# Transition Metal Embedded Two-Dimensional 

# $\mathrm{C}_{3} \mathrm{~N}_{4}$-graphene Nanocomposite: A Multifunctional Material 

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Figure S1. (a) Top and (b) side-view of optimized structure of g-C3N4@graphene. Notice that adhesion of buckled g-C ${ }_{3} \mathrm{~N}_{4}$ on top of graphene makes it (i.e. $\mathrm{g}-\mathrm{C}_{3} \mathrm{~N}_{4}$ ) planar.


Figure S2: Band structure of $g-\mathrm{C}_{3} \mathrm{~N}_{4} @$ graphene are plotted considering high symmetry Kpoints $(\Gamma(0,0,0), \mathrm{M}(1 / 2,1 / 2,0), \mathrm{K}(2 / 3,1 / 3,0)$. Fermi level is scaled to zero. The zoomed picture of dirac-cone at $\Gamma$-point (graphene fold K-point) clearly shows band-gap opening. Symbols: blue dashed and green dotted lines show Fermi levels and high symmetry K-points, respectively.


Figure S3. $2 \times 2$ supercell of $\mathrm{TM}-\mathrm{C}_{3} \mathrm{~N}_{4} @$ graphene. This cell has been used to find out the magnetic ground state of these sheets. $\mathrm{d}_{\mathrm{TM} \_ \text {TM_hori }}$ and $\mathrm{d}_{\text {TM_TM_dia }}$ are the distances between two TM atoms at horizontal and diagonal direstion, respectively. $\mathrm{d}_{\mathrm{M} \text {-Nedge }}(a-\mathrm{f})$ are the distances between $\mathrm{N}_{\text {edge }}$ and TM.


Figure S4. Demonstrations of (a) ferromagnetic and (b) antiferromagnetic coupling between Fe atoms of $\mathrm{Fe}-\mathrm{C}_{3} \mathrm{~N}_{4}$.@graphene. Isosurface at a value of $0.1 \mathrm{e} / \AA^{3}$ is taken. Up and down spin densities are represented as orange and yellow coloured surfaces, respectively.

## Detail of d-p exchange in TM-g-C $\mathbf{3}_{\mathbf{3}} \mathbf{N}_{\mathbf{4}} @$ graphene:

According to Goodenough-Kanamori-Anderson rules, when the interacting magnetic dorbitals of TM couple with p orbitals of ligand, depending upon their symmetry, ferromagnetic or antiferromagnetic ground state appear. When the orbitals of interest interact but total overlap is zero due to symmetry of lobes, ferromagnetism appears Among various such type of situations, one common scenario is when the interacting spin-polarized dorbitals and $p_{\pi}\left(i . e . p_{z} / p_{y}\right)$ orbitals remain perpendicular to each other, they produce a zero overlap and consequently a ferromagnetic coupling.

For Cr- $\mathrm{C}_{3} \mathrm{~N}_{4} @$ graphene, $\mathrm{d}_{\mathrm{xy}}$ and $\mathrm{d}_{\mathrm{x}}{ }^{2}-\mathrm{y}{ }^{2}$ of Cr and $\mathrm{p}_{\mathrm{z}}$ and $\mathrm{p}_{\mathrm{y}}$ of $\mathrm{N}_{\text {edge }}$ atoms get involved in magnetic coupling interaction. The pictorial representations of orbital overlaps as well as resulting magnetic coupling natures are tabulated below. Among four type of overlap, only one kind results antiferromagnetic coupling whereas other three gives ferromagnetic interaction. As a result, the Cr atoms at $\mathrm{Cr}-\mathrm{C}_{3} \mathrm{~N}_{4} @$ graphene show a ferromagnetic ground state.

