

Nanolevitation Phenomena in Real Plane-Parallel Systems due to the Balance between Casimir and Gravity Forces

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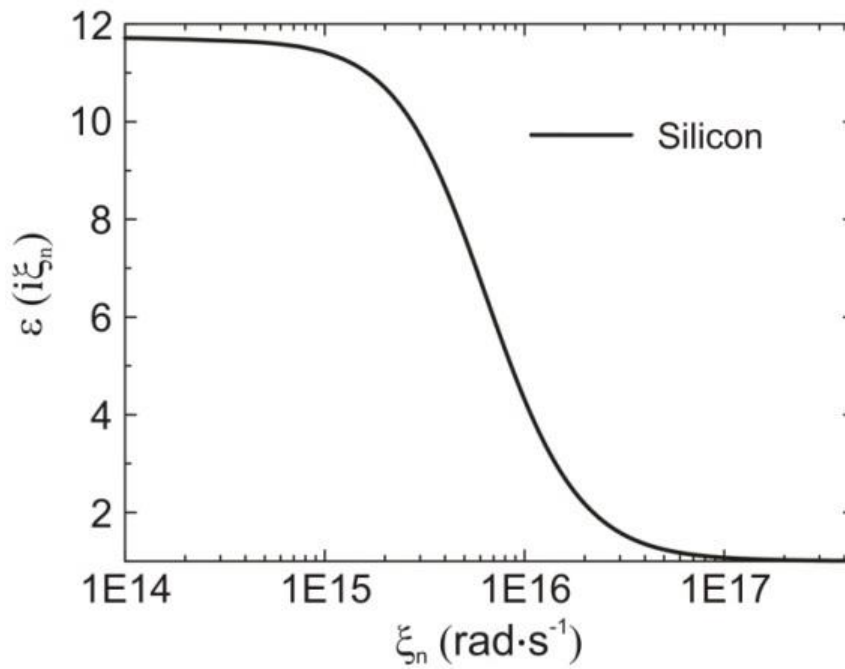


Figure S1: Dielectric permittivity evaluated at Matsubara frequency, $\epsilon(i\xi_n)$, of Si (with a doping level $1.1 \cdot 10^{15} \text{ cm}^{-3}$ and resistivity $0.077 (\Omega \cdot \text{cm})^{-1}$) for an extended frequency range.

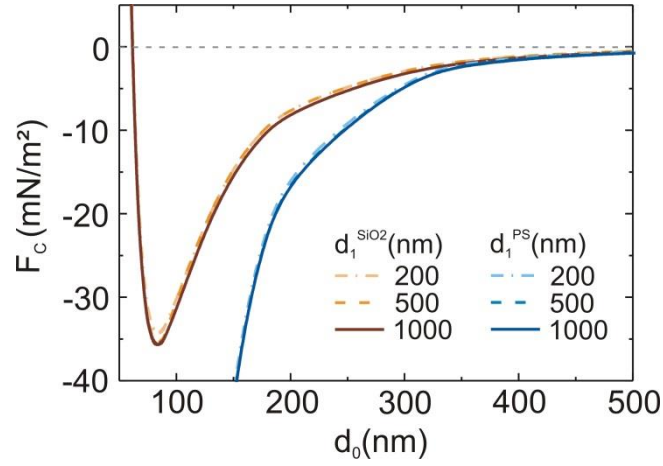


Figure S2: Casimir force (at room temperature) acting on thin films of silica (SiO_2) and polystyrene (PS) immersed in glycerol over a silicon (Si) substrate, as a function of the separation distance, d_0 . Film thicknesses of $d_1 = 200, 500, 1000$ nm are considered.

