## Supporting Information for:

## Novel porous polymorphs of zinc cyanide with rich thermal and mechanical behavior

Fabien Trousselet,<sup>1</sup> Anne Boutin,<sup>1, \*</sup> and François-Xavier Coudert<sup>2,†</sup>

<sup>1</sup>École Normale Supérieure, PSL Research University, Département de Chimie, Sorbonne Universités – UPMC Univ Paris 06, CNRS UMR 8640 PASTEUR, 24 rue Lhomond, 75005 Paris, France <sup>2</sup>PSL Research University, Chimie ParisTech – CNRS, Institut de Recherche de Chimie Paris, 75005 Paris, France

## **Complement to Section V.A**

The distributions of radial distances  $g(r_{Zn-C/N})$ , shown in Fig. S5, allow to distinguish different types of transitions. For either dense or porous structures, the first two peaks (labeled 1 and 2 on the Figure, at  $\simeq$  2 and  $\simeq$  3.1 Å respectively) are associated to the cyanide ions coordinating the zinc, and show very little evolution under pressure. For the *dia-c* structure, a third (very broad) peak at distances  $r \le 5.0$  Å is also visible; it corresponds to the distances between a given Zn, in one of the interpenetrated *dia* frameworks, and the nearest cyanide ions from the other framework. With increasing pressure, this peak shifts significantly to shorter distances (from 4.9 at zero pressure to 4.4 Å at P = 2.0 GPa) and simultaneously broadens. This shift is reminiscent of the experimental observation by Collings et al. that, in the high-pressure phase of  $Zn(CN)_2$ , the corresponding distances (shown in Fig. 9 of Ref.1) decrease with increasing pressure, with a tendency towards 6-fold coordination of Zn at very large pressures. Our approach fails to describe reliably the high-pressure phase, but we suspect that both features are indeed connected and propose the following mechanism: an increase in pressure tends to amplify the *buckling vibrations* (or  $\langle \theta^2 \rangle$ ) of a cyanide ion around its Zn–Zn axis. While such vibrations have little impact on the positions of peaks 1 and 2 in  $g_{Zn-C/N}(r)$  (at order 2 only in  $\theta$ ), they impact much more the distances corresponding to the 3rd peak (at order 1 in  $\theta$ ), accounting for its broadening and shift towards smaller distances. The latter shift is favored by (a combined effect of van der Waals interactions and) compression. The distributions of Zn-C/N radial distances r' evaluated on time-averaged atomic positions — see Fig. S<sub>5</sub>(c) where the third peak seems is converted for P > 1.2 GPa into a very broad feature — indicate that this transition is accompanied by a substantial increase of disorder in the relative atomic positions of distinct frameworks.

This effect should also happen in the *mok* phase where  $g_r(Zn-C/N)$  presents also a broad peak, of similar nature, around 5 Å, but not in the porous phases where such a peak is absent (see Fig. S5-(a)). This may explain the fact that pressure-induced transitions are of different nature in porous structures (at lower pressures, and with a clear discontinuity in volume). The suspected driving mechanism for these transitions is rather a mode softening, that can be evidenced by the evolution of elastic constants.

<sup>\*</sup> Email: anne.boutin@ens.fr

<sup>&</sup>lt;sup>†</sup> Email: fx.coudert@chimie-paristech.fr. Twitter: @fxcoudert. Web: http://coudert.name/

<sup>&</sup>lt;sup>1</sup> I. Collings et al., J. Am. Chem. Soc. 135, 7610(2013).

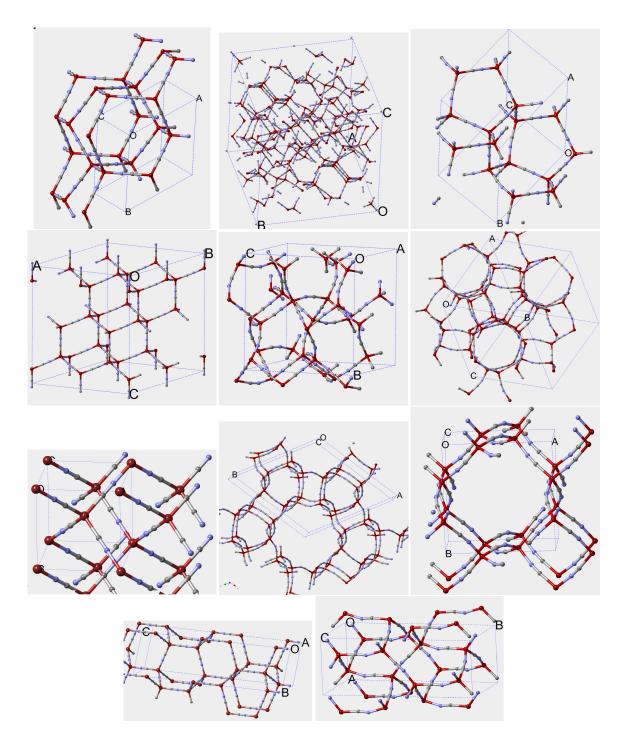


Figure S1. The eleven T = oK structures of  $Zn(CN)_2$  studied in this work, as obtained from energy-minimzation. Unit cells are indicated, and a bit more than 1 unit cell is shown to better give account of the structure's topology. Color code: Zn in red, C in grey, N in blue. Top line, from left to right: *lon, che, unj*; Second line: *dia, gsi, lcs*; Third line: *dia-c, una, unc*; Last line: *cfc, mok*.

