

# SUPPORTING INFORMATION FOR:

## Physically crosslinked polybutadiene by quadruple hydrogen bonding through side-chain incorporation of ureidopyrimidinone with branched alkyl side chains

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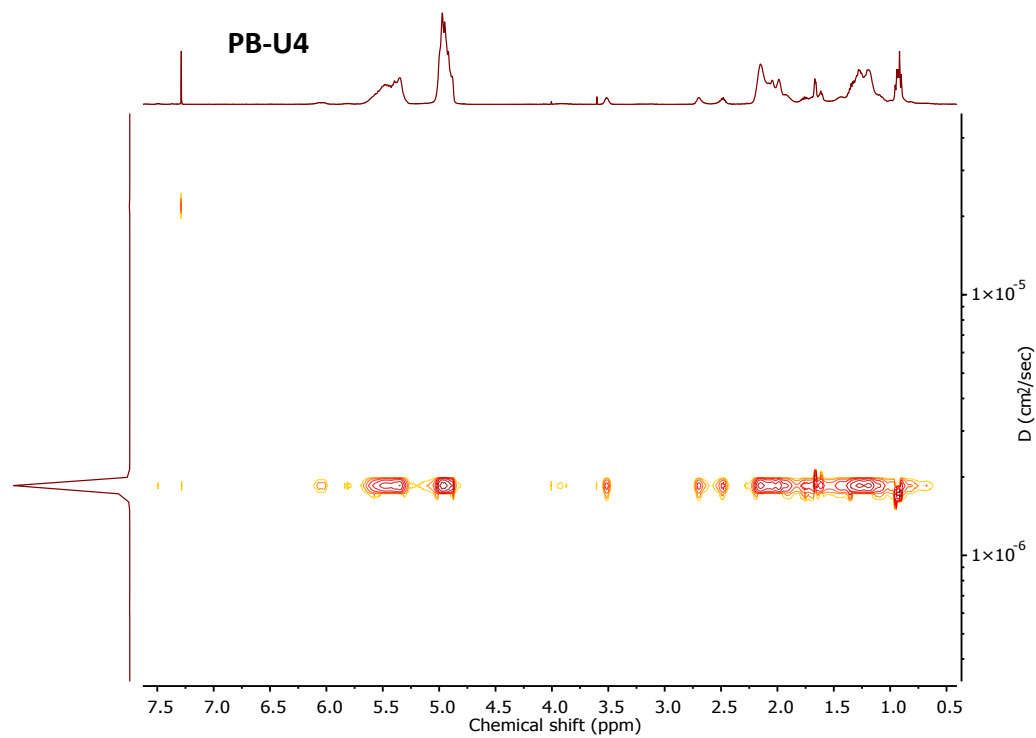
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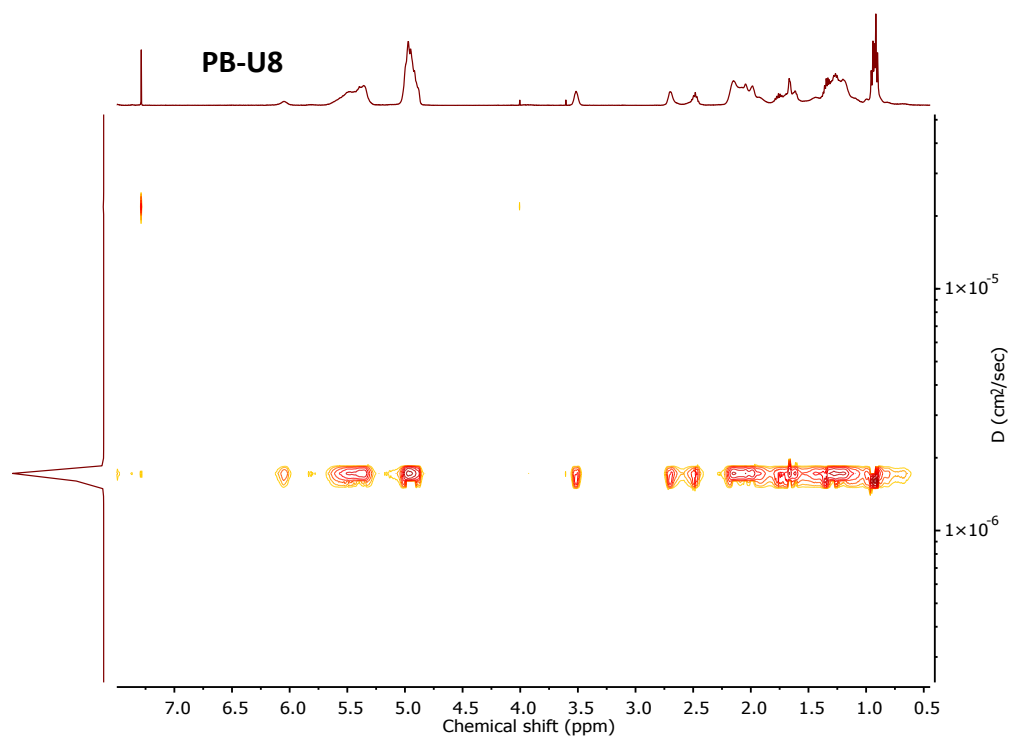
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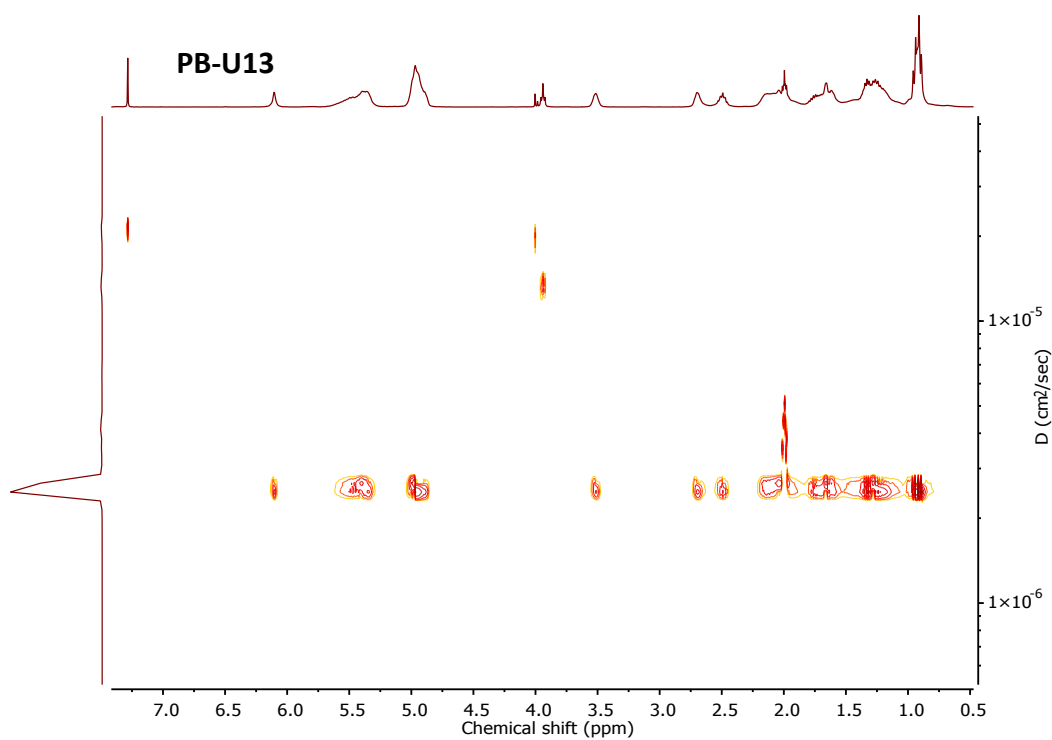
## Additional information and Characterization



