

**Supplementary Information**

**Low-resistance p-type ohmic contacts to  
ultra-thin WSe<sub>2</sub> by using a monolayer dopant**

Abhinandan Borah,<sup>\*,†</sup> Ankur Nipane,<sup>†</sup> Min Sup Choi,<sup>‡</sup> James Hone,<sup>‡</sup> and James  
T Teherani<sup>†</sup>

<sup>†</sup>*Department of Electrical Engineering, Columbia University, New York, NY, USA*

<sup>‡</sup>*Department of Mechanical Engineering, Columbia University, New York, NY, USA*

E-mail: a.borah@columbia.edu

# Comparison of our UV-O<sub>3</sub> method with O<sub>2</sub> plasma oxidation

Table S1: O<sub>2</sub> plasma doping from literature compared to our UV-O<sub>3</sub> doping

	This work	Ref <sup>1</sup>	Ref <sup>2</sup>	Ref <sup>3</sup>	Ref <sup>4</sup>	Ref <sup>5</sup>	Ref <sup>6</sup>
Number of layers	3	8	3–50	6–15	7	2–13	>50
Monolayer dopant	yes	no	no	no	no	no	no
Process temperature (°C)	Room temp	250	N/A	Room temp	Room temp	N/A	250
$2R_C$ (k $\Omega$ . $\mu$ m)	0.6	1	N/A	N/A	4.3	N/A	N/A
Doping density (cm <sup>-2</sup> )	$4 \times 10^{13}$	$8 \times 10^{12}$	N/A	N/A	N/A	N/A	N/A
$2R_C$ improvement	300×	100×	N/A	N/A	100×	N/A	N/A
Area-selective doping	yes	no	no	no	no	no	no
Low- $R_C$ at low-temperature	yes	N/A	N/A	N/A	N/A	N/A	N/A

## AFM characterization of 2L WSe<sub>2</sub>

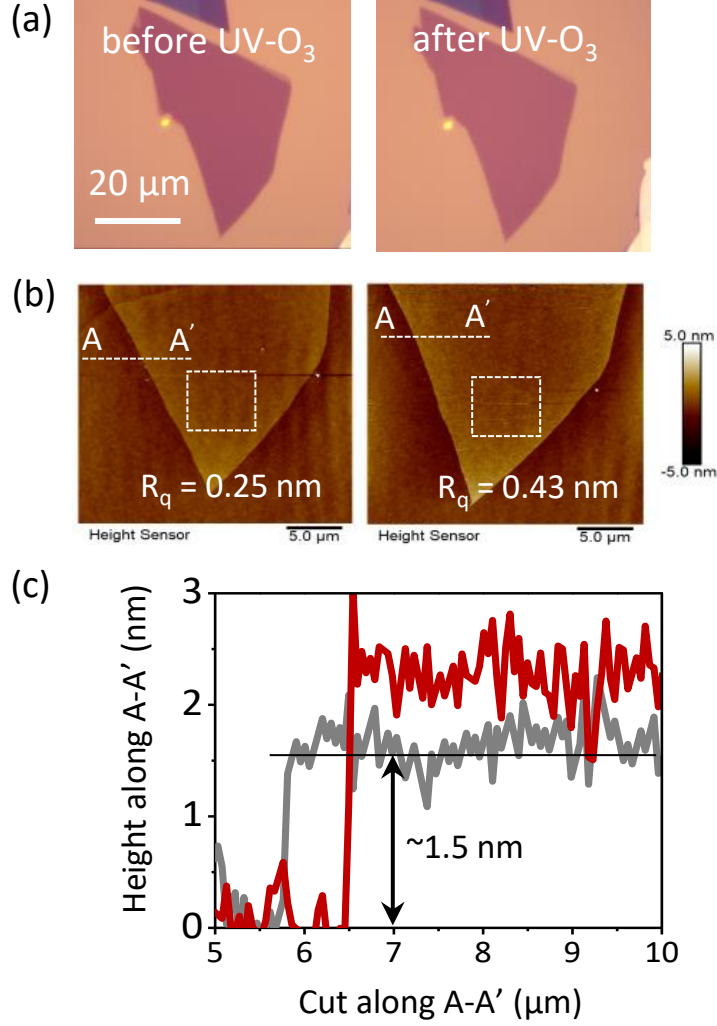


Figure S1: AFM characterization of the bilayer WSe<sub>2</sub> referred in Figure 1 of the main manuscript. (a) Optical microscope image of the flake before (left) and after (right) oxidation. No non-uniformity seen optically after oxidation. (b) AFM topography shows uniform and homogenous oxidation of the entire flake with a slight increase ( $< 2 \text{ \AA}$ ) in root mean square value of surface roughness,  $R_q$ , after oxidation. (c) The height maps along the A-A' cut on the flake before (gray) and after (red) oxidation. Post-oxidation map shows more than 1 nm increase in the height after oxidation due to different bonding configuration of TOS from WSe<sub>2</sub>. Further details on this characterization can be found in a different work from our group<sup>7</sup>.

