#### **Supporting Information**

# Water-wettable polypropylene fibers by facile surface treatment based on soy proteins

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#### Johannsmann model to calculate the adsorbed mass

In the Johannsmann model, the shift in QCM frequency ( $\Delta f$ ) is related to the complex shear compliance of the thin film ( $\hat{J}(\omega)$ ) by Eq. 1:<sup>1</sup>

$$\Delta \hat{f} \approx f_0 \frac{1}{\pi Z_q} (\omega \rho d + \hat{J}(\omega) \frac{\omega^3 \rho^2 d^3}{3}$$
 (1)

where  $Z_q = \sqrt{\rho_q \mu_q} = 8.8 * 10^{-6} kg * m^{-2} s^{-1}$  is the bulk acoustic impedance of crystalline AT-cut quartz,  $\omega = 2\pi f$  and f is the resonance frequency of the crystal;  $\rho$  is the density of the fluid;  $\rho_q$  is the specific density of the quartz,  $\mu_q$  is the elastic shear modulus of quartz and d is the thickness of the film.

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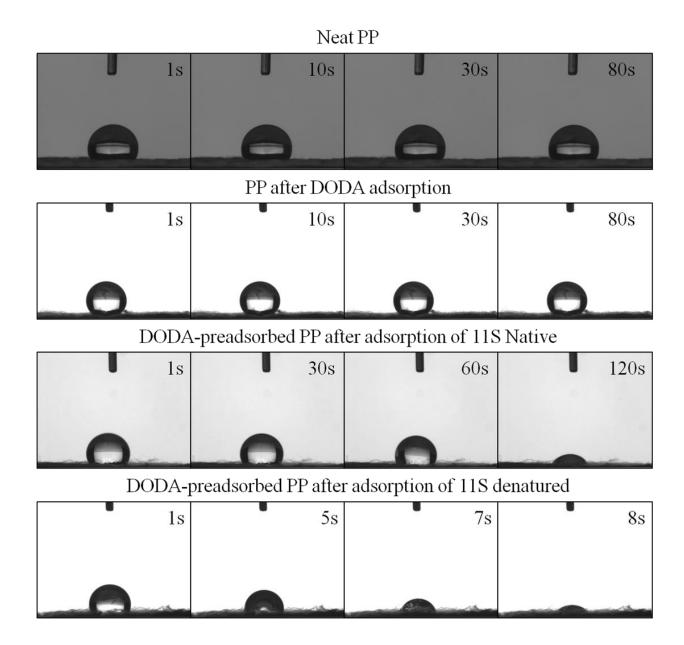
Equation **A** can be rearranged by using the equivalent mass  $m^*$  (Eq. 2), which can be used to calculate the true adsorbed mass under the assumption that  $\hat{J}(\omega)$  is independent of  $\omega^2$  resulting in equation 3:

$$m^* = -\frac{Z_q}{2f_o} \frac{\Delta_f}{f} \tag{2}$$

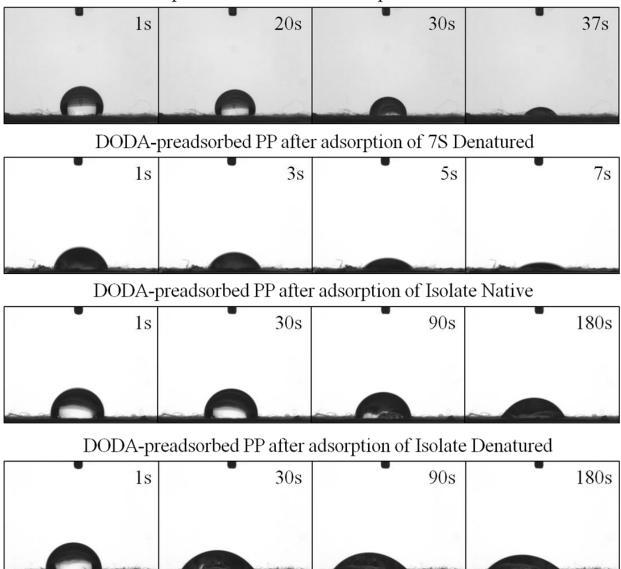
$$m^* = m^o \left( 1 + \hat{J}(\omega) \frac{\omega^2 \rho d^2}{3} \right) \Rightarrow m^* = m^o \left( 1 + \hat{J}(\omega) \frac{f^2 \rho d^2}{3} \right)$$
 (3)

From equation 3 the true mass can be calculated from a plot of  $m^*$  againts  $f^2$  to give  $m^o$  as the intercept.

## Contact angle images of nonwoven surfaces



## DODA-preadsorbed PP after adsorption of 7S Native



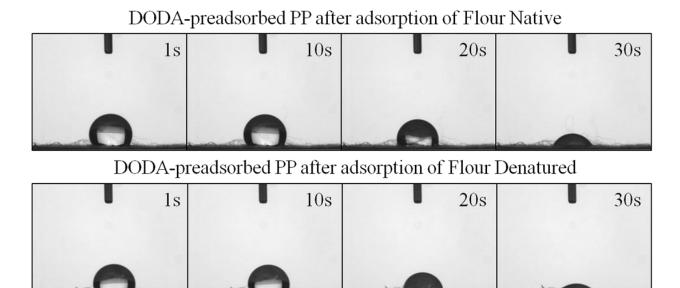


Figure SI-1. Contact angle images for the PP substrates after different treatments studied.

#### References

- 1. Johannsmann, D.; Mathauer, K.; Wegner, G.; Knoll, W., *Phys. Rev. B: Condens. Matter* **1992,** 46, 7808-7815.
- 2. Naderi, A.; Claesson, P. M., Langmuir 2006, 22, 7639-7645.