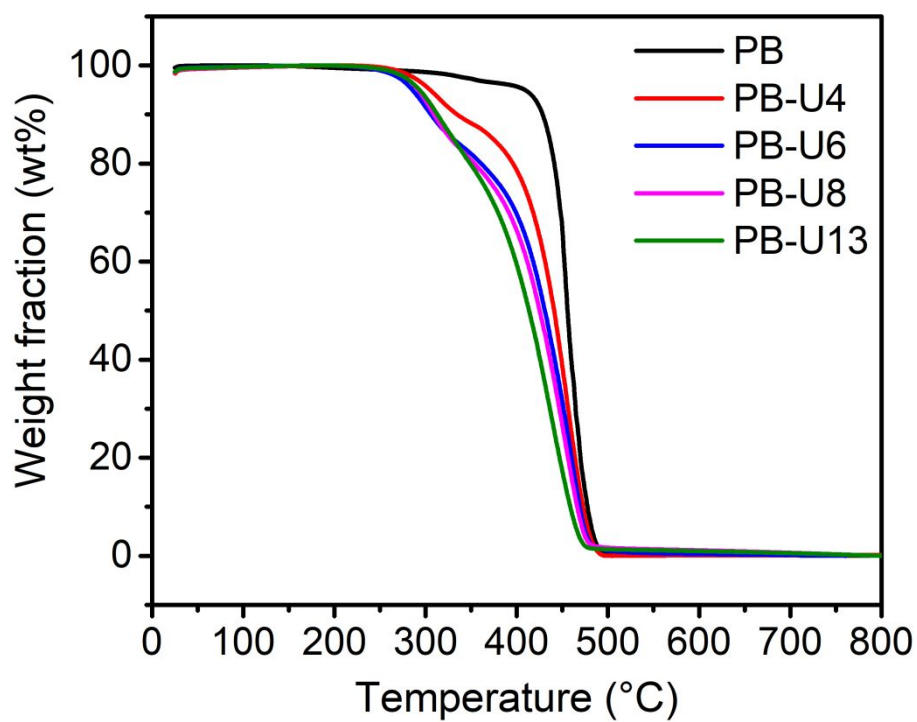
**Figure S1.** DOSY NMR spectrum of PB-U4.



