

Supporting information for

Electronic Structure and Photoelectrochemical Properties of an Ir-Doped SrTiO₃ Photocatalyst

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Materials and methods

Figure S1 shows the X-ray diffraction (XRD) patterns of Ir(x at %):SrTiO₃ (x = 1, 3, 5) powders used in this study, together with reference data for IrO₂, Sr₃Ti₂O₇, and SrTiO₃. The XRD patterns of the powder samples match the SrTiO₃ reference but may also contain a small amount of the layered perovskite Sr₃Ti₂O₇. No diffraction peaks of IrO₂ were observed. These powders were processed to pellet shapes for use as targets in pulsed laser deposition (PLD) growth of epitaxial thin films.

Typical XRD patterns of Ir(5 at %):SrTiO₃ thin film samples deposited on SrTiO₃ (001) substrates are shown in Figure S2. The samples were deposited at 700 °C under oxygen pressures of 10⁻¹ or 10⁻⁶ Torr, producing Ir⁴⁺ and Ir³⁺ doped SrTiO₃ films, respectively. All films were epitaxially grown on SrTiO₃ (001) substrates. The slight shift of the XRD peaks of the film grown at 10⁻⁶ Torr is caused by the presence of oxygen defects that occur due to the low growth pressure¹ and also due to the difference of ionic radii of Ir³⁺ (0.68 Å) and

Ir^{4+} (0.625 Å).² The presence of oxygen vacancies stabilizes the Ir^{3+} dopant state at the Ti^{4+} lattice site in SrTiO_3 by maintaining charge neutrality.

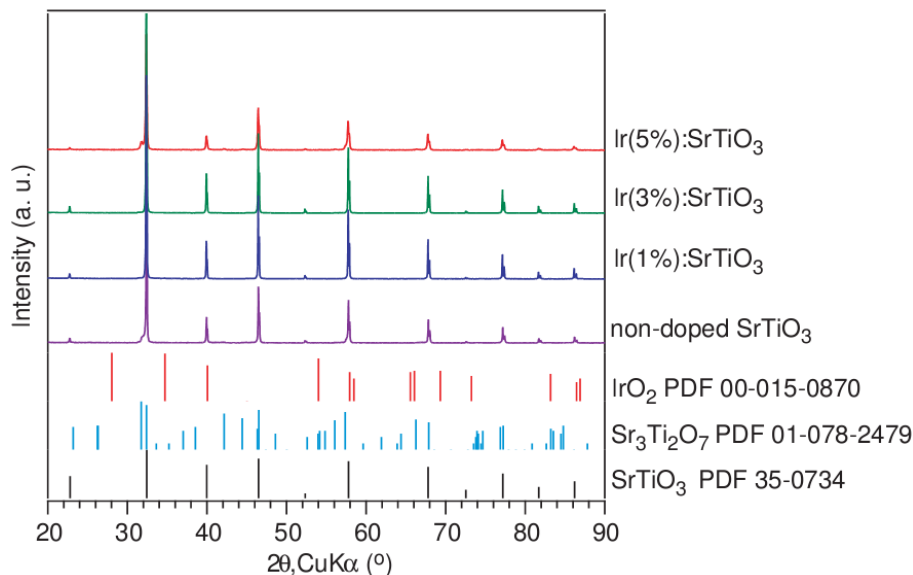


Figure S1. XRD patterns of Ir:SrTiO₃ powder samples. From top to bottom, data for Ir⁴⁺(5, 3, 1%):SrTiO₃ and non-doped SrTiO₃ are shown together with reference data for IrO₂, Sr₃Ti₂O₇, and SrTiO₃.

