

## SUPPORTING INFORMATION

# Twisted Light-Enhanced Photovoltaic Effect

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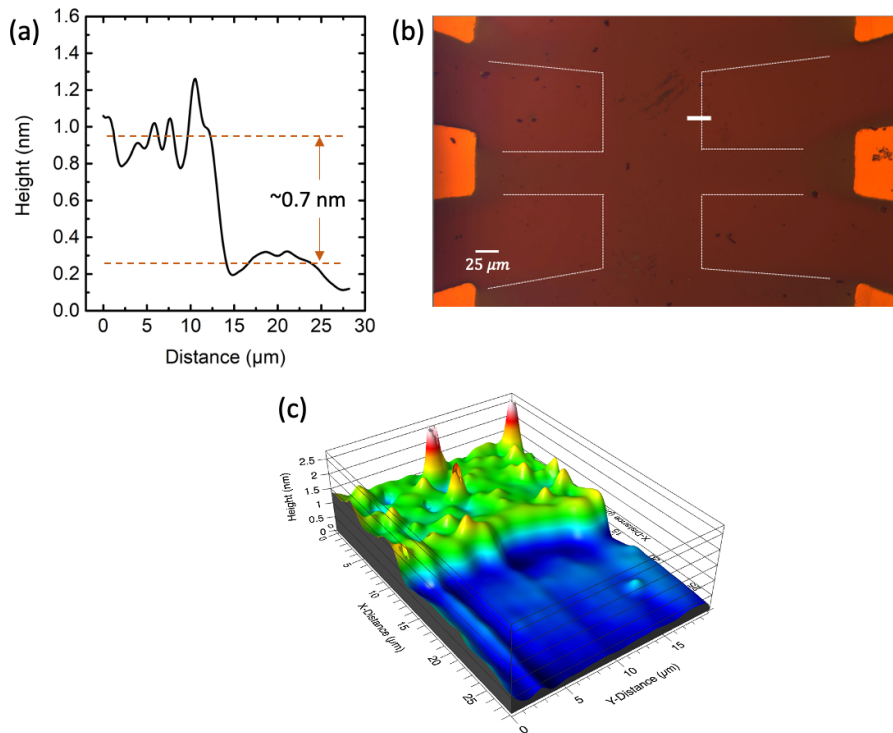
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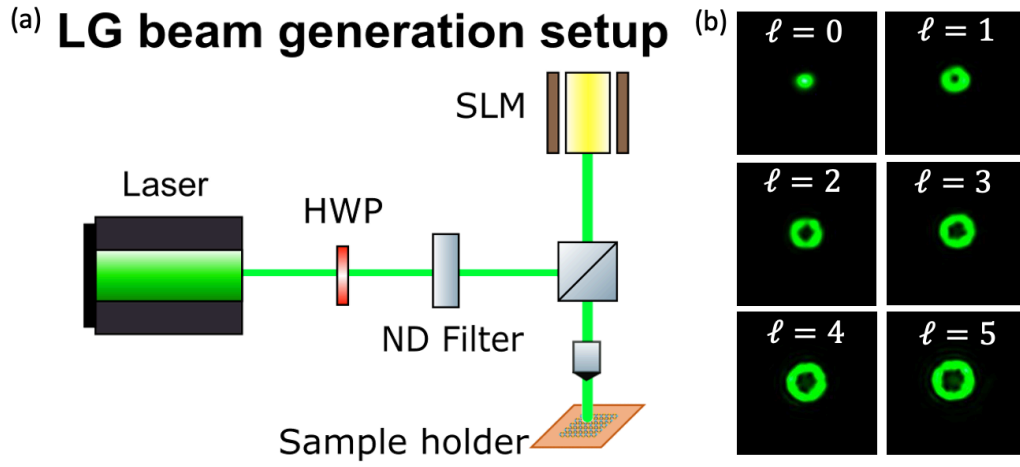
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### Atomic force microscopy measurement on the monolayer MoS<sub>2</sub> channel