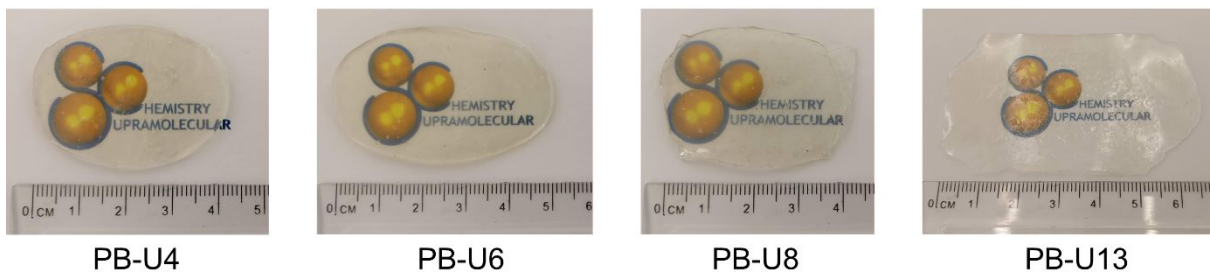
**Figure S2.** DOSY NMR spectrum of PB-U8.



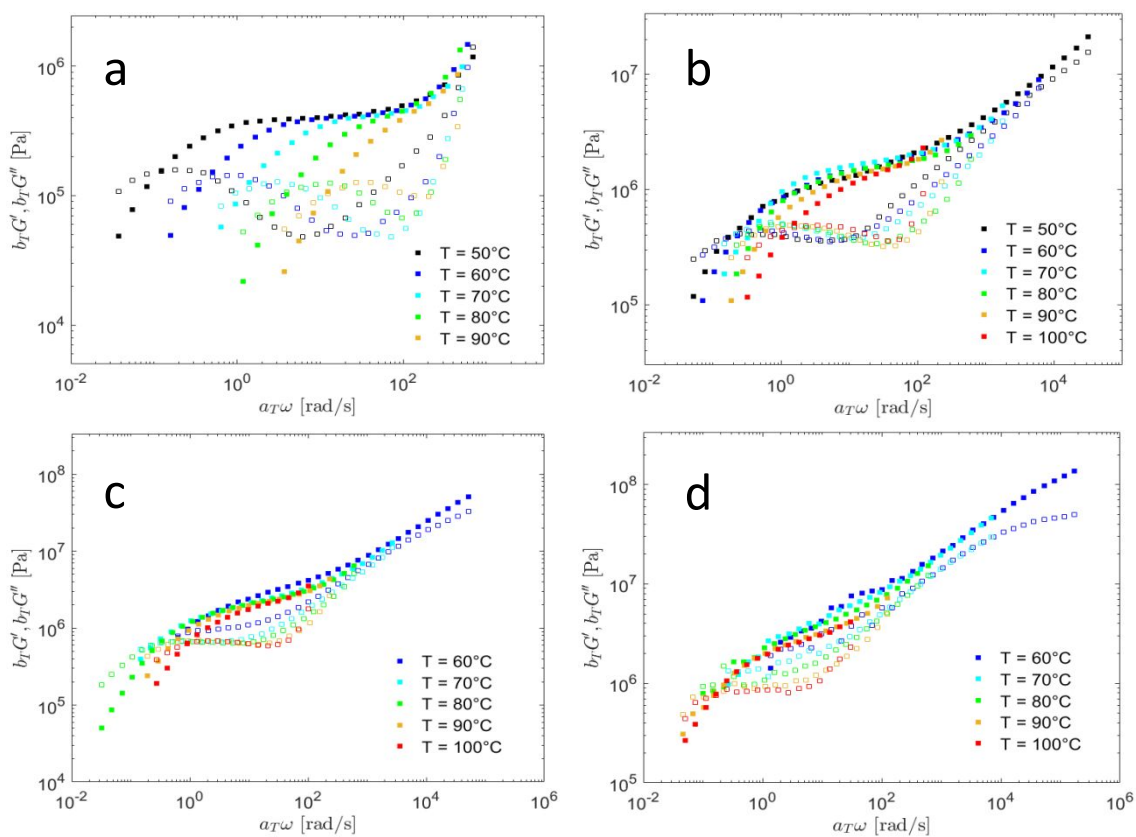
**Figure S3.** DOSY NMR spectrum of PB-U13.