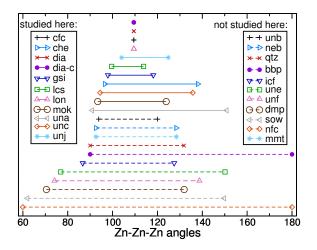


Figure S2. Minimal and maximal values of Zn–Zn–Zn angles in (i) the nets considered in this study (11 uppermost points/continuous lines) and (ii) other 4-coordinated nets found in the RCSR database, but not considered in the study.

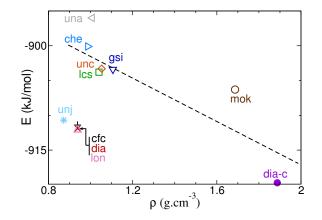


Figure S<sub>3</sub>. Energy of zinc cyanide polymorphs (per  $Zn(CN)_2$  unit) versus density. For each polymorph, the structure was obtained by re-optimizing, with the MD force field (and near-zero Kelvin dynamics), the corresponding structure found by DFT.

Framework	$\alpha_a$	$\alpha_b$	$\alpha_c$	$\alpha_V$
cfc	-12.62	-12.74	-12.58	-37.94
che	-13.50	-13.54	-13.37	-40.41
dia	.,	-13.15		· · ·
dia-c		-12.18		
-				
gsi	-	-13.31		
lcs	-14.90	-14.90	-14.82	-44.62

Table S1. Table giving, for each  $Zn(CN)_2$  framework listed, the linear thermal expansion coefficients  $\alpha_i = \frac{1}{l_i} \left( \frac{dI_i}{dT} \right)$ , with  $l_i = a, b, c$  the cell parameters; additionally the volume thermal expansion coefficient  $\alpha_V = \frac{1}{V} \left( \frac{dV}{dT} \right)$  is also given. Values are in MK<sup>-1</sup>. Note that, apart from the *mok* case, for a given framework the 3 linear expansion coefficients differ from each other by 5 percent at most.

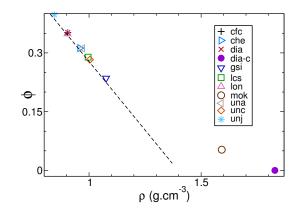


Figure S4. Porosity  $\phi$  as a function of density  $\rho$  for the structures listed, at T = 0 K and P = 0.

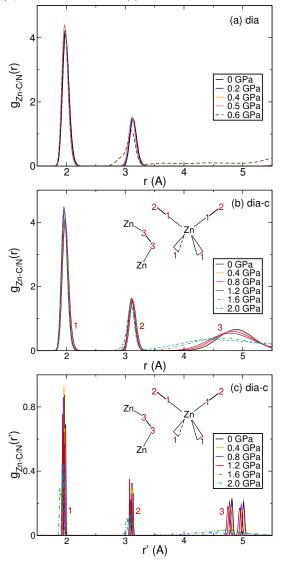


Figure S5. Distributions of Zn-X distances r (in ångströms) between a Zn atom and other types of atoms (X=C or N), for the *dia* (a) and *dia-c* (b) structures, at T = 300 K and pressures P indicated in the legend. For the *dia-c*, peaks refer to the type of C/N atoms, as indicated in the cartoon (1: nearest atom of a cyanide coordinating the Zn atom; 2: other atom of such a cyanide; 3: atoms of a further neighbor cyanide, not coordinating this Zn).(c): distributions of distances r' of the same type as in (b), but evaluated on time-averaged structures.

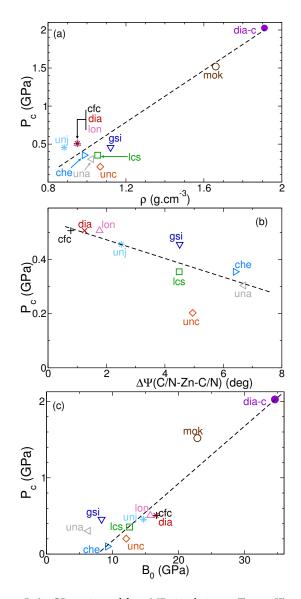


Figure S6. Plots showing the critical pressure  $P_c$  (in GPa, estimated from MD simulations at T= 300 K) versus (a) the framework density (estimated in MD at the same temperature and zero pressure); (b) the root mean square deviation  $\Delta \Psi$  of C/N–Zn–C/N angles (in degrees, estimated from *T* = 0 K structures); or (c) the bulk modulus  $B_0$  at (in GPa, estimated from MD simulations at T= 300 K).

