

Supporting Information:

Flat Phonon Bands based Mechanism of

Amorphization of MOF-5 at *Ultra-low* Pressures

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Crystal structure of MOF-5

The starting structure of MOF-5 crystal considered for this study is obtained from Ref.^{S1} The relaxed structure obtained from our *first-principles* DFT calculations at zero pressure is provided at the end of this document.

MOF-5 exists in a face-centered cubic structure and consists of secondary building units (SBUs) positioned at the Bravais sites connected by benzene linkers along $\langle 100 \rangle$ directions (Figure S1). The *fcc* structure with $Fm\bar{3}m$ symmetry is corroborated by simulated XRD pattern showing high peak-intensity for (110), (121) and (220) peaks.

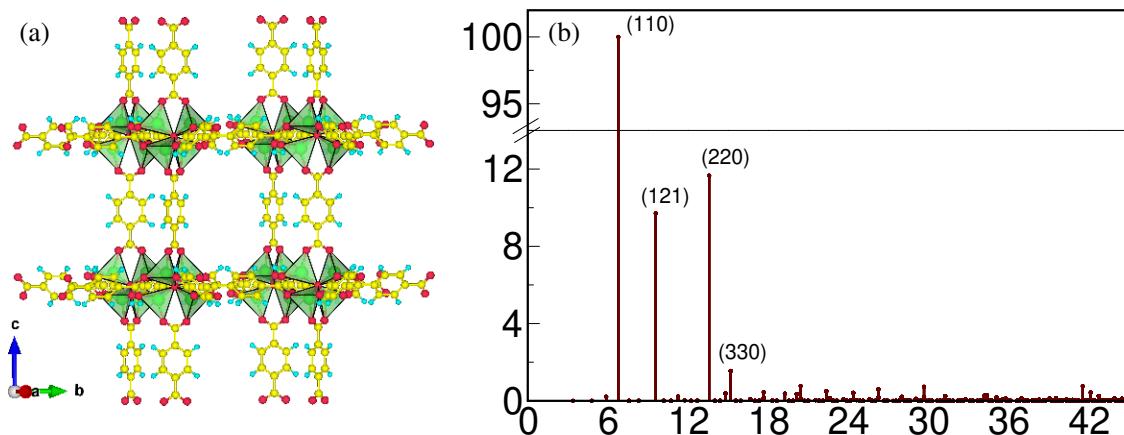


Figure S1: (a) Cubic structure of MOF-5 crystal viewed along [110] direction. Adjacent SBUs are rotated by 90° about each other, and are bonded through organic linkers (i.e. benzene rings) along $\langle 100 \rangle$. Dashed box delineates the 1,4-benzodicarboxylate or the BDC group,^{S2} and (b) simulated XRD shows high peak-intensity primarily along (110), (121), (220) planes, and the corresponding q -values agree closely with the experimental observations for single-crystal MOF-5.^{S3}

Frequencies and irreps of the low-energy normal modes of bulk and strained ($\varepsilon_h = -0.003$) MOF-5 structures, along with symmetry of the structures they generate are detailed in Table S1.

Table S1: Frequencies and symmetry labels of low-frequency phonon modes of bulk MOF-5 subjected to compressive strains; $\varepsilon_h = 0.0$ and $\varepsilon_h = -0.003$. Column III shows the subgroups of $Fm\bar{3}m$ generated by the symmetry labels (Column IV) of the normal modes of MOF-5 crystal.

$\varepsilon_h = 0.0$	$\varepsilon_h = -0.003$	Resulting structure	Irrep
ω (THz)	ω (THz)		
3.1	-1.7	P-1 (2)	T_{1g}
3.1	-1.7	P-1 (2)	
3.1	-1.7	C2/m (12)	
1.6	-1.4	C2 (5)	T_{2u}
1.6	-1.4	C2 (5)	
1.6	-1.4	Imm2 (44)	
2.9	-1.2	P-1 (2)	T_{2g}
2.9	-1.2	P-1 (2)	
2.9	-1.2	C2/m (12)	
1.5	0.4	P1 (1)	T_{1u}
1.5	0.4	Cm (8)	
1.5	0.4	P1 (1)	
0.6	1.1	$Fm\bar{3}$ (202)	A_{2g}
1.4	1.2	C2 (5)	T_{2u}
1.4	1.2	C2 (5)	
1.4	1.2	Imm2 (44)	
0.7	1.3	P-1 (2)	T_{1g}
0.7	1.3	P-1 (2)	
0.7	1.3	C2/m (12)	
0.8	1.4	Fmmm (69)	E_g
0.8	1.4	I4/mmm (139)	

Structural instability of MOF-5 crystal upon expansion

At $\varepsilon_h = 0.01$, the MOF-5 structure undergoes a lattice instability, as characterized by unstable optical modes in its vibrational spectrum (Figure S2).