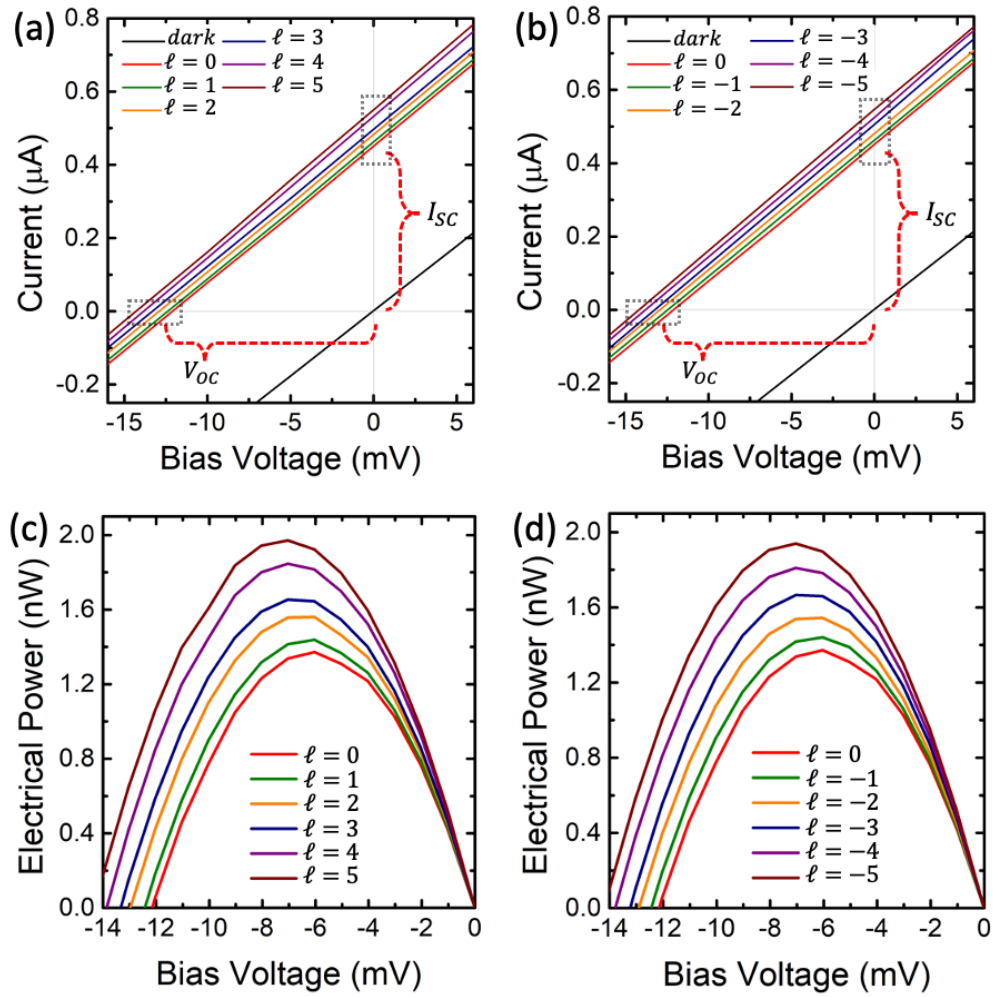
**Figure S1. Atomic Force Microscopy measurement on the monolayer MoS<sub>2</sub> channel.** (a) The height profile of the region crossing the border between the MoS<sub>2</sub> film and the SiO<sub>2</sub> substrate as indicated by a white bar in (b). The white bar has a length of ~25 μm. (c) The 3D image of the AFM profile in the vicinity of the same sample portion indicated in (b) [3D image processed using ProfilimOnline, KLA Corporation, 2021].