As can be seen in Figure S , the interacting Fe and ligand orbitals in Fe $\mathrm{C}_{3} \mathrm{~N}_{4} @$ graphene are $\mathrm{d}_{\mathrm{xz}} / \mathrm{d}_{\mathrm{yz}}$ and $\mathrm{p}_{\mathrm{y}}$, respectively. It is evident from orbital overlap pictures that the effective d-p overlap is zero here and consequently the exchange becomes ferromagnetic in nature.

| Interacting Orbital | Magnetic <br> Coupling <br> Nature | Interacting Orbital | Magnetic <br> Coupling <br> Nature |  |
| :---: | :---: | :---: | :---: | :---: |
|  | AFM | FM | FM |  |

Orbital pictures for d-p overlaps in Cr-C $\mathbf{C}_{3} \mathbf{N}_{4} @$ graphene

| Interacting Orbital | Magnetic <br> Coupling <br> Nature | Interacting Orbital | Magnetic <br> Coupling <br> Nature |
| :---: | :---: | :---: | :---: |
|  |  |  |  |

Orbital pictures for d-p overlaps in Fe- $\mathbf{C}_{3} \mathrm{~N}_{4} @$ graphene


Figure S5. (a) pDOS of d-orbitals on the Fe atom of $\mathrm{Fe}-\mathrm{C}_{3} \mathrm{~N}_{4} @$ graphene (symbols: solid black, dashed red, dashed doted maroon, dashed double doted deep green and double dashed doted cyan represent $\mathrm{d}_{\mathrm{xy}}, \mathrm{d}_{\mathrm{yz}}, \mathrm{d}_{\mathrm{z}}^{2}, \mathrm{~d}_{\mathrm{xz}}$ and $\mathrm{d}_{\mathrm{x}-\mathrm{y}}{ }^{2}$ orbitals, respectively). (b) pDOS of p-orbitals on the $\mathrm{N}_{\text {edge }}$ of same structure has been plotted. Note, up and down red coloured arrows represent majority and minority spins, respectively.


Figure S6. Calculated band structures for $\mathrm{TM}-\mathrm{g}-\mathrm{C}_{3} \mathrm{~N}_{4}$ where TM is (a) V , (b) Cr , (c) Mn , (d) Co , (e) Ni and (f) Cu .

## Adhesion Energy

To know the strength of interaction between TM@g-C3 $\mathrm{N}_{4}$ and graphene, we have calculated interface adhesion energy using following formula,

$$
\mathrm{E}_{\text {adhesion }}=\mathrm{E}_{\text {composite }}-\left(\mathrm{E}_{\text {g-C3N4-TM }}+\mathrm{E}_{\text {graphene }}\right)
$$

where $\mathrm{E}_{\text {composite }}$, $\mathrm{E}_{\mathrm{g}-\mathrm{C} 3 \mathrm{~N} 4-\mathrm{TM}}$ and $\mathrm{E}_{\text {graphene }}$ are the energies of the $\mathrm{TM}-\mathrm{g}-\mathrm{C}_{3} \mathrm{~N}_{4} @$ graphene, TM@g-C $\mathrm{C}_{3} \mathrm{~N}_{4}$ and graphene, respectively. Calculated values are tabulated in Table S1.

Table S1. Structural details of fully optimized geometry of TM- $\mathrm{C}_{3} \mathrm{~N}_{4} @$ graphene. Distances between two layers, two TMs, TM and $\mathrm{N}_{\text {edge }}$ ( $\mathrm{M}-\mathrm{N}$; denoted as a-f as can be seen in Figure S3) at $g-\mathrm{C}_{3} \mathrm{~N}_{4}$ layer are given. The Fermi level shift due to charge transfer and adhesion energies/unit cell are also tabulated for all systems.