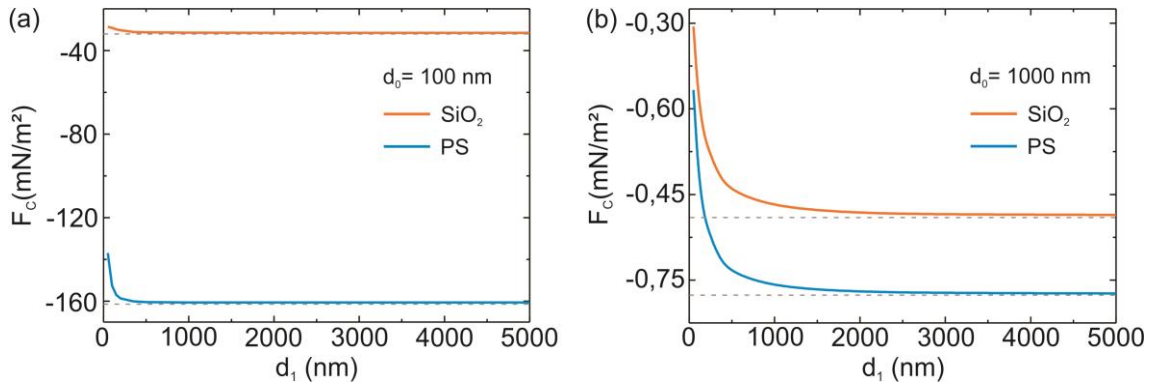


Figure S3: Casimir force as a function of thickness for a single slab of SiO_2 (orange) or PS (blue) immersed in glycerol over a silicon substrate at (a) $d_0 = 100$ nm, and (b) $d_0 = 1000$ nm. The interaction for thick enough slabs tends to the limiting case of two semi-infinite media (displayed with dashed grey lines) as d_1 increases.

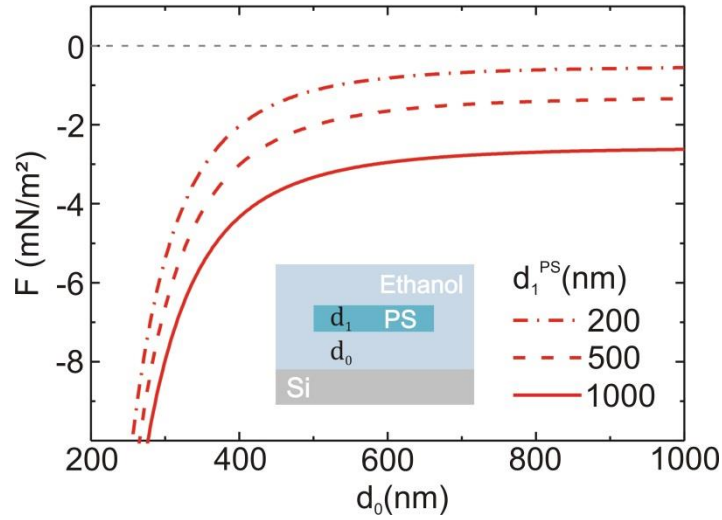


Figure S4: Total force (at room temperature) acting on thin films of PS immersed in ethanol over a Si substrate, as a function of the separation distance, d_0 . Film thicknesses of $d_1 = 200, 500, 1000$ nm are considered. $F(d_0, T) = 0$ is marked with light grey lines. In the calculations we consider, $\rho_{PS} = 1.05$ g/cm³, $\rho_{ethanol} = 0.79$ g/cm³, and the permittivity reported in Ref.[1-3], Ref.[4,5] and Ref.[6] for PS, Si, and ethanol, respectively. In all cases, an attractive (negative) force is displayed, and no equilibrium distances are found.

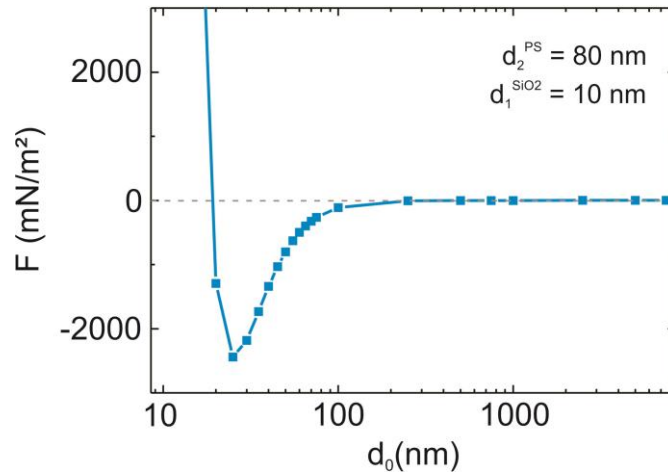


Figure S5: Total force (at room temperature) acting on a bilayer film with a bottom layer made of PS and a thin film of SiO₂ on top immersed in glycerol over a silicon substrate. The thickness of the SiO₂ ($d_2^{SiO_2}$) is fixed to 10 nm and the thickness of the PS (d_1^{PS}) to 80 nm. The figure shows how the sign of the total force changes with the separation distance (d_0). In this system we find two values of d_{eq} , ($d_{eq}^{stable} \sim 20$ nm and $d_{eq}^{unstable} \sim 1600$ nm) where the lower corresponds to a stable equilibrium position (F changes from positive to negative values) and the other one to an unstable equilibrium distance (F changes from negative to positive values).