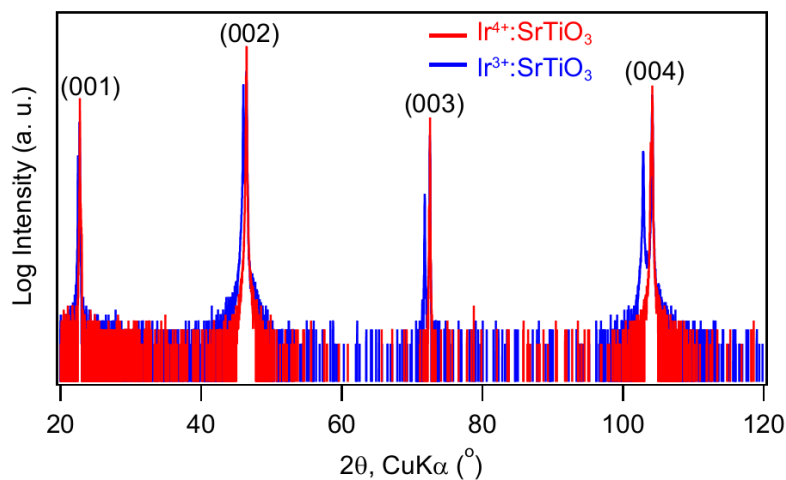


Figure S2. XRD patterns of Ir(5 at %):SrTiO₃ thin film samples grown on SrTiO₃ (001) substrates. The samples were deposited at 700 °C and an oxygen pressure of 10⁻¹ (red) and 10⁻⁶ Torr (blue).

X-ray photoelectron spectroscopy (XPS)

High resolution X-ray photoelectron spectroscopy was performed using the synchrotron light source at BL2 of the Photon Factory (Tsukuba, Japan). Figure S3 shows the atomic subshell photoionization cross sections of Ir 5d, O 2p, and Ti 3d electrons as a function of the incident photon energy, calculated from literature values.³ The Ir 5d / O 2p and Ir 5d / Ti 3d spectral ratios increase monotonically with the increase of the photon energy from 200 to 1600 eV. Therefore, element-selective photoelectron spectra for Ir 5d can be obtained by changing the photon energy from 600 to 1400 eV and looking for spectral features that increase in intensity at higher photon energies (Figure 5 in the main text).

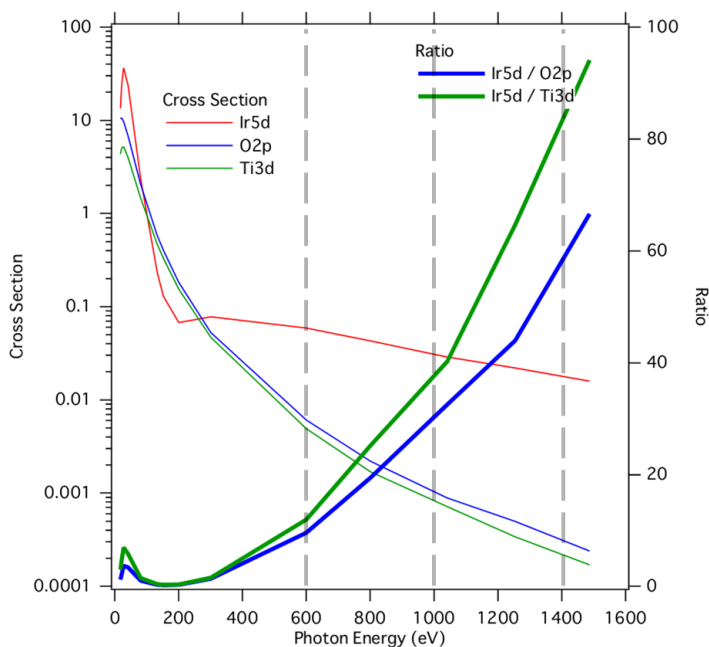


Figure S3. Atomic subshell photoionization cross sections for Ir 5d, O 2p, and Ti 3d electrons as a function of the incident X-ray photon energy.

Radiation damage was clearly observed when films were irradiated by the synchrotron light at $h\nu = 600$ eV in ultrahigh vacuum ($\sim 5 \times 10^{-10}$ Torr). Figure S4 shows the first and 10th scans of the Ir 4f core level photoelectron spectra for a film sample deposited at 700 °C and 10^{-1} Torr. The Ir 4f core-level spectrum is a doublet consisting of Ir 4f_{7/2} and Ir 4f_{5/2} components. Ten scans were measured in a sequence after selecting a fresh point on the film surface. The measurement time for a single scan was 1.5 min. It is clear that the Ir⁴⁺ is easily reduced to Ir³⁺ under strong X-ray irradiation probably due to the photodesorption of

oxygen from the crystal lattice. Therefore, the film composition was analyzed by measuring Ir 4f core level spectra with a laboratory Mg K α x-ray source (JEOL; JPS-9010MC, $h\nu = 1253.6$ eV) instead of a synchrotron light source.