## AFM characterization of 3L WSe<sub>2</sub>

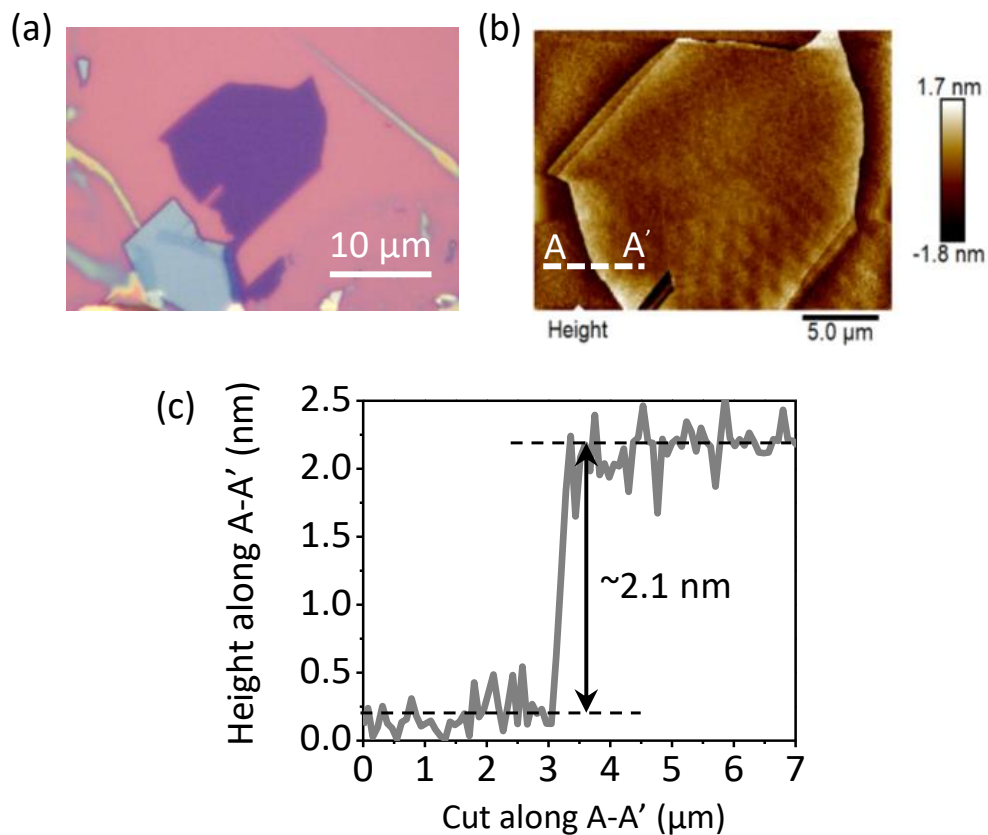


Figure S2: AFM characterization of the WSe<sub>2</sub> used in fabricating the blanket exposure device shown in Figure 2a of main manuscript. (a) The optical microscope image of the flake after exfoliation. (b) Topography of the flake measured before etching into a hall-bar pattern. (c) Height map shows  $\sim 2.1$  nm thick flake, which is the approximated thickness for 3L WSe<sub>2</sub><sup>8</sup>.