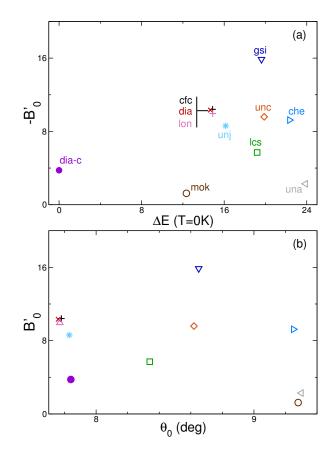


Figure S7. Coefficients  $B'_{0}$  estimated from MD simulations at T = 300K and subsequent fits to Eq. (4) in the paper, for the 11 structures indicated; they are plotted against (a) the DFT-obtained energy  $\Delta E(T = 0K)$  (in kJ/mol); and (b) the angle  $\theta_{0}$  (in degrees).

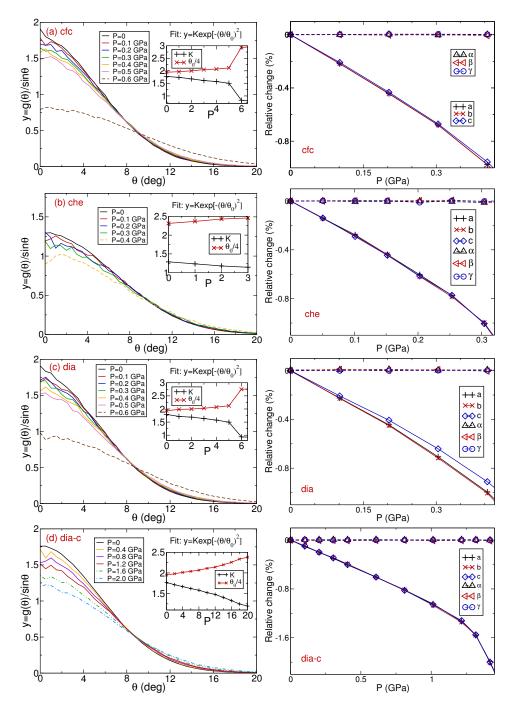


Figure S8. Evolution of structures (a) *cfc*, (b) *che*, (c) *dia*, and (d) *dia-c* under pressure, at T = 300 K (in MD simulations of stepwise compression). Left: functions  $g(\theta)/\sin(\theta)$  with  $g(\theta)$  the distribution of angle  $\theta$  (see inset in Fig. 4 of the main text) at pressures indicated in the captions. Right: Relative evolution of cell parameters *a*, *b*, *c* and cell angles  $\alpha$ ,  $\beta$ ,  $\gamma$  versus pressure.

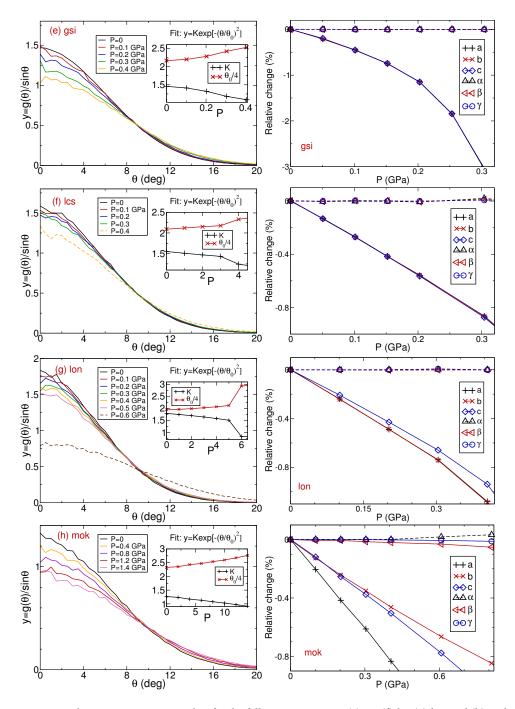


Figure S9. Similar content as in Fig. S8, but for the following structures: (e) gsi, (f) lcs, (g) lon, and (h) mok.

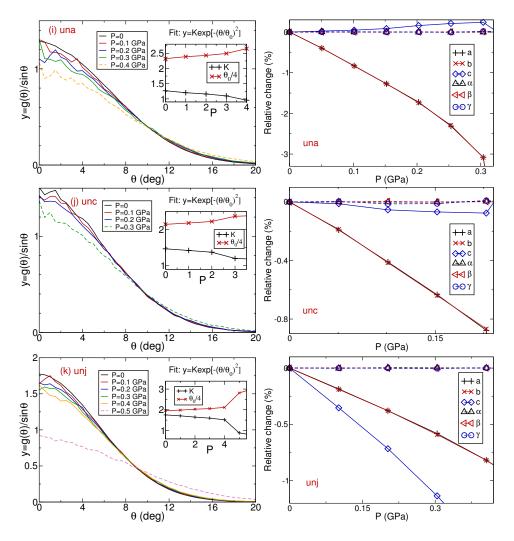


Figure S10. Similar content as in Fig. S8, but for the following structures: (i) una, (j) unc, and (k) unj.