

# **Identifying the Manipulation of Individual Atomic-Scale Defects for Boosting Thermoelectric Performances in Artificially Controlled $\text{Bi}_2\text{Te}_3$ Films**

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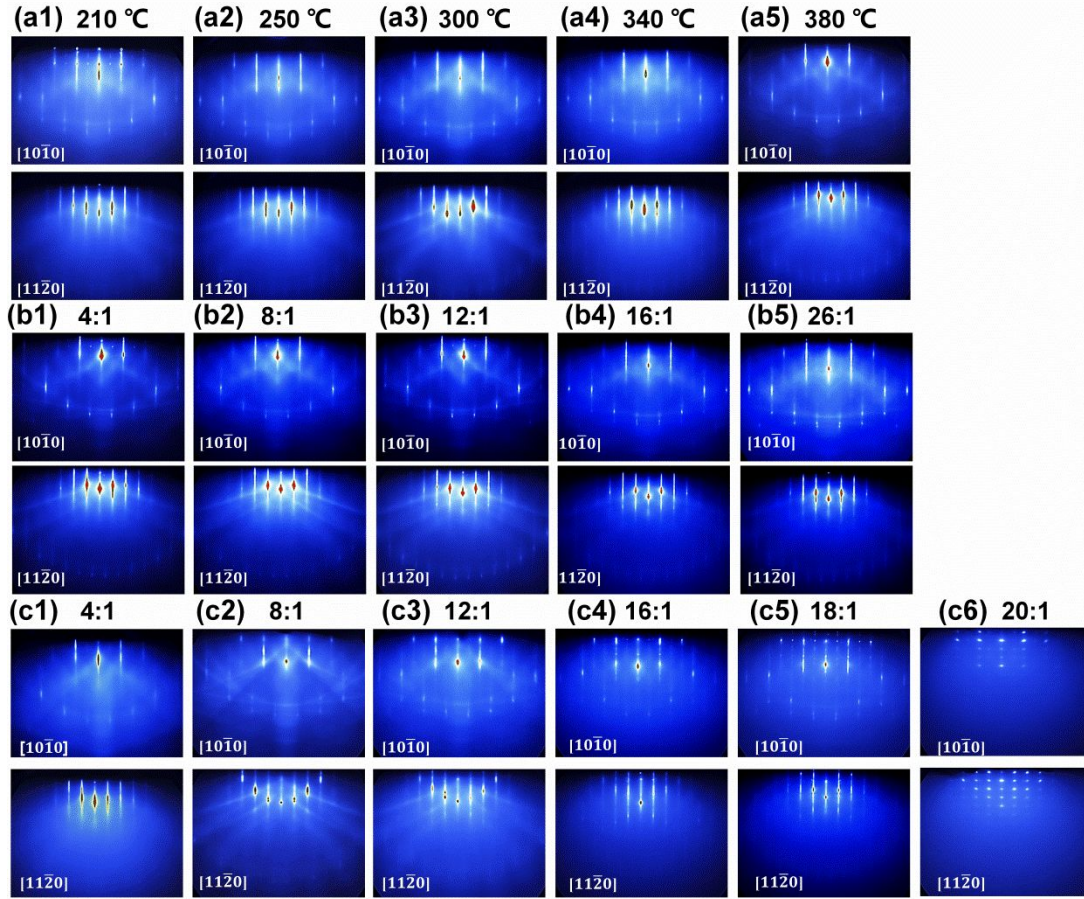
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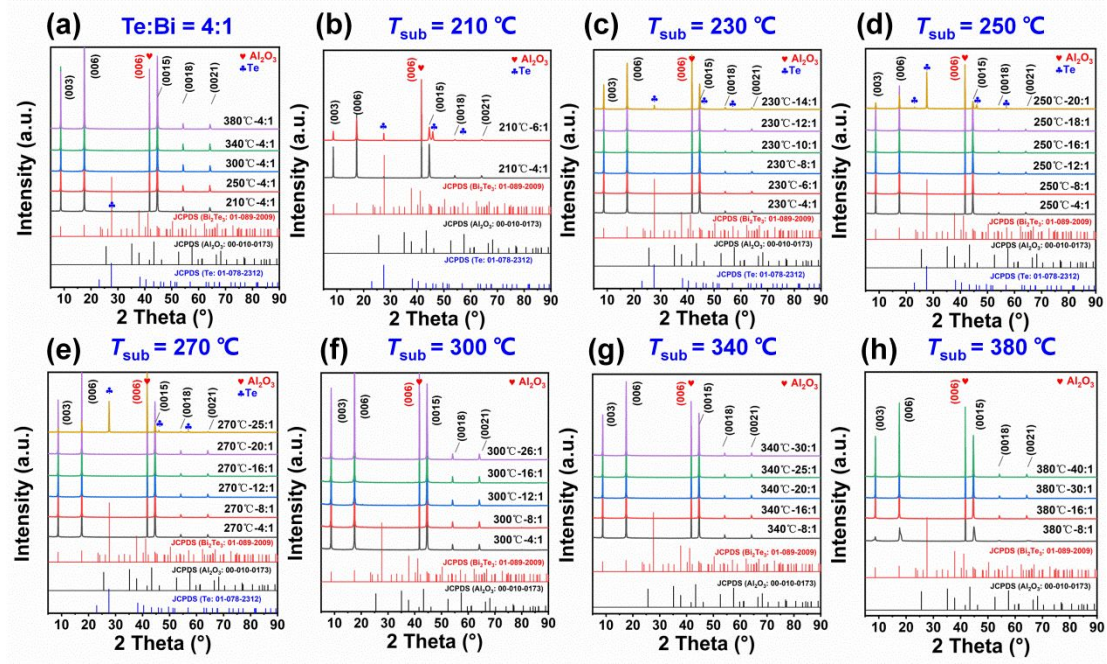
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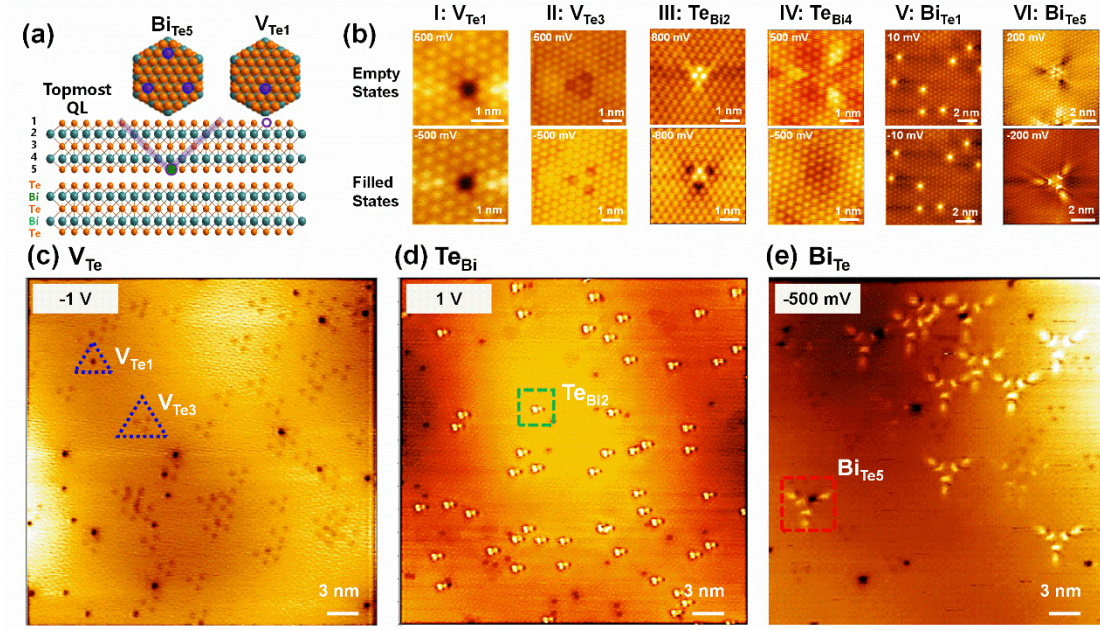


