## **COMPUTATIONAL DETAILS AND FIGURES**

All calculation in this work were carried out by first-principle calculation based on the periodic density functional theory (DFT), as implemented in CASTEP code<sup>1,2</sup>. The interactions between the ions and valence electrons were described by the Vanderbilt ultrasoft pseudopotential<sup>3</sup>, using the generalized gradient approximation (GGA) proposed by Perdew, Burke, and Ernzerhof (PBE)<sup>4</sup>. The atomic configurations of Ti, C, O and Li generated from the ultrasoft pseudopotential were  $3s^23p^63d^24s^2$ ,  $2s^22p^2$ ,  $2s^22p^4$  and  $1s^22s^1$ , respectively.

The Broyden-Fletcher-Goldfarb-Shannon  $(BFGS)^5$  algorithm was adopted to minimize the total energy of whole models and relax their structure to the ground state. For geometry optimization, the final convergent values for self-consistence are set at  $2.0 \times 10^{-5}$ eV/atom for total energy, 0.05eV/Å for force, 0.1GPa for maximum stress, and 0.002Å for the maximum displacement, respectively.

A kinetic energy cut-off value of 385eV was used in all models which contain TiC, including bulk model, surface slab models,  $Li_2O_2$  cluster adsorption models, and  $Li_2O_2$ /TiC interface models. A cutoff energy of plane wave basis was set to 405eV for calculating the properties of  $Li_2O_2$  crystal and its surface. And all of cutoff energy we applied was already tested on convergence.

The k-point meshes were set at  $11 \times 11 \times 11$  for the TiC bulk unit cell, and  $11 \times 11 \times 1$  for its surface slabs using Monkhorst-Pack scheme<sup>6</sup>, which was utilized in all calculation. A 12Å vacuum is employed on both surface for TiC surface slabs, which was sufficient to ensure vanishing wave function overlap across the vacuum region. A  $7 \times 7 \times 7$  k-point mesh was set for the Li<sub>2</sub>O<sub>2</sub> bulk unit cell, and  $7 \times 7 \times 1$  for its surface slabs, including a ~10Å vacuum region on both sides, which was test to be convergent. For calculating the adsorption of Li<sub>2</sub>O<sub>2</sub> clusters, slab with 9Å vacuum region, and  $3 \times 3 \times 1$  k-point mesh was employed in  $2 \times 2$  TiC(111) supercells. Note that the same mesh size was employed by Zhang et al.<sup>7</sup> for their study O<sub>2</sub> adsorption on TiC surface. For optimizing the geometry of the interface between Li<sub>2</sub>O<sub>2</sub> and TiC surface, a thicker vacuum 14 Å and larger k-point mesh  $7 \times 7 \times 1$  were employed in

 $2 \times 2 \text{ Li}_2 O_2(0001)/\text{Ti}(111)$  supercell. To avoid polarization effect,  $\text{Li}_2 O_2$  coating was applied to both sides of a nine atomic layer symmetric TiC slab.

To correct well-known  $O_2$  overbinding problem, we follow the method described by Hummelshoj et al.<sup>8</sup> We used the gas phase reaction  $2H_2 + O_2 \rightarrow 2H_2O$ , with the DFT calculated  $H_2$  and  $H_2O$  energies and the experimental formation energy of water to establish the  $O_2$  energy.

The Mulliken population analysis was used to reveal the nature of the chemical bonds at the interface. The average bond length  $\overline{L}$  and the average overlap bond population  $\overline{n}$  are defined as follows:

$$\overline{L}(A-B) = \frac{\sum L_i N_i}{\sum N_i}$$
(4a)

$$\overline{n}(A-B) = \frac{\sum n_i N_i}{\sum N_i}$$
(4b)

$$N_i = \sum_j n(A_i - B_j) \tag{4c}$$

where  $L_i$  and  $n_i$  are the bond length and overlap bond population between atom  $A_i$ and atom  $B_j$ , respectively.  $N_i$  is the total overlap bond population of  $A_i$  contributed from all atom  $B_j$ .

The adsorption energies  $(E_{ad})$  were calculated using the following formula:

$$E_{ad} = (E_{sub-ad} - E_{sub} - E_{mol})/N_{mol}$$

where  $E_{sub-ad}$ ,  $E_{sub}$  and  $E_{mol}$  are the total energy of relaxed adsorption model, TiC substrate slab and Li<sub>2</sub>O<sub>2</sub> cluster, respectively.  $N_{mol}$  represents the number of adsorbed Li<sub>2</sub>O<sub>2</sub> clusters.

The surface energies  $(\gamma)$  were calculated using the following equation:

$$\gamma = \frac{1}{2A} (E_{slab} - N_A \mu_A - N_B \mu_B)$$
(1a)

where  $E_{slab}$  is the total energy of slab and  $E_{bulk}$  is the total energy of bulk.  $N_A$  and  $N_B$  are the number of the atom A and the atom B in the slab,  $\mu_A$  and  $\mu_B$  correspond to their chemical potential, respectively. The fraction of 1/2 accounts for the two surfaces of the slab. A stands for the surface area.

The ideal work of adhesion  $W_{ad}$  is a rough approximation to the energy required to separate the coating from the substrate at the interface, which is defined as:

$$W_{ad} = \frac{\sum_{i} E_{i} - E_{TiC/Li_{2}O_{2}}}{\Omega}$$
(5)

where  $E_i$  is the total energy of slab i (i=TiC, Li<sub>2</sub>O<sub>2</sub>).  $E_{TiC/Li2O2}$  is the total energy of the interface system.  $\Omega$  is the interface area.

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Figure s1. Side and top view of a) double  $Li_2O_2$  adsorption on C-terminated TiC{111} surface, before (left) and after relaxation (right); b) single  $Li_2O_2$  (left) and double  $Li_2O_2$  (right) adsorption on Ti-terminated TiC{111} models before relaxation.



Figure s2. a) Cubic crystallographic structure of TiC, a = 4.329 Å, b) Hexagonal crystallographic structure of  $\text{Li}_2\text{O}_2$ , a = 3.177 Å, and c = 7.751 Å, c) Planar structure of  $\text{Li}_2\text{O}_2$  cluster, O-O bond length = 1.56Å.



Figure s3. Labeling scheme for unique atom layer in the  $Li_2O_2$  structure.



Figure s4. Angular momentum projected LDOS distribution for atoms at interface and sub-interface. a) One  $Li_2O_2$  adsorption on Ti terminated Ti{111} surface, planar  $Li_2O_2$  in parallel with TiC{111}, and b) Two  $Li_2O_2$  adsorption on Ti terminated Ti{111} surface, planar Li2O2 in parallel with TiC{111}.