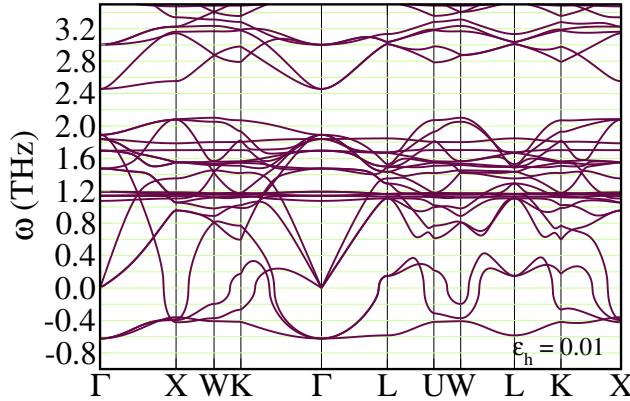


Figure S2: Low-frequency vibrational spectrum of expanded MOF-5 crystal, $\varepsilon_h = 0.01$, exhibiting unstable optical phonon bands.

Nature of normal modes of MOF-5 that destabilize the SBUs on compression

At very small compressive strains $\varepsilon_h \sim -0.003$, MOF-5 crystal undergoes lattice instabilities. The unstable phonon modes show maximum overlap with normal modes of bulk MOF-5 in the mid-frequency range, $\omega \in \{1.6-3.2 \text{ THz}\}$. The corresponding eigen-vectors show rotation and in-plane displacements of the BDC groups, as shown in Figure S3.