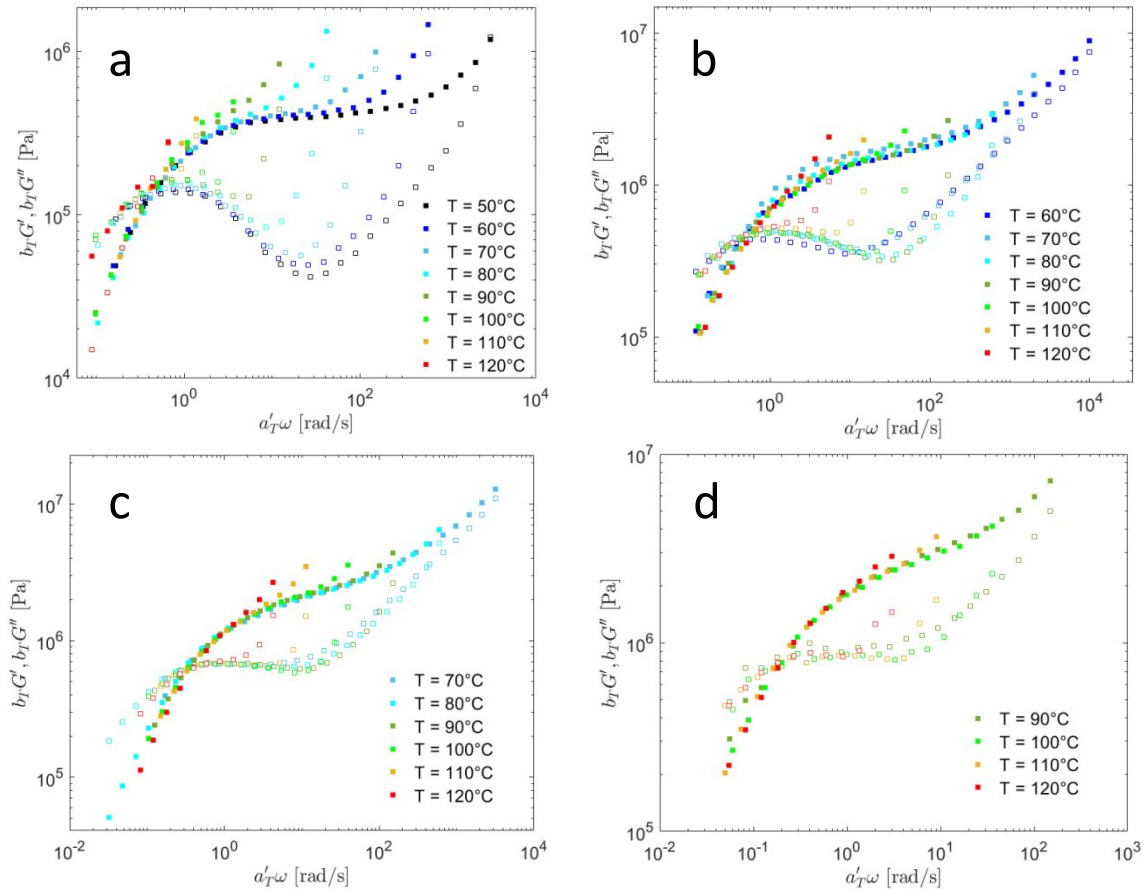
**Figure S4.** TGA curves of all PB-UPy polymers.