## Optical setup: Laguerre Gaussian beam generation



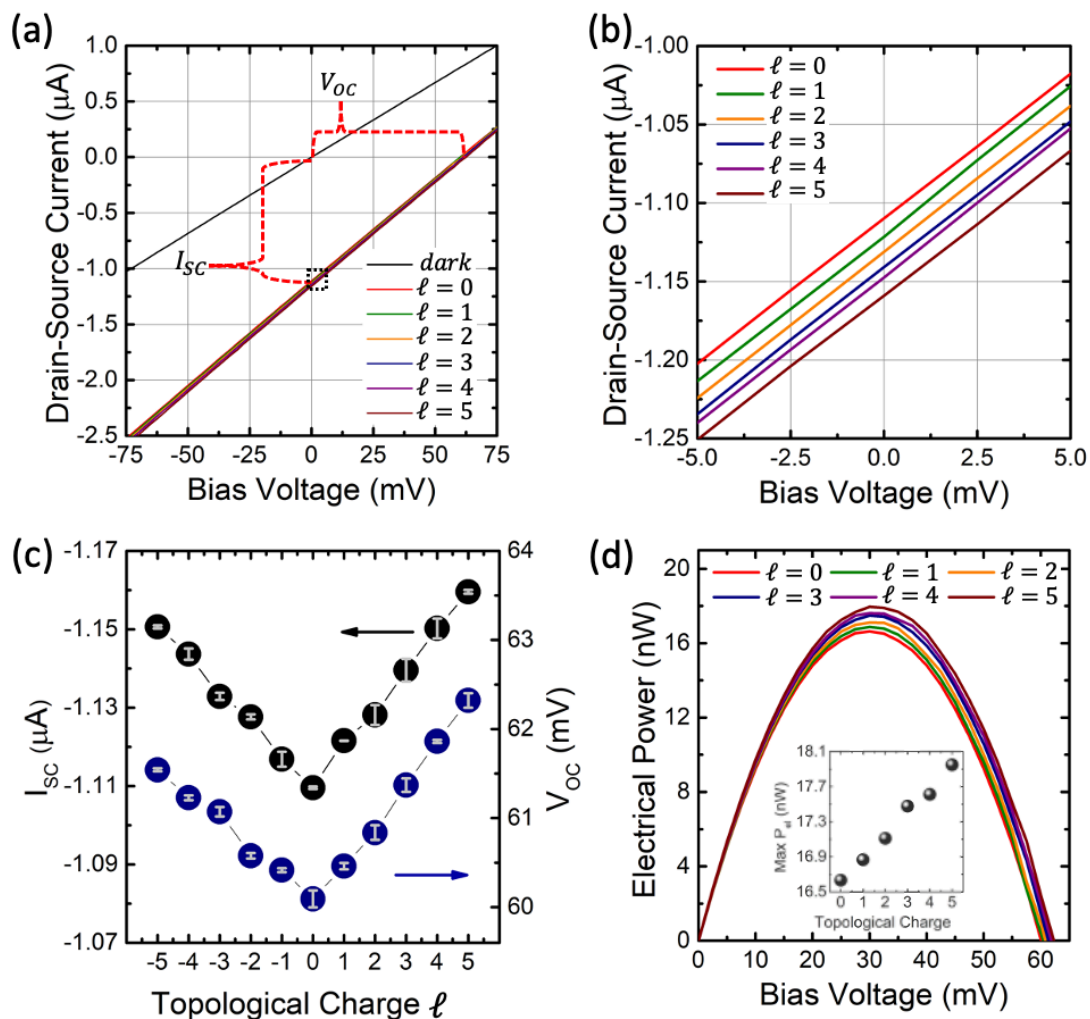
**Figure S2. Laguerre-Gaussian beam generation setup.** (a) The general optical setup for the generation of Laguerre Gaussian (LG) beam. The critical component in converting a Gaussian beam to LG light is the Spatial Light Modulator (SLM), which modulates the phase of the incident beam using computer holograms. HWP and ND Filter refers to the half-wave plate and the neutral density filter, respectively. (b) The resulting LG beam spots are presented according to selected values of  $\ell$  with a constant radial index equal to zero.