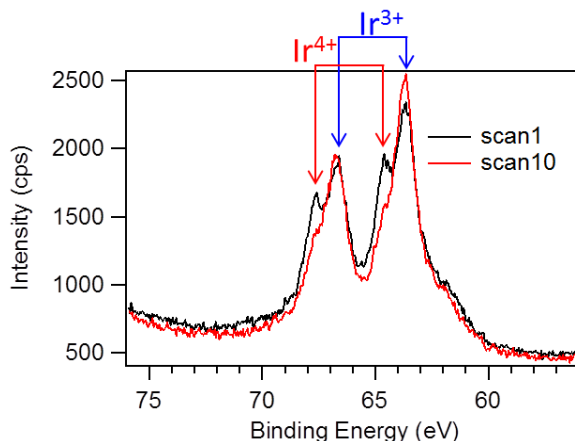


Figure S4. The first and the 10th scan of the Ir 4f core level photoelectron spectra for a film sample deposited at 700 °C and 10⁻¹ Torr. The scans were taken continuously at the same sample position, with each scan lasting 1.5 min. The photon energy was 600 eV.

The Fermi level position of doped SrTiO₃ can be inferred from the valence band spectral feature shifts in XPS spectra. Figure S5 shows synchrotron X-ray photoelectron spectra of the valence band of Ir(5%):SrTiO₃ films deposited at 700 °C, 10⁻¹ Torr and 700 °C, 10⁻⁶ Torr, an oxygen-deficient Rh(5%):SrTiO₃ film deposited at 700 °C, 10⁻¹ Torr and annealed at 600 °C, 10⁻⁶ Torr for 2 hours, and a Nb(0.05wt%):SrTiO₃ substrate. The spectral intensity differences among the samples are caused by cross section differences that depend on the incident photon energy, but these differences can be ignored when discussing the energy level positions of the O2p- π and O2p- σ components. In this case, the Rh:SrTiO₃ sample was intentionally reduced and contained both Rh³⁺ and Rh⁴⁺ (Rh³⁺/Rh⁴⁺ \approx 1) features. It showed a -0.5 eV shift of the VB from the other spectra, indicating that the Fermi level of Rh:SrTiO₃ is 0.5 eV deeper than that of Nb:SrTiO₃, which leads to the p-typeness observed in this material. However, the binding energy positions of O2p- π and O2p- σ of Ir:SrTiO₃ were at the same positions of those of Nb:SrTiO₃, regardless of the Ir valence, showing that the Fermi level positions of Ir:SrTiO₃ and Nb:SrTiO₃ are similar, although the first-principles calculations predicted that the Fermi level of Ir:SrTiO₃ should be deeper than that of non-doped SrTiO₃ (Figure 4). One possible reason for the difference between the results

obtained from XPS and simulations might be related to oxygen defect formation in the samples, especially in the surface layer of the films.

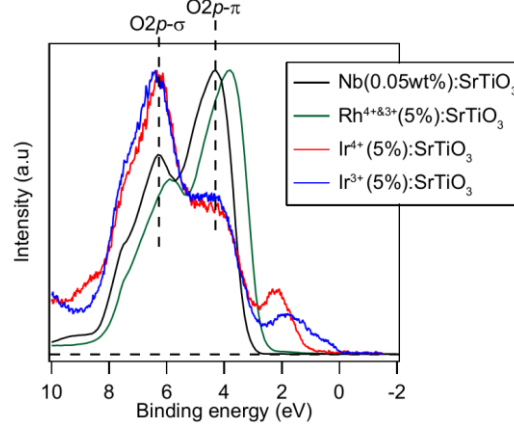


Figure S5. Synchrotron X-ray photoelectron spectra of the valence band of Ir(5%):SrTiO₃ films deposited at 700 °C, 10⁻¹ Torr (red) and 700 °C, 10⁻⁶ Torr (blue), Rh(5%):SrTiO₃ film deposited at 700 °C, 10⁻¹ Torr and annealed at 600 °C, 10⁻⁶ Torr for 2 hours (green), and Nb(0.05wt%):SrTiO₃ substrate. The photon energy was set at 1000 eV for Ir^{4+/3+}:SrTiO₃ and 160 eV for Rh:SrTiO₃ and Nb:SrTiO₃ samples. The binding energy was referenced to the Au Fermi level position at 0 eV.

