthanum(III) ion bound in the central $\mathrm{O}_{6}$ cavity and the smaller copper(II) ions in the outer $\mathrm{N}_{2} \mathrm{O}_{2}$ sites. The lanthanum(III) ion is ten-coordinate, with distorted decahexahedral geometry. All three copper(II) ions are all five-coordinate and are fairly regular square pyramids ${ }^{1}(\tau$ for $\mathrm{Cu}(1)=\mathrm{Cu}(3)=$ $0.15 ; \mathrm{Cu}(2)=0.20)$. One of the nitrate anions is bound in a bidentate fashion to the lanthanum(III) ion, one bridges the lanthanum(III) ion and a copper(II) in a $\mu$-mode and the remaining anion is not coordinated. One methanol molecule is coordinated to the lanthanum(III) ion and two are coordinated in the apical position of two copper(II) ions. The macrocycle is quite curved; the angles made by the plane of an arbitrarily chosen catecholate ring with the planes of the other two rings are $18.8^{\circ}$ and $39.4^{\circ}$.

Refinement details for $\left[\mathrm{Cu}_{3} \mathrm{La}\left(\mathbf{L}^{\mathrm{Bu}}\right)\left(\mathbf{N O}_{3}\right)_{2}(\mathbf{M e O H})_{3}\right]\left(\mathbf{N O}_{3}\right)$ : All hydrogen atoms were placed in calculated positions and rode on the attached non- H atom with $\mathrm{U}(\mathrm{H})=1.2 \mathrm{U}(\mathrm{X})$. In the case of $\mathrm{X}=\mathrm{O}$, the hydrogen selected was the one that formed the best hydrogen bond. All non- H atoms were refined anisotropically except for two catecholate carbon atoms ( $\mathrm{C} 4, \mathrm{C} 16$ ) that presented negative ellipsoid parameters unless refined isotropically. An ISOR restraint was used on four carbon atoms (C6, C14, C20, C27) with a low principal mean square atomic displacement. No disorder is present. For further details see the cif file.

Refinement details for $\left[\mathrm{Cu}_{3} \mathbf{T b}\left(\mathrm{~L}^{\mathrm{Bu}}\right)\left(\mathrm{NO}_{3}\right)_{2}(\mathrm{DMF})\left(\mathrm{H}_{2} \mathrm{O}\right)\right]\left(\mathrm{NO}_{3}\right) \cdot$ DMF: All hydrogen atoms were placed in calculated positions and rode on the attached non-H atom with $\mathrm{U}(\mathrm{H})=1.2 \mathrm{U}(\mathrm{X})$. In the case of $\mathrm{X}=\mathrm{O}$, the hydrogen selected was the one that formed the best hydrogen bond. All non- H atoms were refined anisotropically except for two partial-occupancy carbonyl oxygen atoms (O80, O90) of the non-coordinated DMF solvent molecules. For further details see the cif file.

## Supplementary magnetic data



Figure S2: Scheme of the spin topology in $\left[\mathbf{C u}_{\mathbf{3}} \mathbf{L a}\left(\mathbf{L}^{\mathbf{B u}}\right)\right]$ and its energy diagram in the case of ferromagnetic superexchange interactions.


Figure S3: Temperature dependence of the $\chi T$ product of $\left[\mathbf{C u}_{3} \mathbf{L a}\left(\mathbf{L}^{\mathrm{Bu}}\right)\right]$ in a 1000 Oe dc field (circles) and a fit to the data (redline) based on the model shown in Figure S2.


Figure S4: $M$ vs. $H$ (left) and $M$ vs. $H / T$ (right) data for $\left[\mathbf{C u} \mathbf{u}_{3} \mathbf{T b}\left(\mathbf{L}^{\mathrm{Bu}}\right)\right]$ at the indicated temperatures, scanning at about $100-200 \mathrm{Oe} \mathrm{min}.{ }^{-1}$ The solid lines are the calculated magnetization from the ab initio calculations described in the text.


Figure S5: Field dependence of magnetisation for $\left[\mathbf{Z n}_{3} \mathbf{T b}\left(\mathbf{L}^{\mathrm{Bu}}\right)\right]$ (left) and $\left[\mathrm{Cu}_{3} \mathbf{L a}\left(\mathbf{L}^{\mathrm{Bu}}\right)\right]$ (right) at the temperatures indicated, scanning at about $100-200 \mathrm{Oe} \mathrm{min}^{-1}$.