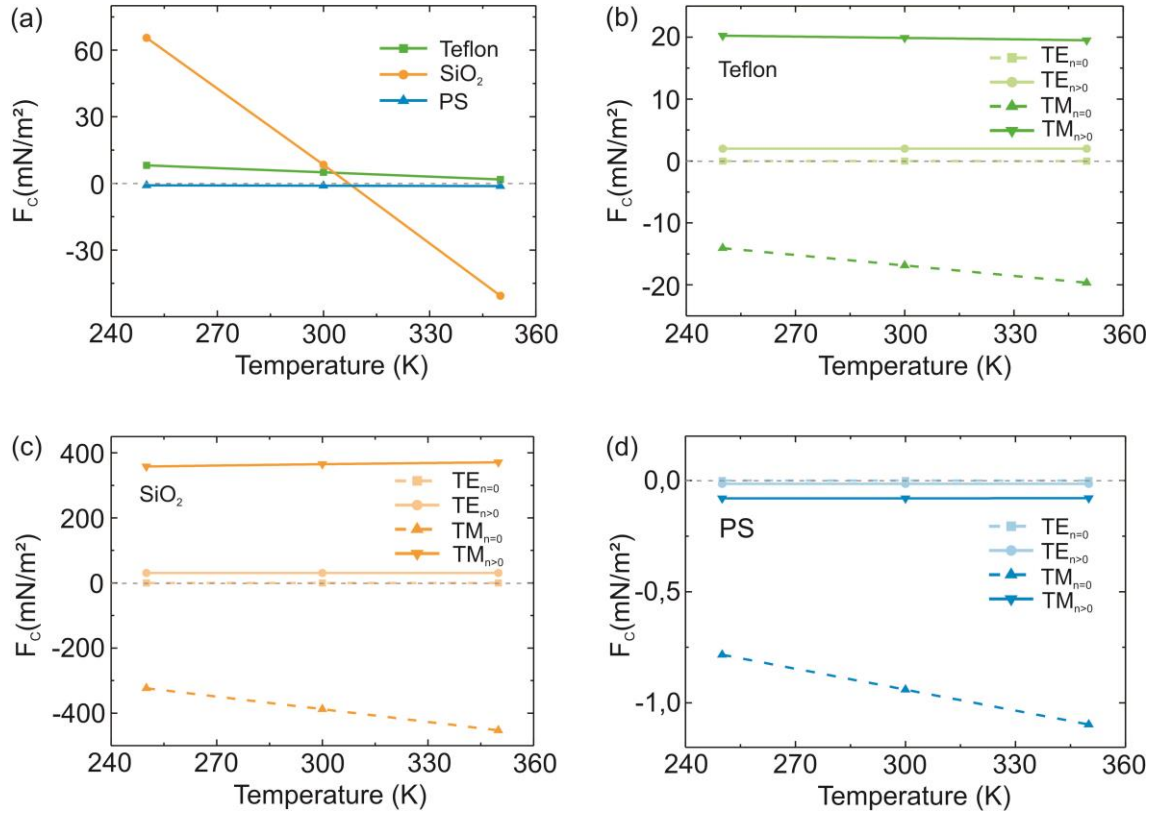


Figure S6: (a) Casimir force as a function of temperature for single layers with thickness fixed to 500 nm immersed in glycerol over a silicon substrate for Teflon, SiO₂ and PS. The gap distance is chosen to be, in each case, that of equilibrium at 300 K in Figure 4 in the main manuscript: $d_0 = 175$ nm for teflon (green), $d_0 = 60$ nm for SiO₂ (orange), and $d_0 = 450$ nm for PS (blue). Casimir force contributions of the *TE* and *TM* polarizations at $n = 0$ and $n > 0$ for (b) teflon, (c) SiO₂ and (d) PS.