| Metal |  | V | Cr | Mn | Fe | Co | Ni | $\mathbf{C u}$ | Zn |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Interlayer distance$(\AA)$ |  | 2.98 | 2.97 | 2.98 | 2.97 | 2.95 | 2.97 | 2.96 | 2.95 |
| $\mathrm{d}_{\text {TM_TM_hori }}(\AA)$ |  | 7.28 | 7.28 | 7.29 | 7.29 | 7.28 | 7.32 | 7.29 | 7.29 |
| d TM _TM_dia $(\AA)$ |  | 7.28 | 7.28 | 7.29 | 7.29 | 7.27 | 7.30 | 7.29 | 7.29 |
| $\mathrm{d}_{\text {M-Nedge }}$ | $\begin{aligned} & \mathrm{a}(1-\mathrm{TM}) \\ & \mathrm{b}(2-\mathrm{TM}) \\ & \mathrm{c}(3-\mathrm{TM}) \\ & \mathrm{d}(4-\mathrm{TM}) \\ & \mathrm{e}(5-\mathrm{TM}) \\ & \mathrm{f}(6-\mathrm{TM}) \end{aligned}$ | $\begin{aligned} & \hline 2.44 \\ & 2.44 \\ & 2.44 \\ & 2.44 \\ & 2.44 \\ & 2.44 \end{aligned}$ | $\begin{aligned} & \hline 2.72 \\ & 2.57 \\ & 2.53 \\ & 2.18 \\ & 2.16 \\ & 2.63 \end{aligned}$ | $\begin{aligned} & \hline 2.46 \\ & 2.46 \\ & 2.45 \\ & 2.45 \\ & 2.45 \\ & 2.46 \end{aligned}$ | $\begin{aligned} & \hline 2.19 \\ & 2.43 \\ & 2.70 \\ & 2.72 \\ & 2.49 \\ & 2.21 \end{aligned}$ | $\begin{aligned} & \hline 2.71 \\ & 2.75 \\ & 2.65 \\ & 2.18 \\ & 2.16 \\ & 2.44 \end{aligned}$ | $\begin{aligned} & \hline 1.95 \\ & 2.51 \\ & 3.00 \\ & 3.00 \\ & 2.58 \\ & 1.95 \end{aligned}$ | $\begin{aligned} & 2.47 \\ & 2.48 \\ & 2.47 \\ & 2.46 \\ & 2.46 \\ & 2.46 \end{aligned}$ | $\begin{aligned} & \hline 2.47 \\ & 2.48 \\ & 2.48 \\ & 2.47 \\ & 2.46 \\ & 2.46 \end{aligned}$ |
| Transferred electrons to graphene ( $\mathrm{e}^{-}$) |  | 0.18 | 0.15 | 0.23 | 0.16 | 0.15 | 0.12 | 0.10 | 0.14 |
| Fermi level shift$(\mathrm{eV})$ |  | 0.41 | 0.38 | 0.51 | 0.43 | 0.37 | 0.33 | 0.27 | 0.36 |
| Adhesion Energy$(\mathrm{eV})$ |  | -0.22 | -0.24 | -0.29 | -0.23 | -0.31 | -0.28 | -0.29 | -0.25 |

## Calculation of Magnetic Coupling Constants:

We have calculated magnetic coupling constant $J$ by using following Heisenberg Hamiltonian,

$$
H=\sum_{\langle i j\rangle} J_{i j}\left(S_{i} . S_{j}\right)
$$

considering rhombic ( $2 \times 2$ ) supercell and imposing periodic boundary condition.
The H turns out to be
$H=J\left(\mathrm{~s}_{1} \mathrm{~s}_{2}+\mathrm{s}_{2} \mathrm{~s}_{3}+\mathrm{s}_{3} \mathrm{~s}_{4}+\mathrm{s}_{4} \mathrm{~s}_{1}+\mathrm{s}_{2} \mathrm{~s}_{4}+\mathrm{s}_{3} \mathrm{~s}_{4}\right)$
Now, we can write total spin as,

$$
\mathrm{S}_{\mathrm{T}}^{2}=\left(\mathrm{s}_{1}+\mathrm{s}_{2}+\mathrm{s}_{3}+\mathrm{s}_{4}\right)^{2}=\mathrm{s}_{1}{ }^{2}+\mathrm{s}_{2}{ }^{2}+\mathrm{s}_{3}{ }^{2}+\mathrm{s}_{4}{ }^{2}+2\left(\mathrm{~s}_{1} \mathrm{~s}_{2}+\mathrm{s}_{2} \mathrm{~s}_{3}+\mathrm{s}_{3} \mathrm{~s}_{4}+\mathrm{s}_{4} \mathrm{~s}_{1}+\mathrm{s}_{2} \mathrm{~s}_{4}+\mathrm{s}_{3} \mathrm{~s}_{4}\right)
$$