**Electrical measurements on the monolayer MoS<sub>2</sub> while illuminated with light of different OAM conditions**



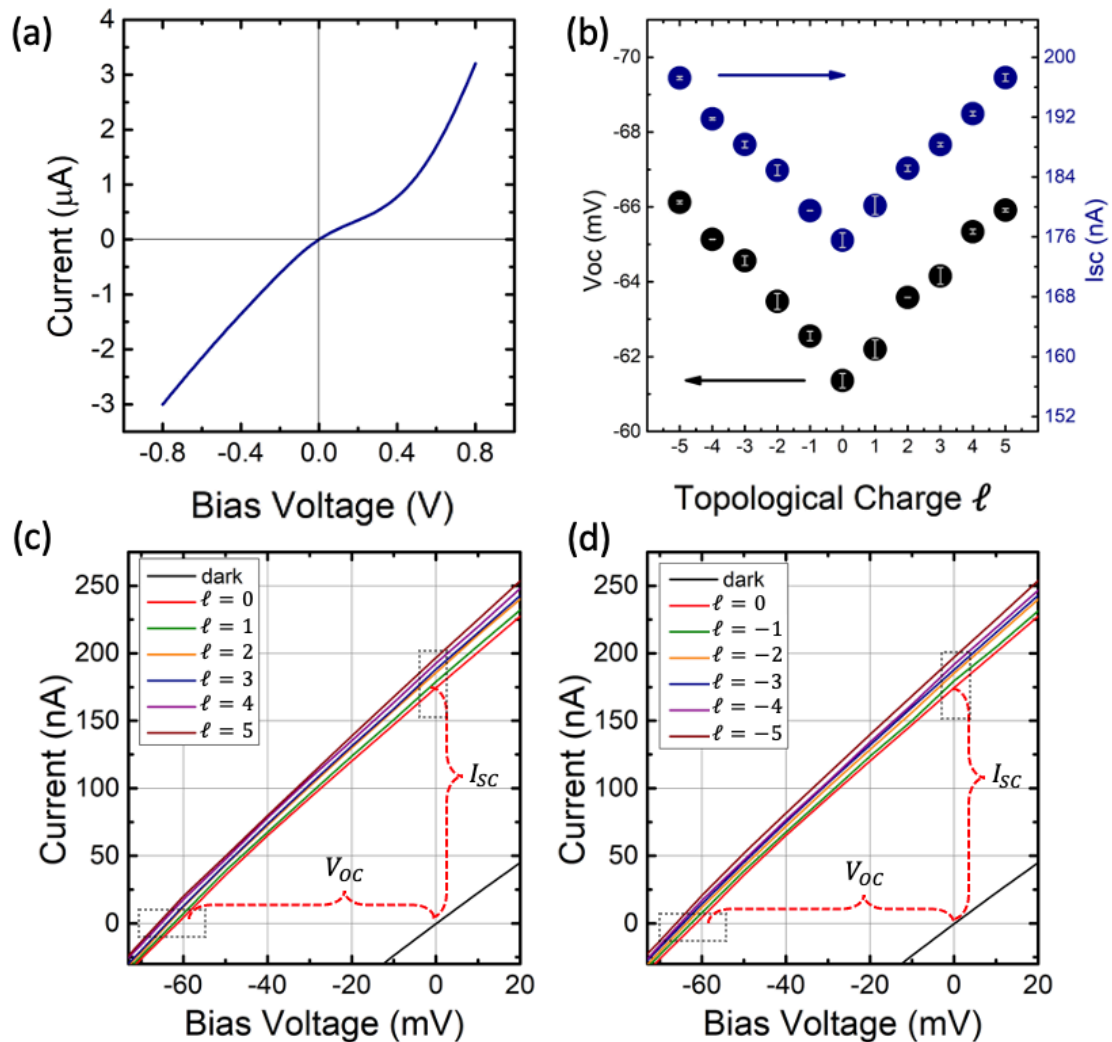
**Figure S3. IV Characteristics.** The IV characteristic curves measured while the device channel is illuminated by light with (a) positive and (b) negative values of  $\ell$ . The electrical power versus bias voltage chart for (c) positive and (d) negative  $\ell$  conditions.