## Electrical characterization of multiple devices

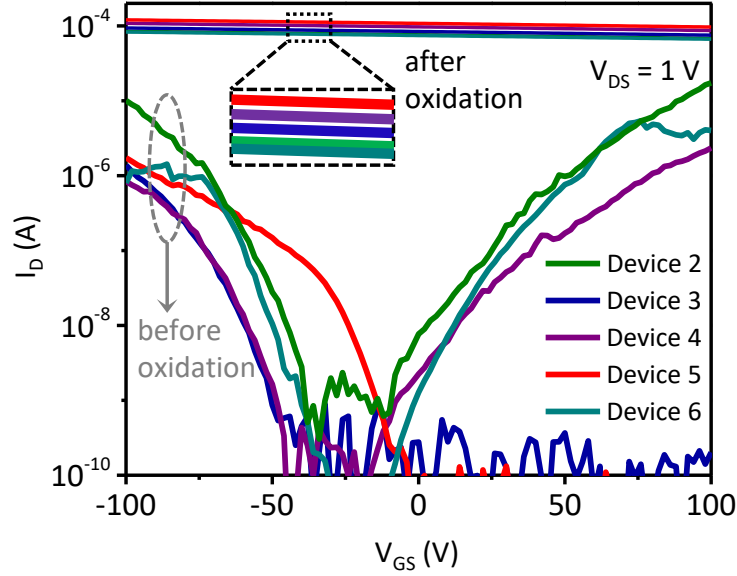


Figure S3: Transfer characteristics of multiple devices before (3L WSe<sub>2</sub>) and after oxidation (2L WSe<sub>2</sub> + 1L TOS). Data for device 1 is presented in the Figure 2c in the main manuscript. The channel length and width for all devices are 10  $\mu\text{m}$  and 4  $\mu\text{m}$  respectively. The variation in device-characteristics before oxidation is expected for 2D semiconductors and arises during standard lithography steps. Post-oxidation positive shift in threshold voltage and increase in the on-current shows high hole doping. The variation in the current levels post-doping is due to mobility variation (44–60  $\text{cm}^2/\text{V.s}$ ) among the channels, however the slopes of the transfer curves remain almost same which indicates similar levels of doping in each device ( $3\text{--}4 \times 10^{13} \text{ cm}^{-2}$ ).

## Schottky barrier height extraction

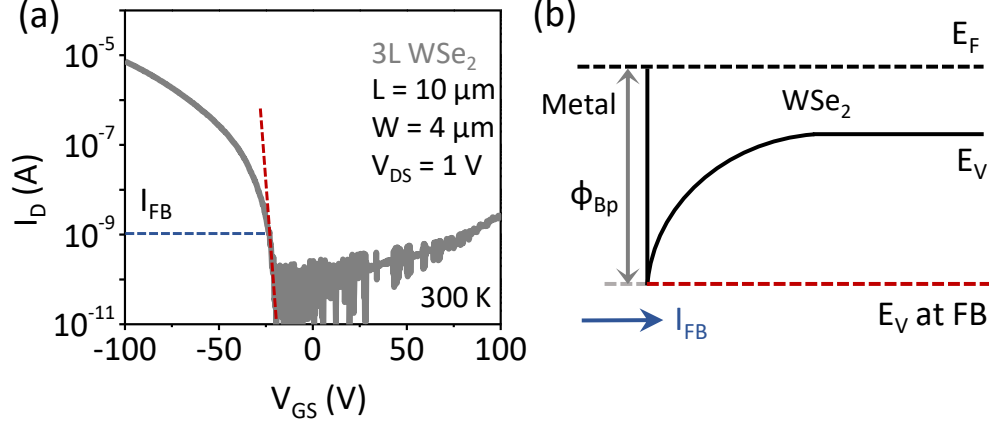


Figure S4: Schottky barrier height extraction for the device shown in Figure 2a in the main manuscript before oxidation (a) Transfer characteristics of the device. The dashed red line is used to find the current at flatband condition,  $I_{FB}$ . The  $V_{GS}$  at which the transfer characteristics starts deviating from exponential subthreshold behavior is identified as the flatband voltage. The current corresponding to this  $V_{GS}$  is the  $I_{FB}$  (for this device  $\sim 1$  nA), which we use to determine the Schottky barrier height before oxidation. (b) The tentative energy band diagram on the source side before doping.  $\Phi_{Bp}$  is the Schottky barrier height,  $E_F$  is the Fermi level and  $E_V$  is the valence band edge. The red dashed line shows the valence band edge at flatband condition.

The hole thermionic emission current over a Schottky barrier in a 2D semiconductor is given by<sup>9,10</sup>

$$I_{2D} = W A_{2D}^* T^{3/2} \exp\left(\frac{-q\Phi_{Bp}}{kT}\right) (1 - \exp(\frac{qV_D}{kT})) \quad (S1)$$

where  $W$  is the width of the semiconductor ( $4 \mu\text{m}$  for our device),  $A_{2D}^*$  is the modified 2D Richardson's constant,  $T$  is the temperature (300 K),  $\Phi_{Bp}$  is the hole Schottky barrier height,  $k$  is the Boltzmann constant,  $q$  is a unit charge and  $V_D$  is the applied bias (1 V for our measurements). The modified 2D Richardson can be calculated as

$$A_{2D}^* = \frac{q\sqrt{8\pi m^* k^3}}{h^2} \quad (S2)$$

where  $m^*$  is the hole effective mass ( $0.44 \times \text{mass of an electron for WSe}_2$ )<sup>11</sup> and  $h$  is the Placnk's constant. At the flatband condition the tunneling current through the barrier can be assumed zero and any transport is due to the thermionic emission over the barrier. This flatband condition is signified by the point in the transfer characteristics beyond which the current in the subthreshold region starts deviating from exponential behavior<sup>12</sup> (red dashed line in Figure S4). The current at this point is the flatband current,  $I_{FB}$ , which is found to be  $\sim 1$  nA for our device and plugging this  $I_{FB}$  for  $I_{2D}$  in Eq. (S1), we extract the Schottky barrier height as  $\sim 0.35$  eV.

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