Figure S6: Temperature dependence of the in-phase (left) and out-of-phase (right) components of the ac susceptibility for $\left[\mathbf{C u}_{3} \mathbf{T b}\left(\mathbf{L}^{\mathrm{Bu}}\right)\right]$ at the frequencies indicated in zero-dc field.


Figure S7: Temperature dependence of the in-phase and out-of-phase components of ac susceptibility for $\left[\mathbf{Z n}_{3} \mathbf{T b}\left(\mathbf{L}^{\mathbf{B u}}\right)\right]$ at 1000 Hz in zero-dc field.

## Additional ab initio results

Table S1. A comparison of low-lying level energies (in $\mathrm{cm}^{-1}$ ) corresponding to the free ion $J=6$ multiplet of central Tb ions in $\left[\mathbf{C u}_{3} \mathbf{T b}\left(\mathbf{L}^{\mathrm{Pr}}\right)\right]$ and $\left[\mathbf{C u}_{3} \mathbf{T b}\left(\mathbf{L}^{\mathrm{Bu}}\right)\right]$. The states defining the axiality are shown with bold: the axiality of Tb ion increases with the decrease of the energy gap.

| $\left[\mathrm{Cu}_{3} \mathrm{~Tb}\left(\mathrm{~L}^{\mathrm{Pr}}\right)\right]$ | $\left[\mathrm{Cu}_{3} \mathrm{~Tb}\left(\mathrm{~L}^{\mathrm{Bu}}\right)\right]$ |
| :---: | :---: |
| $\mathbf{0 . 0 0 0}$ | $\mathbf{0 . 0 0 0}$ |
| $\mathbf{4 . 2 4 8}$ | $\mathbf{0 . 1 7 4}$ |
| 59.496 | 42.060 |
| $\mathbf{9 1 . 3 2 1}$ | 42.688 |
| $\mathbf{1 0 2 . 0 1 8}$ | 76.038 |
| $\mathbf{1 5 3 . 3 0 4}$ | 86.706 |
| $\mathbf{1 6 4 . 9 1 4}$ | 109.163 |
| $\mathbf{2 1 1 . 5 4 9}$ | 110.691 |
| $\mathbf{2 2 0 . 3 6 6}$ | 128.768 |
| $\mathbf{2 6 4 . 6 9 1}$ | 182.476 |
| $\mathbf{2 6 7 . 0 4 9}$ | 185.891 |
| $\mathbf{4 7 2 . 0 3 3}$ | 406.847 |
| $\mathbf{4 7 2 . 9 6 8}$ | 407.827 |
| $\ldots$ | $\ldots$ |

Table S2. Calculated exchange energies (in K ) and the corresponding $g$ tensors of the low-lying exchange Kramers doublets of the previously reported $\left[\mathbf{C u} \mathbf{3} \mathbf{T b}\left(\mathbf{L}^{\mathrm{Pr}}\right)\right]$ complex ${ }^{2}$ and the present $\left[\mathbf{C u}_{3} \mathbf{T b}\left(\mathbf{L}^{\mathrm{Bu}}\right)\right]$ compounds.