Another set of electrical measurements on the monolayer MoS<sub>2</sub> channel device at fixed 100  $\mu$ W laser power



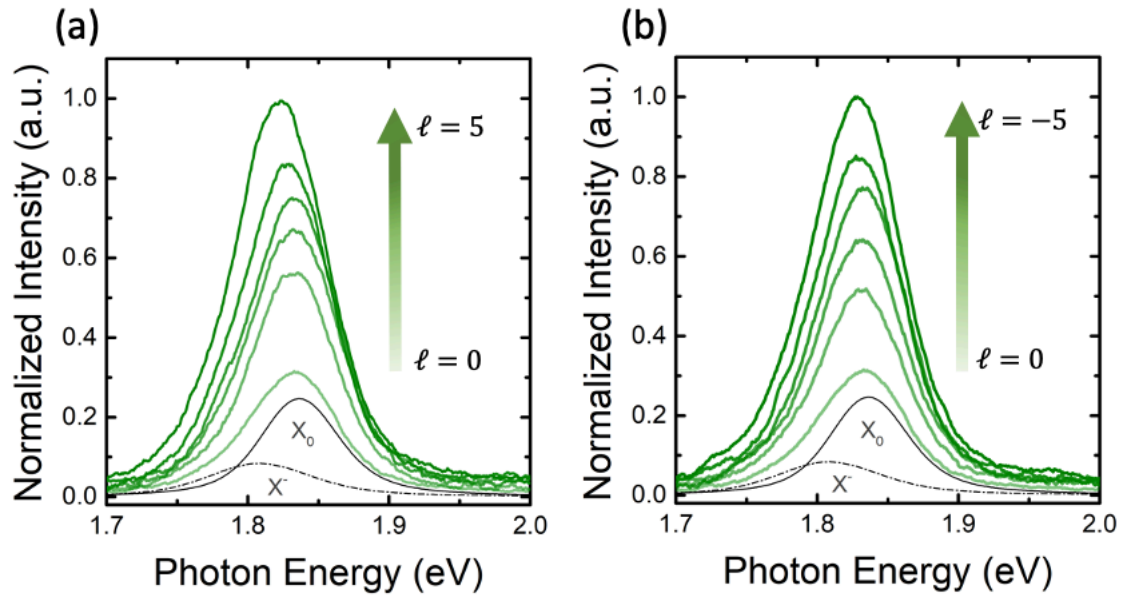
**Figure S4. Enhancement of the photovoltaic effect in monolayer MoS<sub>2</sub>.** (a) Drain-source current *versus* bias voltage characteristic of the device in the dark and under illumination of light with different  $\ell$  (4 $\times$  objective lens, 100  $\mu$ W laser power, 532 nm wavelength). (b) Zoomed-in portion of the region enclosed by the dotted black square in (a). (c) The corresponding  $I_{sc}$  and  $V_{oc}$  at various  $\ell$ , and the (d) generated electrical power *versus* bias voltage of the device extracted from (a).

**A set of electrical measurements on the second monolayer MoS<sub>2</sub> channel device at fixed 100  $\mu$ W optical power**