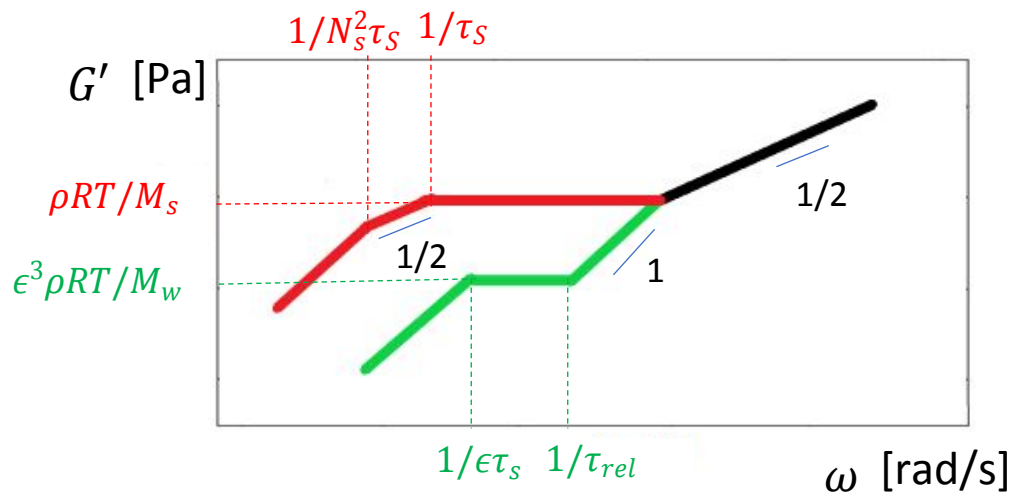
**Figure S5.** PB-UPy films obtained after compression molding.



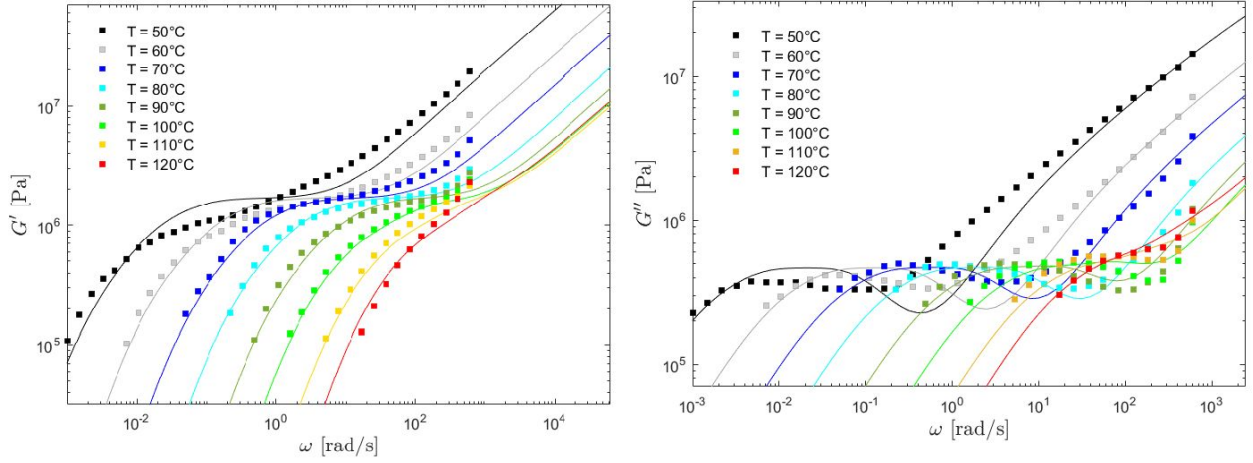
**Figure S6.** Test of time-temperature superposition of the storage and loss moduli,  $G'$  and  $G''$ , of a) PB-U4 b) PB-U6 c) PB-U8 d) PB-U13 at  $T_{ref} = 80^\circ\text{C}$ .