**Figure S1.** RHEED patterns along  $[10\bar{1}0]$  and  $[11\bar{2}0]$  azimuthal directions for  $\text{Bi}_2\text{Te}_3$  films grown at different substrate temperatures  $T_{\text{sub}}$  and Te/Bi flux ratios  $R$ . (a1-a5):  $R = 4$  and  $T_{\text{sub}} = 210$  °C (a1), 250 °C (a2), 300 °C (a3), 340 °C (a4), and 380 °C (a5); (b1-b5):  $T_{\text{sub}} = 300$  °C and  $R = 4:1$  (b1), 8:1 (b2), 12:1 (b3), 16:1 (b4), and 26:1 (b5);  $T_{\text{sub}} = 250$  °C and  $R = 4:1$  (c1), 8:1 (c2), 12:1 (c3), 16:1 (c4), 18:1 (c5), and 20:1 (c6). Well-defined and sharp RHEED streaks along the  $[10\bar{1}0]$  and  $[11\bar{2}0]$  azimuthal directions reveal the single crystalline nature of the MBE grown  $\text{Bi}_2\text{Te}_3$  films. The crystallinity of  $\text{Bi}_2\text{Te}_3$  films grown at lower  $T_{\text{sub}}$  and higher  $R$  becomes poor, which is evident from the transition of streaky patterns to spotty patterns, as shown in (c5) and (c6).