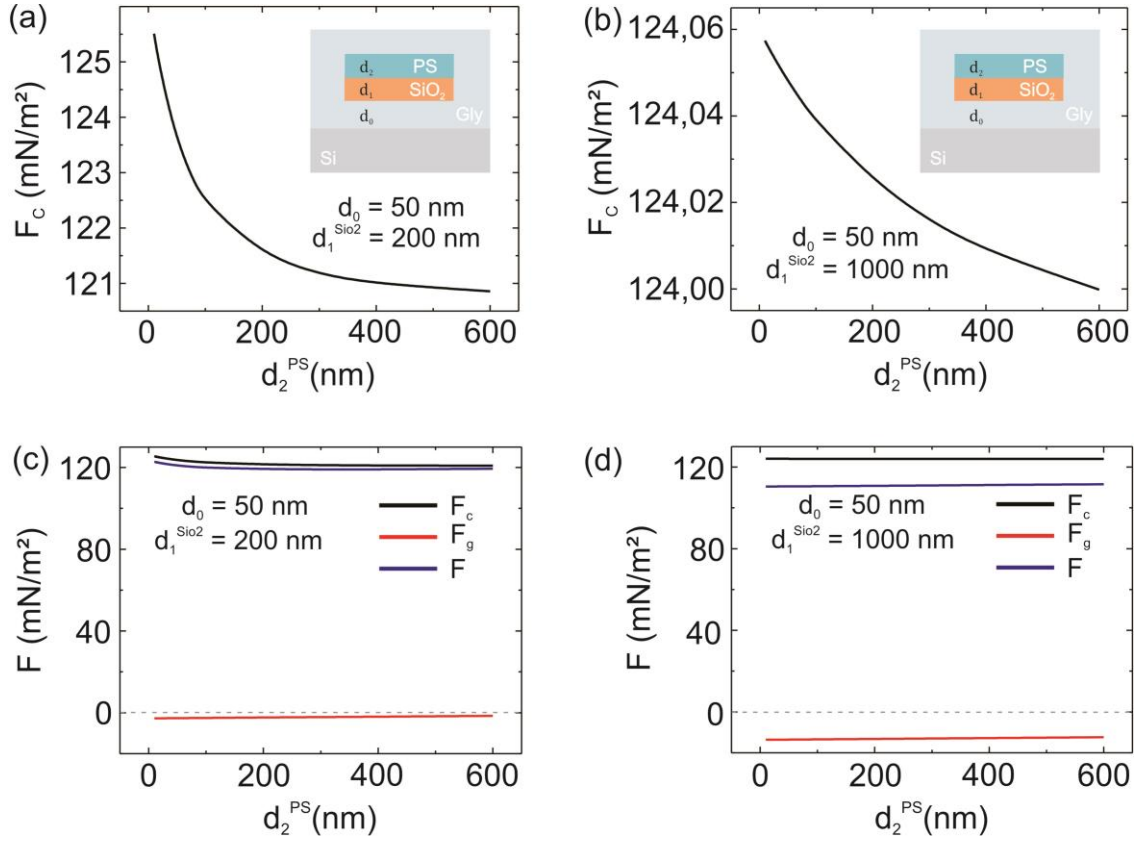


Figure S7: For a bilayer system with a lower SiO₂ layer, and a PS layer on top, immersed in glycerol over a Si substrate a distance $d_0 = 50$ nm, Casimir force as a function of PS thickness for (a) $d_1 = 200$ nm and (b) $d_1 = 1000$ nm. Panels (c) and (d) display the total force, Casimir force and gravity force for the same systems as in panels (a) and (b), respectively.