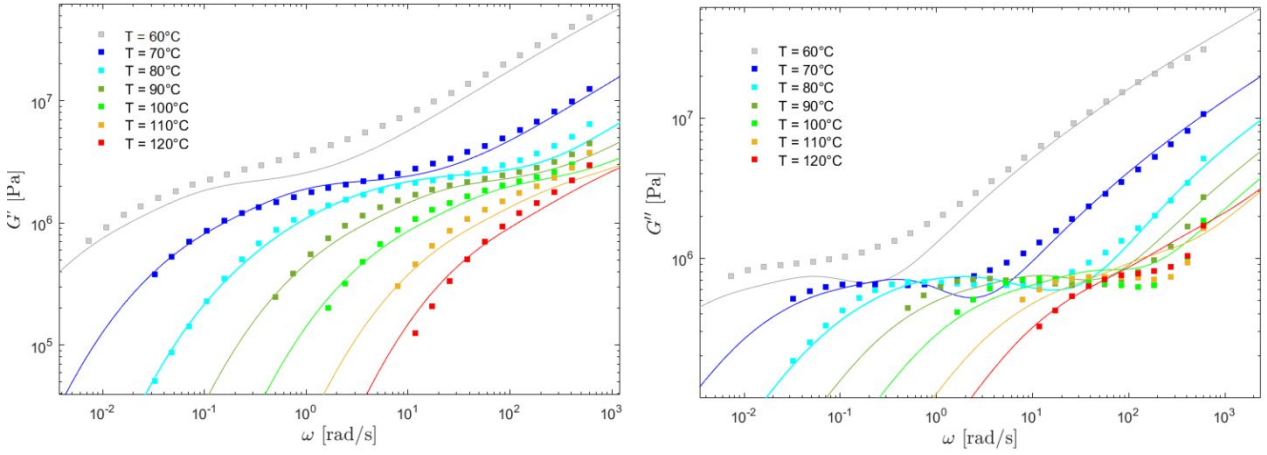
**Figure S7.** Test of time-temperature superposition of the storage and loss moduli,  $G'$  and  $G''$ , of a) PB-U4 b) PB-U6 c) PB-U8 d) PB-U13 at  $T_{ref} = 60^\circ\text{C}$ .



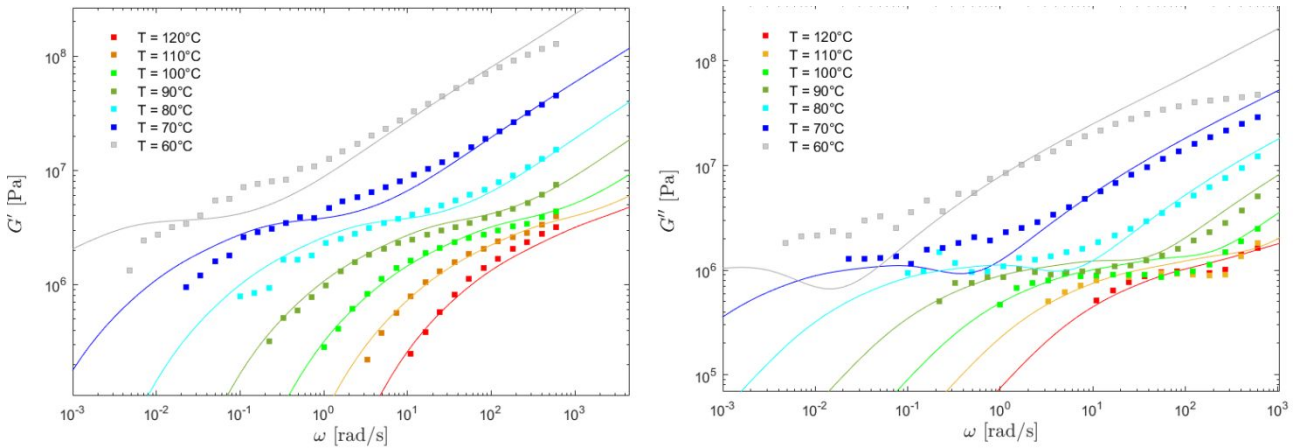
**Figure S8.** Scaling laws of the storage modulus for reversible network ( $\epsilon > 1$ ) and mixture of sol and gel ( $0 < \epsilon < 1$ ) (with mean-field percolation statistics).