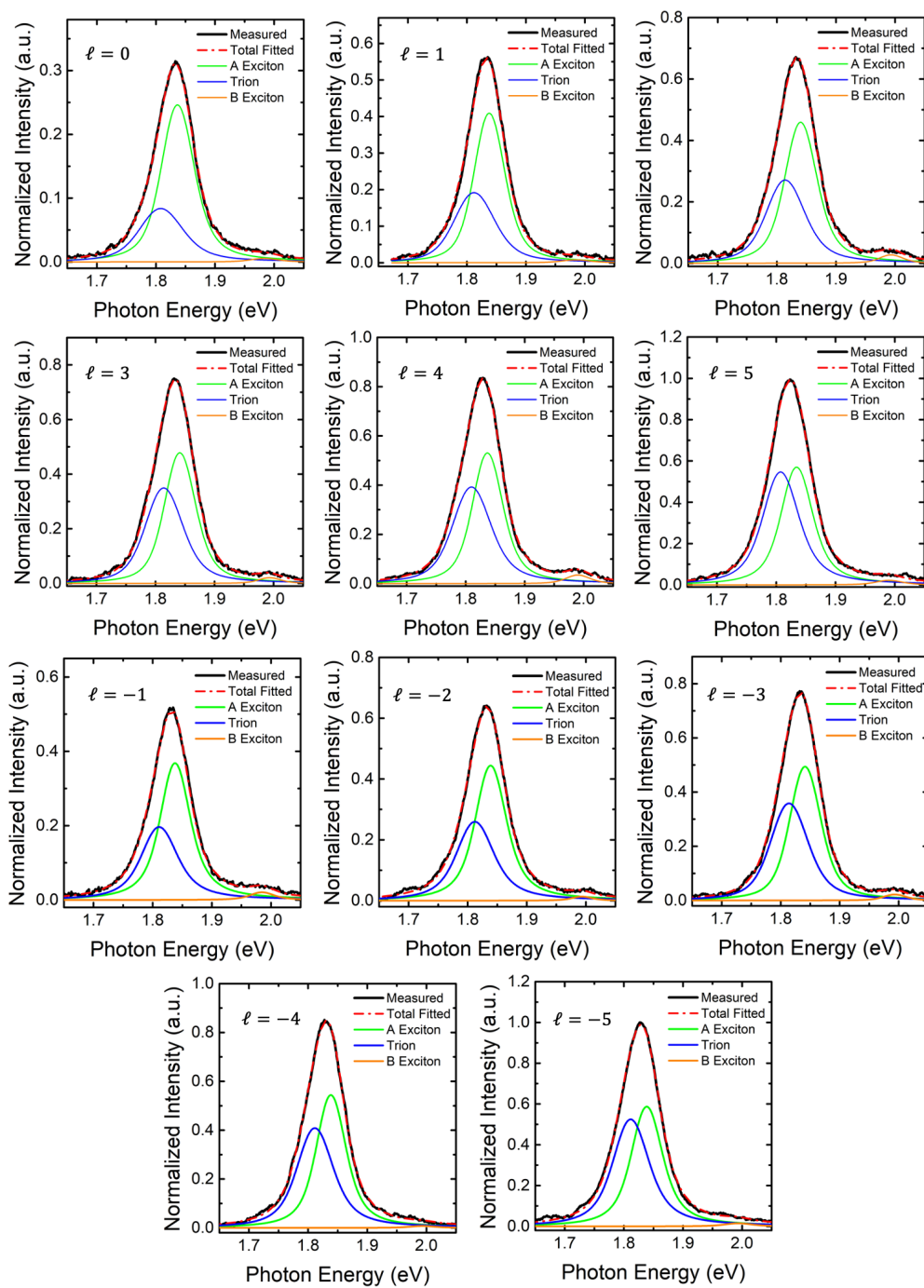
**Figure S5. IV curve measurements on another device of similar structure as the one in the manuscript. (a)** The IV curve of the device under dark condition. This second device apparently have higher channel resistance, manifested by a lower device current, compared to the device featured in the manuscript. This may be due to the presence of more defects in the device, giving it a more “diode-like” behavior. **(b)** The recorded  $I_{sc}$  and  $V_{oc}$  from this device at different light conditions. The overall increase in  $I_{sc}$  is relatively lower in this device more likely due to the presence of more defects. The error bars were derived from three measurements for each light condition. The representative IV curves measured while the device is illuminated with **(c)** positive and **(d)** negative  $\ell$  conditions (4 $\times$  objective lens, 100  $\mu$ W laser power, 532 nm wavelength).