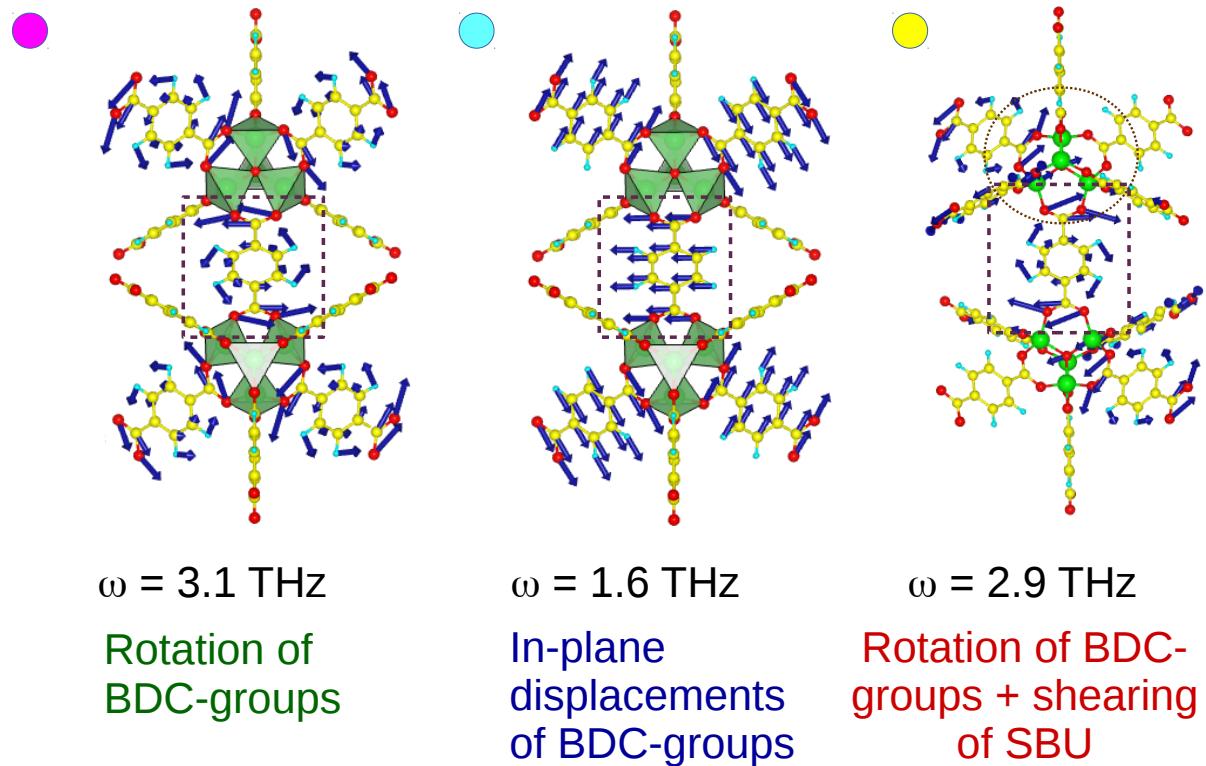


Figure S3: Eigenvectors of normal modes in frequency range $\omega \sim \{1.6\text{-}3.1 \text{ THz}\}$ that become unstable at $\varepsilon_h \sim -0.003$.

Freezing-in these normal modes in strained MOF-5 crystals shows harmonic dependence of strain-energy $E(\varepsilon_h, u_v) - E(\varepsilon_h, u_v=0)$ on amplitudes u_v for $|\varepsilon_h| < 0.3\%$ to quartic in u_v for $|\varepsilon_h| \geq 0.3\%$ (Figure S4).

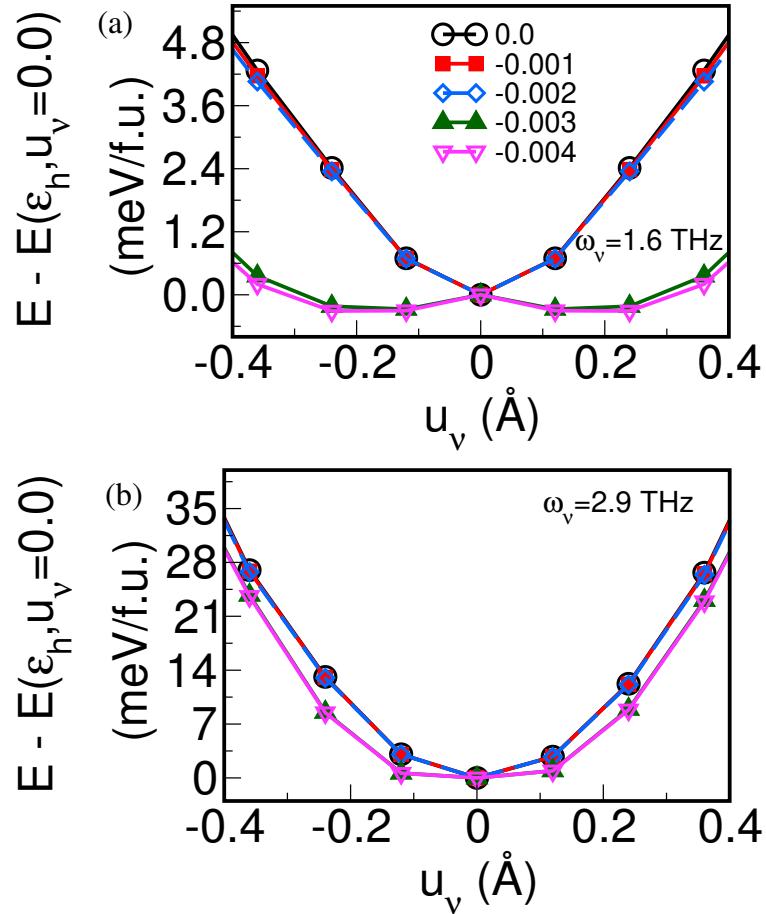


Figure S4: Dependence of internal strain energy on amplitude $|u_v|$ of triply degenerate modes at (a) $\omega = 1.6$ THz and (b) $\omega = 2.9$ THz.

Relaxation of structural distortions in compressed MOF-5

At low compressive strains, $|\varepsilon_h| < 0.01$, the structures distorted with random combinations of modes in flat unstable phonon bands relax to lower-energy states, and exhibit localized atomic displacements with respect to pristine MOF-5 crystal (Figure S5(a)). However, at moderately larger compressive strains $|\varepsilon_h| > 0.015$, relaxation of structural distortions in MOF-5 crystal give rise to sliding of nano-scaled domains relative to each other and structural inhomogeneities, leading to its irreversible amorphization (Figure S5(b)).