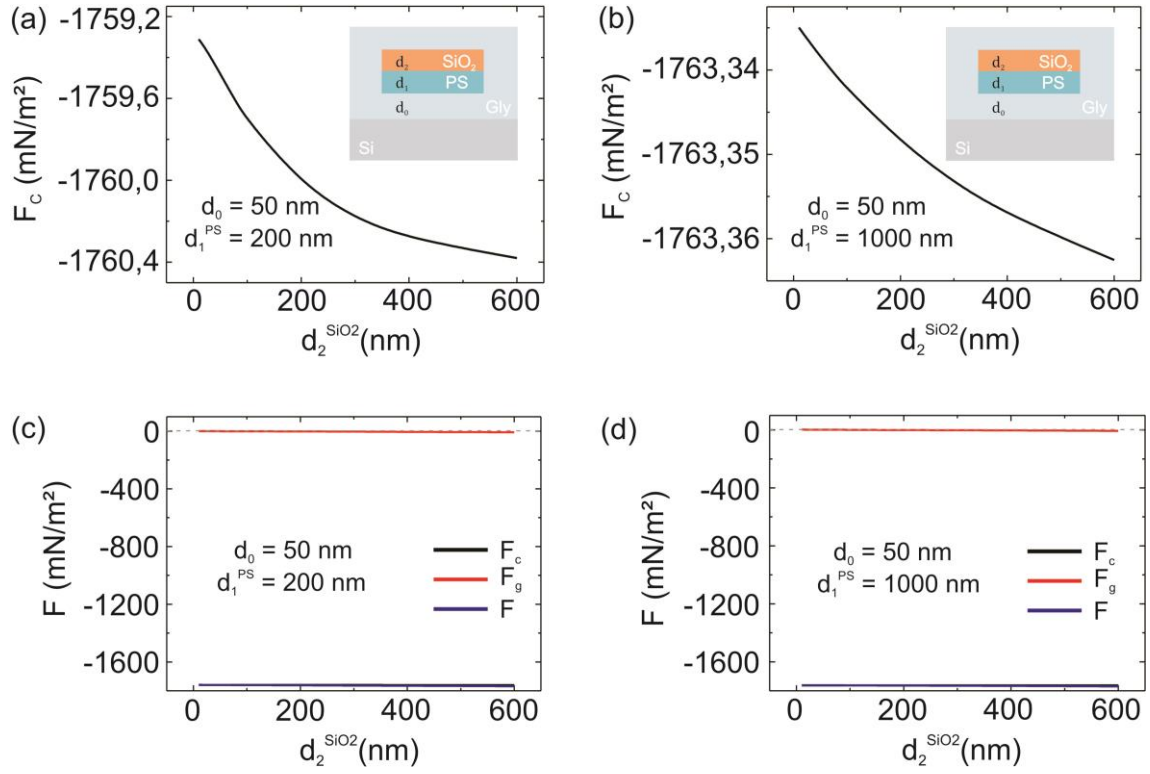


Figure S8: For a bilayer system with a lower PS layer, and a SiO₂ layer on top, immersed in glycerol over a Si substrate a distance $d_0 = 50$ nm, Casimir force as a function of PS thickness for (a) $d_1 = 200$ nm and (b) $d_1 = 1000$ nm. Panels (c) and (d) display the total force, Casimir force and gravity force for the same systems as in panels (a) and (c), respectively.

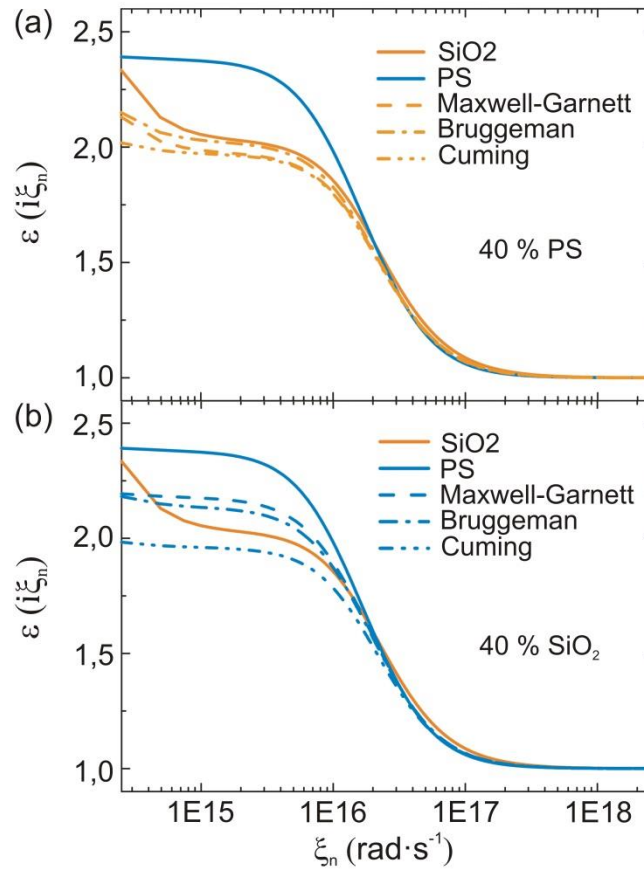


Figure S9: Permittivity evaluated at Matsubara frequencies, $\varepsilon_{eff}(i\xi_n)$, according to Maxwell-Garnett, Bruggeman, and Cuming models, for $f = 0.4$ in (a) a SiO_2 matrix with PS inclusions, and (b) the complementary system, i.e., a PS matrix with SiO_2 inclusions. For comparison, $\varepsilon(i\xi_n)$ of SiO_2 and PS are also displayed.

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