Density of states calculations

A $3 \times 3 \times 3$ cell was used in the calculations. One of the Ti atoms was substituted by Ir⁴⁺ to calculate the density of states (DOS) of an isovalent sample. While the GGA/PBE96⁴ functional underestimated the band-gap width (ca. 1.5 eV for SrTiO₃) and failed to describe the in-gap states, the hybrid functional HSE06 accurately reproduced the band-gap width (ca. 3.0 eV for SrTiO₃) and successfully described both the valence band and mid-gap states (Figures S6 and S7). The HSE06 functional was therefore used for all calculations presented in this paper.

The PAW method^{5,6} was used for effective atomic potentials with the cutoff energy of the plane wave basis set at 400 eV. After a comparison of $2 \times 2 \times 2$ and $4 \times 4 \times 4$ Monkhorst-Pack k-point meshes,⁷ the $2 \times 2 \times 2$ mesh was adopted. Gaussian smearing was applied with a width of 0.2 eV. The lattice constant was set to the experimental value of 3.905 Å. Structure optimization was done using the HSE06 functional until the maximum force became less than 0.03 eV/Å. While the Ti⁴⁺-O distance in a SrTiO₃ crystal was 1.95 Å, Ir⁴⁺-O and Ir³⁺-O distances were 1.98 Å and 2.01 Å, respectively. The DOS data were

calculated based on these optimized structures. The $\text{Ir}^{3+}:\text{SrTiO}_3$ system was modeled by injecting an excess electron with the same amount of uniform background counter charge. Molecular energy level diagrams for the Ir^{4+} and Ir^{3+} dopants are shown in Figure S8.

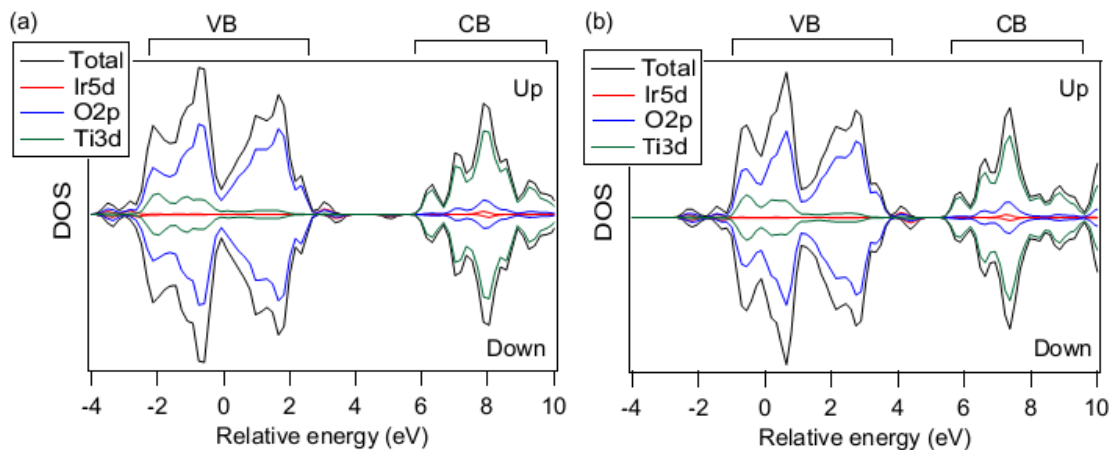


Figure S6. Total DOS (black) and PDOS of Ir 5d (red), O 2p (blue), and Ti 3d (green) for $\text{Ir}^{4+}:(3.7 \text{ at\%}): \text{SrTiO}_3$ obtained by first-principles calculations using (a) the HSE06 functional and (b) the GGA/PBE96 functional. Up and down spin states are shown in the upper and lower parts of the figures, respectively. Up and down spin states are distinguished in this system due to the existence of an unpaired electron at the Ir^{4+} site.