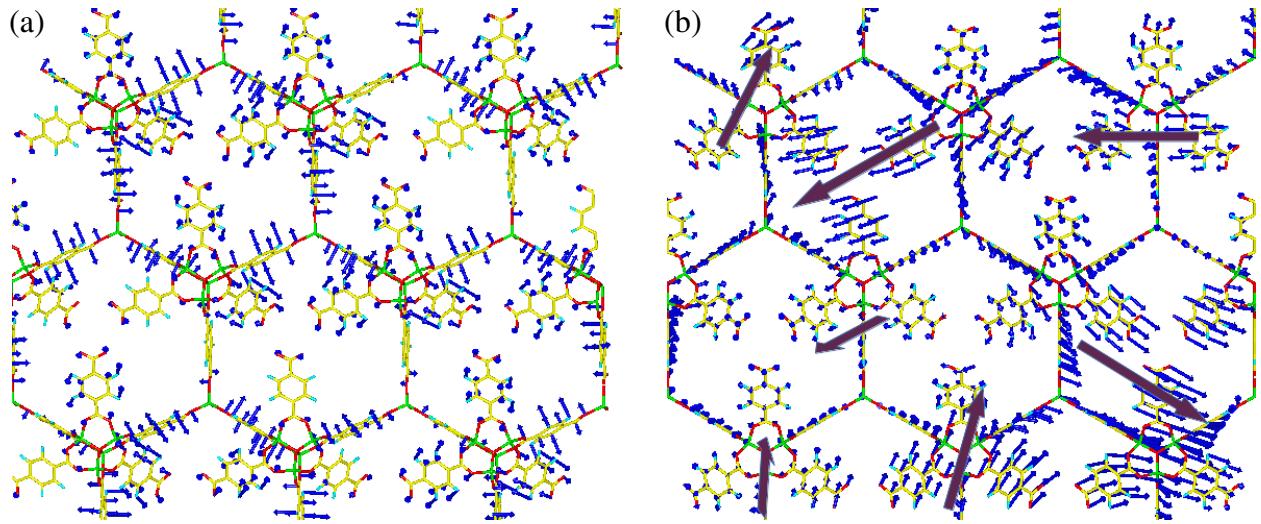


Figure S5: Atomic structural changes in distorted ($3 \times 3 \times 3$) supercells optimized to minimum-energy configuration with respect to pristine MOF-5 crystal: (a) low-strain regime, $\epsilon_h = -0.003$, showing random localized displacements, and (b) larger strains, $\epsilon_h = -0.05$, exhibiting long-ranged displacements and structural inhomogeneities.

Lattice dynamics of MOF-5 at moderately large compressive strains

At $\epsilon_h = -0.015$, the coupling of unstable optical modes with hydrostatic strain result in loss of elastic rigidity of MOF-5, as characterized by unstable transverse acoustic (TA) branches along various crystallographic directions (see Figure S6).

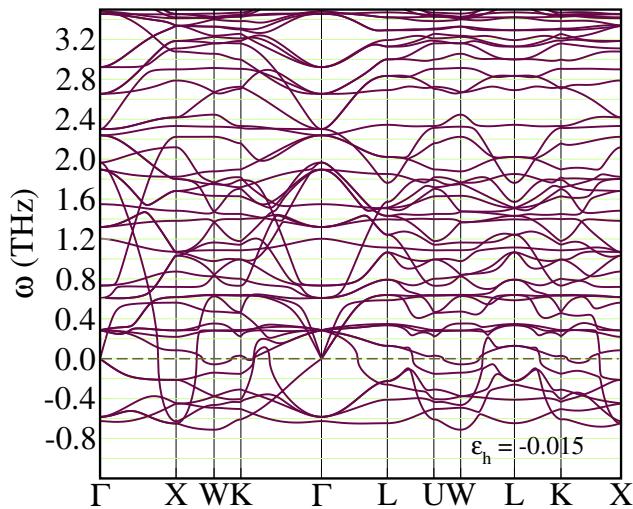


Figure S6: Low-frequency vibrational spectrum of compressed MOF-5 crystal, $\epsilon_h = -0.015$, exhibiting instability of softest transverse acoustic (TA) branches and unstable optical modes.

MOF-5 structure

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| O11 | O  | 0.36680875 | 0.36680875 | 0.19492224 | 24.0 |
| C35 | C  | 0.38912961 | 0.38912961 | 0.11087038 | 12.0 |
| C47 | C  | 0.44640641 | 0.44640641 | 0.05359358 | 12.0 |
| C59 | C  | 0.47337519 | 0.47337519 | 0.09235429 | 24.0 |
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References

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