|  | $\left[\mathrm{Cu}_{3} \mathrm{~Tb}\left(\mathrm{~L}^{\mathrm{Pr}}\right)\right]^{4 \mathrm{~d}}$ |  |  | $\left[\mathrm{Cu}_{3} \mathrm{~Tb}\left(\mathrm{~L}^{\mathrm{Bu}}\right)\right]$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $J(C u-$ <br> $\mathbf{C u}) / \mathbf{k}_{\mathrm{B}}$ <br> $J\left(\mathrm{Cu}^{-}\right.$ <br> $\mathrm{Tb}) / \mathbf{k}_{\mathrm{B}}$ | 1.41 K |  |  |  | 1.37 K |  |
| Kramers Doublets | $\begin{aligned} & \text { energy } \\ & \mathrm{cm}^{-1} \end{aligned}$ | $g$ tensor |  | $\begin{aligned} & \text { energy } \\ & \mathrm{cm}^{-1} \end{aligned}$ | $g$ tensor |  |
| 1 | 0.000 | $\mathrm{g}_{\mathrm{X}}$ <br> $\mathrm{g}_{\mathrm{Y}}$ <br> $\mathrm{g}_{\mathrm{z}}$ |  | 0.000 | $\mathrm{g}_{\mathrm{X}}$ $\mathrm{g}_{\mathrm{Y}}$ $\mathrm{g}_{\mathrm{Z}}$ |  |
| 2 | 7.105 | $\begin{aligned} & \mathrm{g}_{\mathrm{X}} \\ & \mathrm{~g}_{\mathrm{Y}} \\ & \mathrm{~g}_{Z} \end{aligned}$ | $\begin{array}{r} 2.3663 \\ 2.7211 \\ 16.3099 \\ \hline \end{array}$ | 8.333 | $\begin{aligned} & \mathrm{g}_{\mathrm{X}} \\ & \mathrm{~g}_{\mathrm{Y}} \\ & \mathrm{~g}_{\mathrm{Z}} \end{aligned}$ | $\begin{array}{r} 0.0906 \\ 0.1029 \\ 19.7233 \\ \hline \end{array}$ |
| 3 | 9.064 | $\begin{aligned} & \mathrm{g}_{\mathrm{X}} \\ & \mathrm{~g}_{\mathrm{Y}} \\ & \mathrm{~g}_{\mathrm{Z}} \\ & \hline \end{aligned}$ | $\begin{array}{r} 1.0807 \\ 1.3485 \\ 16.5745 \\ \hline \end{array}$ | 10.220 | $\begin{aligned} & \mathrm{g}_{\mathrm{X}} \\ & \mathrm{~g}_{\mathrm{Y}} \\ & \mathrm{~g}_{\mathrm{I}} \end{aligned}$ | $\begin{array}{r} 0.0433 \\ 0.0509 \\ 19.7362 \\ \hline \end{array}$ |
| 4 | 9.068 | $\mathrm{g}_{\mathrm{X}}$ <br> $\mathrm{g}_{\mathrm{Y}}$ <br> $\mathrm{g}_{\mathrm{z}}$ |  | 10.453 | $\mathrm{g}_{\mathrm{x}}$ <br> $\mathrm{g}_{\mathrm{Y}}$ <br> $\mathrm{g}_{\mathrm{z}}$ |  |
| 5 | 16.026 | $\mathrm{g}_{\mathrm{x}}$ <br> $\mathrm{g}_{\mathrm{Y}}$ <br> $\mathrm{g}_{\mathrm{z}}$ | $\begin{array}{r} 1.8793 \\ 2.6203 \\ 11.9709 \\ \hline \end{array}$ | 16.988 | $\mathrm{g}_{\mathrm{x}}$ <br> $g_{Y}$ <br> $\mathrm{g}_{\mathrm{z}}$ |  |
| 6 | 17.920 | $\mathrm{g}_{\mathrm{x}}$ <br> $\mathrm{g}_{\mathrm{Y}}$ <br> $\mathrm{g}_{\mathrm{z}}$ | $0.9637$ $1.2247$ $12.2555$ | 18.916 | $\mathrm{g}_{\mathrm{x}}$ <br> $\mathrm{g}_{\mathrm{Y}}$ <br> $\mathrm{g}_{\mathrm{z}}$ |  |
| 7 | 17.924 | $\mathrm{g}_{\mathrm{X}}$ <br> $\mathrm{g}_{\mathrm{Y}}$ <br> $\mathrm{g}_{\mathrm{z}}$ | 0.9765 <br> 1.1664 <br> 12.1787 | 19.147 | $\mathrm{g}_{\mathrm{x}}$ <br> $\mathrm{g}_{\mathrm{Y}}$ <br> $\mathrm{g}_{\mathrm{z}}$ |  |
| 8 | 23.608 | $\begin{aligned} & \mathrm{g}_{\mathrm{X}} \\ & \mathrm{~g}_{\mathrm{Y}} \\ & \mathrm{~g}_{\mathrm{Z}} \end{aligned}$ | $\begin{aligned} & \hline 0.1508 \\ & 0.1766 \\ & 9.5883 \\ & \hline \end{aligned}$ | 26.027 | $\mathrm{g}_{\mathrm{X}}$ $\mathrm{g}_{\mathrm{Y}}$ $\mathrm{g}_{\mathrm{Z}}$ | $\begin{aligned} & 0.0066 \\ & 0.0369 \\ & 11.2530 \end{aligned}$ |

## References

1. A. W. Addison, T. N. Rao, J. Reedijk, J. van Rijn and G. C. Vershoor, J. Chem. Soc., Dalton Trans., 1984, 1349-1356.
2. H. L. C. Feltham, R. Clérac, L. Ungur, V. Vieru, L. F. Chibotaru, A. K. Powell and S. Brooker, Inorg. Chem., 2012, 51, 10603-10612.