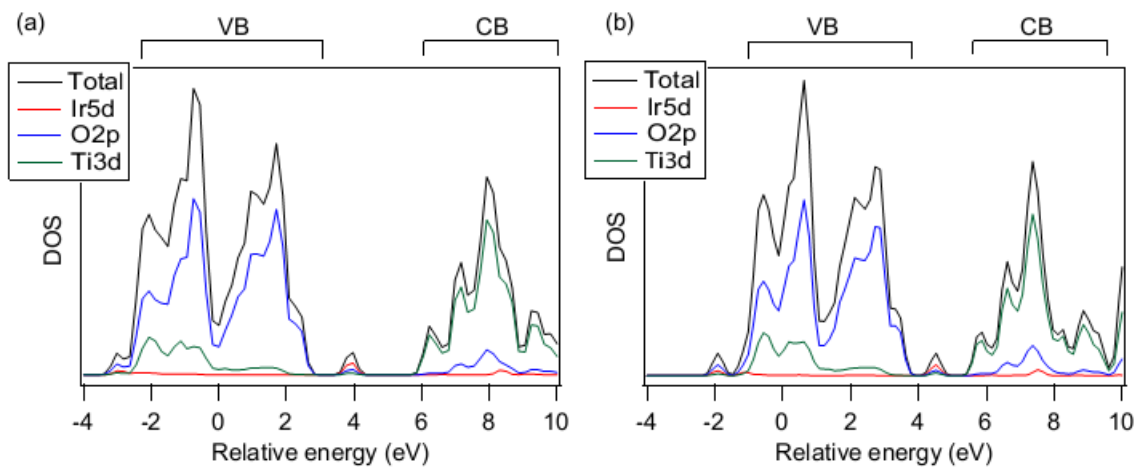


Figure S7. Total DOS (black) and PDOS of Ir 5d (red), O 2p (blue), and Ti 3d (green) for $\text{Ir}^{3+}:(3.7 \text{ at\%}): \text{SrTiO}_3$ obtained by first-principles calculations using (a) the HSE06 functional

and (b) the GGA/PBE96 functional. Up and down spin states are not distinguished in this system, since both up and down spin states are located at the same energy positions due to the absence of any magnetic anisotropy in this system.

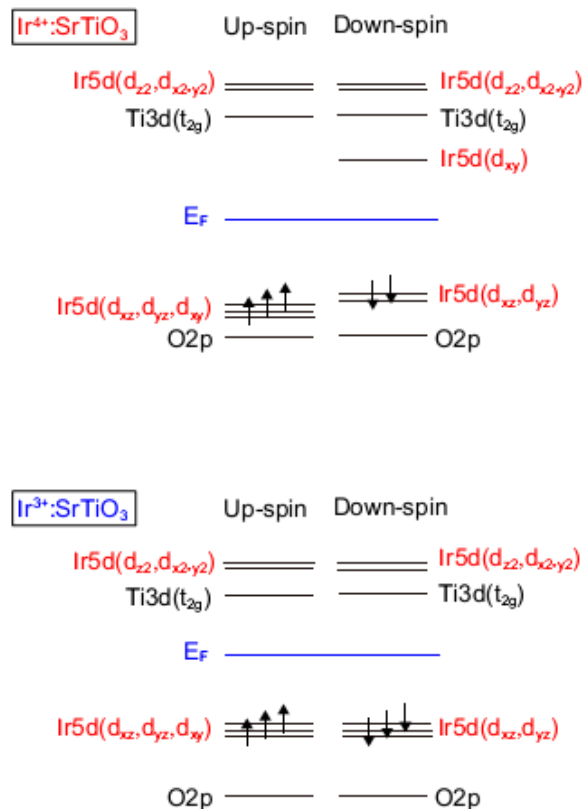


Figure S8. Molecular energy diagram for the Ir^{4+} and Ir^{3+} sites.

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