## PL Measurements on the monolayer MoS<sub>2</sub> using incident light of various OAM



**Figure S6. Photoluminescence spectra.** The Photoluminescence spectra measured simultaneously with the IT curve in Fig. 3(a) for **(a)** positive  $\ell$  and **(b)** negative  $\ell$  conditions (all including  $\ell = 0$  or the fundamental mode). Only the fitted exciton and trion peaks for  $\ell = 0$  is shown in both (a) and (b) to tidy up the graphs.

## Detailed peak fittings of the measured PL spectra



**Figure S7. Peak-fitted PL curves.** Curve fitting results of the PL spectra in Fig. 3(b) in the main manuscript.

## Estimating the absorption efficiency

To estimate the absorption efficiency,  $\eta$ , we used the equation,

$$V_{OC} = \frac{kT}{q} \ln \left( \frac{(N_D + \Delta n + p_t)(\Delta p)}{n_i^2} \right), \quad [1]$$

that considers the intrinsic ( $n_i$ ), doping ( $N_D$ ), and photogenerated ( $\Delta n = \Delta p$ ) carriers, as well as the reported effective photoconductive gain in the form of trapped hole concentrations  $p_t$ . The estimated values of  $n_i$  and  $N_D$  are ( $\sim 10^{10} \text{ cm}^{-2}$ ) [1, 2] and ( $\sim 10^{12} \text{ cm}^{-2}$ ) [3], respectively. The formula for finding  $\Delta n$  is  $\Delta n = \phi \tau_r = \frac{\eta P_{laser} \lambda}{A_{ph} h c} \tau_r$ , where  $\eta$  is the absorption efficiency,  $\lambda$  the wavelength of the incident light (532 nm),  $A_{ph}$  the illuminated area,  $h$  Planck's constant,  $c$  the speed of light,  $\tau_r$  the carrier recombination lifetime ( $\sim 100 \text{ ps}$ ) [4], and  $\phi = (\eta P_{laser} \lambda) / (A_{ph} h c)$  the photon flux. The density of trapped hole states can be expressed as  $p_t = (\phi P_t \tau_r) / \left( \phi \tau_r + P_t \frac{\tau_t}{\tau_g} \right)$ , where  $P_t$  is the total density of hole trap states ( $\sim 10^{10} \text{ cm}^{-2}$ ), and  $\tau_t$  and  $\tau_g$  the hole trapping and release time, respectively, in which ratio  $\frac{\tau_g}{\tau_t}$  is fixed and was assumed to be  $\sim 200$  as adapted from a previous study [4]. Using all the given values, **Eq. 1** is fitted with measured  $V_{OC}$  at different values of the  $\ell$  of the incident light to get the estimated absorption efficiency,  $\eta$ .



## References:

- [1] Hong, J.; Hu, Z.; Probert, M.; Li, K.; Lv, D.; Yang, X.; Gu, L.; Mao, N.; Feng, Q.; Xie, L.; Zhang, J.; Wu, D.; Zhang, Z.; Jin, C.; Ji, W.; Zhang, X.; Yuan, J.; Zhang, Z. Exploring Atomic Defects in Molybdenum Disulphide Monolayers. *Nat. Commun.* **2015**, *6* (1), 6293.
- [2] Chae, W. H.; Cain, J. D.; Hanson, E. D.; Murthy, A. A.; Dravid, V. P. Substrate-Induced Strain and Charge Doping in CVD-Grown Monolayer MoS<sub>2</sub>. *Appl. Phys. Lett.* **2017**, *111* (14), 143106.
- [3] Wu, C.-C.; Jariwala, D.; Sangwan, V. K.; Marks, T. J.; Hersam, M. C.; Lauhon, L. J. Elucidating the Photoresponse of Ultrathin MoS<sub>2</sub> Field-Effect Transistors by Scanning Photocurrent Microscopy. *J. Phys. Chem. Lett.* **2013**, *4* (15), 2508–2513.
- [4] Furchi, M. M.; Polyushkin, D. K.; Pospischil, A.; Mueller, T. Mechanisms of Photoconductivity in Atomically Thin MoS<sub>2</sub>. *Nano Lett.* **2014**, *14* (11), 6165–6170.