For spin state $S, S^{2}$ has eigen value of $S(S+1) \hbar^{2}$. We have considered $\hbar=1$ here after.
Thus, in terms of eigen values of above mentioned Hamiltonian we can write,

$$
\mathrm{E}=J / 2\left[\mathrm{~S}_{\mathrm{T}}\left(\mathrm{~S}_{\mathrm{T}}+1\right)-\mathrm{s}_{1}\left(\mathrm{~s}_{1}+1\right)-\mathrm{s}_{2}\left(\mathrm{~s}_{2}+1\right)-\mathrm{s}_{3}\left(\mathrm{~s}_{3}+1\right)-\mathrm{s}_{4}\left(\mathrm{~s}_{4}+1\right)\right]
$$

This energy equation is quite general and now depending upon TM, we will consider different $s_{1}, s_{2}, s_{3}$ and $s_{4}$ values and derive the exchange coupling constant.

For Cr-g- $\mathbf{C}_{3} \mathbf{N}_{4}$, where $\mathrm{s}_{1}=\mathrm{s}_{2}=\mathrm{s}_{3}=\mathrm{s}_{4}=2$ (as $\mathrm{Cr}^{+2}$ has 4 unpaired electrons);
E comes out to,
$\mathrm{E}=J / 2\left[\mathrm{~S}_{\mathrm{T}}\left(\mathrm{S}_{\mathrm{T}}+1\right)-24\right]$
Therefore we can write, energy for antiferromagnetic configuration taking $\mathrm{S}_{\mathrm{T}}=0$;

$$
\mathrm{E}_{\mathrm{AFM}}=-12 \mathrm{~J}
$$

For ferromagnetic configuration, $\mathrm{S}_{\mathrm{T}}=8$,
Thus, $\quad \mathrm{E}_{\mathrm{FM}}=24 \mathrm{~J}$
So, $\Delta \mathrm{E}_{\text {ex }}=\mathrm{E}_{\mathrm{FM}}-\mathrm{E}_{\mathrm{AFM}}$

$$
=36 \mathrm{~J}
$$

From DFT calculation, $\Delta \mathrm{E}_{\text {ex }}$ for this system appears as -179 meV .
Thus, $\boldsymbol{J}=\mathbf{- 4 . 9 7} \mathbf{~ m e V}$

Next, for $\mathbf{F e}-\mathrm{g}-\mathrm{C}_{3} \mathbf{N}_{4}$, where $\mathrm{s}_{1}=\mathrm{s}_{2}=\mathrm{s}_{3}=\mathrm{s}_{4}=2$ (as Fe ${ }^{+2}$ has 4 unpaired electrons); $\mathrm{E}=J / 2\left[\mathrm{~S}_{\mathrm{T}}\left(\mathrm{S}_{\mathrm{T}}+1\right)-24\right]$

For antiferromagnetic configuration, $\mathrm{S}_{\mathrm{T}}=0$
$\mathrm{E}_{\mathrm{AFM}}=12 \mathrm{~J}$
For ferromagnetic configuration, $\mathrm{S}_{\mathrm{T}}=8$
$\mathrm{E}_{\mathrm{FM}}=24 \mathrm{~J}$
So, $\Delta \mathrm{E}_{\text {ex }}=36 \mathrm{~J}$
From DFT calculation, $\Delta \mathrm{E}_{\text {ex }}=-201 \mathrm{meV}$
Therefore, $\boldsymbol{J}=\mathbf{- 5 . 5 8} \mathbf{~ m e V}$