**Figure S9.** Predictions of the sticky Rouse model (with the loops) for the system PB-U6 at different temperatures. The relaxation has been computed based on the value in **Table 2**, and on equation 3) and 4). To account for temperature effect, the values of  $\tau_0$  and  $\tau_s$  have been modified by the shift factors  $a_T$  and  $a'_T$  shown in **Figure 9**. The symbols are the experimental results and the solid lines are the predictions.



**Figure S10.** Predictions of the sticky Rouse model (with the loops) for the system PB-U8 at different temperatures. The relaxation has been computed based on the value in **Table 2**, and on equation 3) and 4). To account for temperature effect, the values of  $\tau_0$  and  $\tau_s$  have been modified by the shift factors  $a_T$  and  $a'_T$  shown in **Figure 9**. The symbols are the experimental results and the solid lines are the predictions.



**Figure S11.** Predictions of the sticky Rouse model (with the loops) for the system PB-U13 at different temperatures. The relaxation has been computed based on the value in **Table 2**, and on equation 3) and 4). To account for temperature effect, the values of  $\tau_0$  and  $\tau_s$  have been modified by the shift factors  $a_T$  and  $a'_T$  shown in **Figure 9**. The symbols are the experimental results and the solid lines are the predictions.

## Appendix A

Due to the different glass transition temperatures of the systems, the horizontal shift factors of the PB – Upy systems are different and have a smaller temperature dependence with decreasing the UPy concentration. Indeed, as  $T_{ref} - T_g$  increases, the activation energy decreases and so does the dependence of the relaxation times with the temperature. Despite this, as demonstrated by Wagner<sup>1</sup> and Liu *et al.*,<sup>2</sup> it is still possible to draw a mastercurve for the shift factors of all the systems. This can be done by multiplying the shift factors of the systems by a correction factor  $a_{T_g}$  which takes into account the effect of the different  $T_g$ . This shift  $a_{T_g}$  can be computed by a WLF-like equation:

$$\log_{10} a_{T_g} = \frac{-c_1^0 (T_{g,ref} - T_g)}{c_2^0 + (T_{g,ref} - T_g)}, \#(A1)$$

with the glass transition temperature  $T_g$  of the solutions and a reference glass transition temperature  $T_{g,ref}$ . The computed shift  $a_{T_g}$  are shown in Table 6. Here, PB – U13 is considered as the reference material.

**Table A1.** Correction factor  $a_{T_g}$

	PB – U13	PB – U8	PB – U6
$a_{T_g}$	1	0.0791	0.0126

## Appendix B

If a variable  $M$  is log-normally distributed, then its probability density function is given by

$$p(M = m) = \begin{cases} \frac{1}{m\sigma\sqrt{2\pi}} \exp\left(-\frac{(\ln m - \mu)^2}{2\sigma^2}\right), & m > 0 \\ 0, & \text{elsewhere} \end{cases} \#(B1)$$

The parameters  $\mu$  and  $\sigma^2$  can be expressed as a function of the mean  $E[X]$  and the variance  $Var[X]$  of the distribution in the following way.

$$\mu = \ln\left(\frac{E[X]^2}{\sqrt{E[X]^2 + Var[X]^2}}\right) \#(B2)$$

$$\sigma^2 = \ln\left(1 + \frac{Var[X]}{E[X]^2}\right) \#(B3)$$

The mean  $E[X]$  is by definition equal to  $M_n$  and the variance  $Var[X]$  can be deduced from

$$Var[X] = E[X^2] - E[X]^2 = M_w M_n - M_n^2 = M_n^2 (PDI - 1) \#(B4)$$

As a consequence, one can model the molecular weight distribution from the values of the PDI and  $M_w$ .

## References

- (1) Wagner, M. H. Scaling relations for elongational flow of polystyrene melts and concentrated solutions of polystyrene in oligomeric styrene. *Rheol. Acta*, **2014**, 53, 765–777
- (2) Liu, C.-Y., He, J., Keunings, R. & Bailly, C. New Linearized Relation for the Universal Viscosity–Temperature Behavior of Polymer Melts. *Macromolecules*, **2006**, 39, 8867–8869