**Figure S2.** XRD patterns of  $\text{Bi}_2\text{Te}_3$  films grown at different  $T_{\text{sub}}$  and  $R$ :  $R = 4/1$  and  $T_{\text{sub}} = 210 \sim 380$  °C (a);  $T_{\text{sub}} = 210$  °C and  $R = 4/1$  and  $6/1$  (b);  $T_{\text{sub}} = 230$  °C and  $R = 4/1 \sim 14/1$  (c);  $T_{\text{sub}} = 250$  °C and  $R = 4/1 \sim 20/1$  (d);  $T_{\text{sub}} = 270$  °C and  $R = 4/1 \sim 25/1$  (e);  $T_{\text{sub}} = 300$  °C and  $R = 4/1 \sim 26/1$  (f);  $T_{\text{sub}} = 340$  °C and  $R = 8/1 \sim 30/1$  (g);  $T_{\text{sub}} = 380$  °C and  $R = 4/1 \sim 40/1$  (h). We observe obvious extra peaks, which could be indexed to Te precipitates, in  $\text{Bi}_2\text{Te}_3$  films grown at  $T_{\text{sub}} = 210$  °C and  $R = 6/1$ ,  $T_{\text{sub}} = 230$  °C and  $R = 14/1$ ,  $T_{\text{sub}} = 250$  °C and  $R = 20/1$  as well as  $T_{\text{sub}} = 270$  °C and  $R = 25/1$ , respectively, which is due to the excess of Te. For other  $\text{Bi}_2\text{Te}_3$  films, except for the peak at  $\sim 42^\circ$  originating from the sapphire substrate, all other diffraction peaks can be indexed to the standard (00 $l$ ) patterns of  $\text{Bi}_2\text{Te}_3$ , verifying their (00 $l$ ) single crystal characteristics.

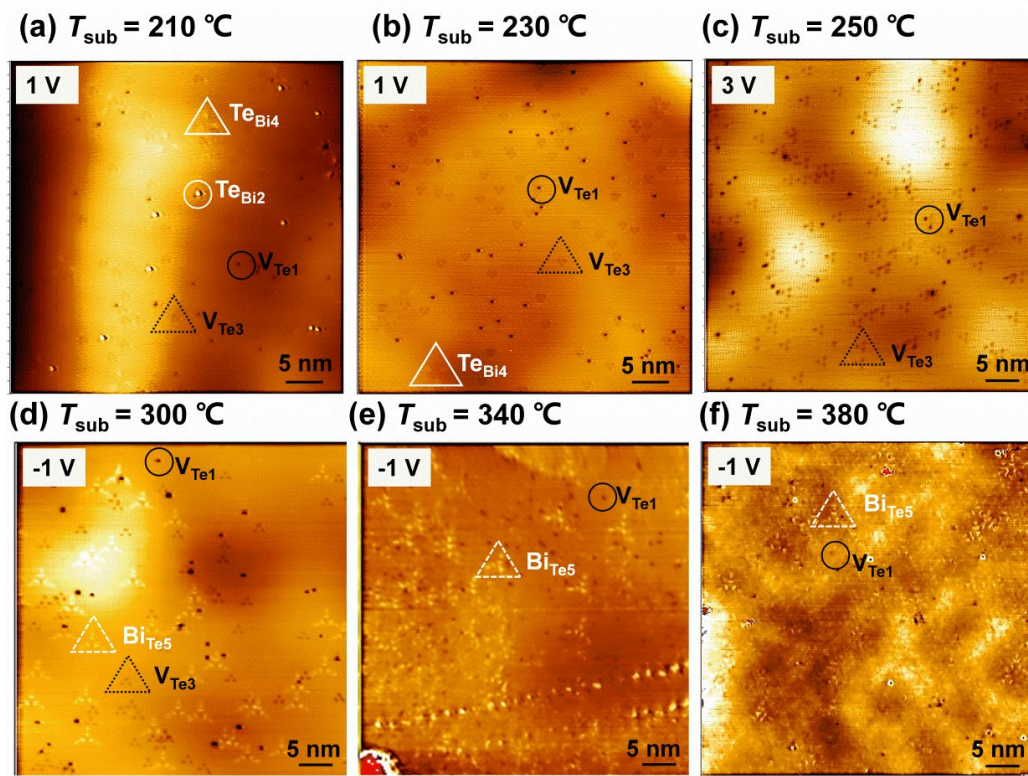




**Figure S3.** (a) Schematic diagram of the stacked quintuple layers for Bi<sub>2</sub>Te<sub>3</sub> and the STM surface topographies of point defects; (b) High resolution STM surface topographies of different intrinsic point defects in Bi<sub>2</sub>Te<sub>3</sub> films at the negative bias (filled states) and at the positive bias (empty states), respectively: (I) V<sub>Te1</sub> - dark holes, (II) V<sub>Te3</sub> - dark triangles, (III) Te<sub>Bi2</sub> - bright and dark triangles, (IV) Te<sub>Bi4</sub> - bright trefoil and dark triangle, (V) Bi<sub>Te1</sub> - bright spots, and (VI) Bi<sub>Te5</sub> - depressed and bright double triangles. The STM images (30 nm × 30 nm) of Bi<sub>2</sub>Te<sub>3</sub> films with the different dominant surface point defects of V<sub>Te</sub> (c), Te<sub>Bi</sub> (d) and Bi<sub>Te</sub> (e) grown at  $T_{\text{sub}} = 250$  °C and  $R = 4/1$ ,  $T_{\text{sub}} = 250$  °C and  $R = 12/1$  as well as  $T_{\text{sub}} = 300$  °C and  $R = 4/1$ , respectively. In (a), the lateral numbers of 1 to 5 indicate the  $N$ th atomic layers in the topmost QL.

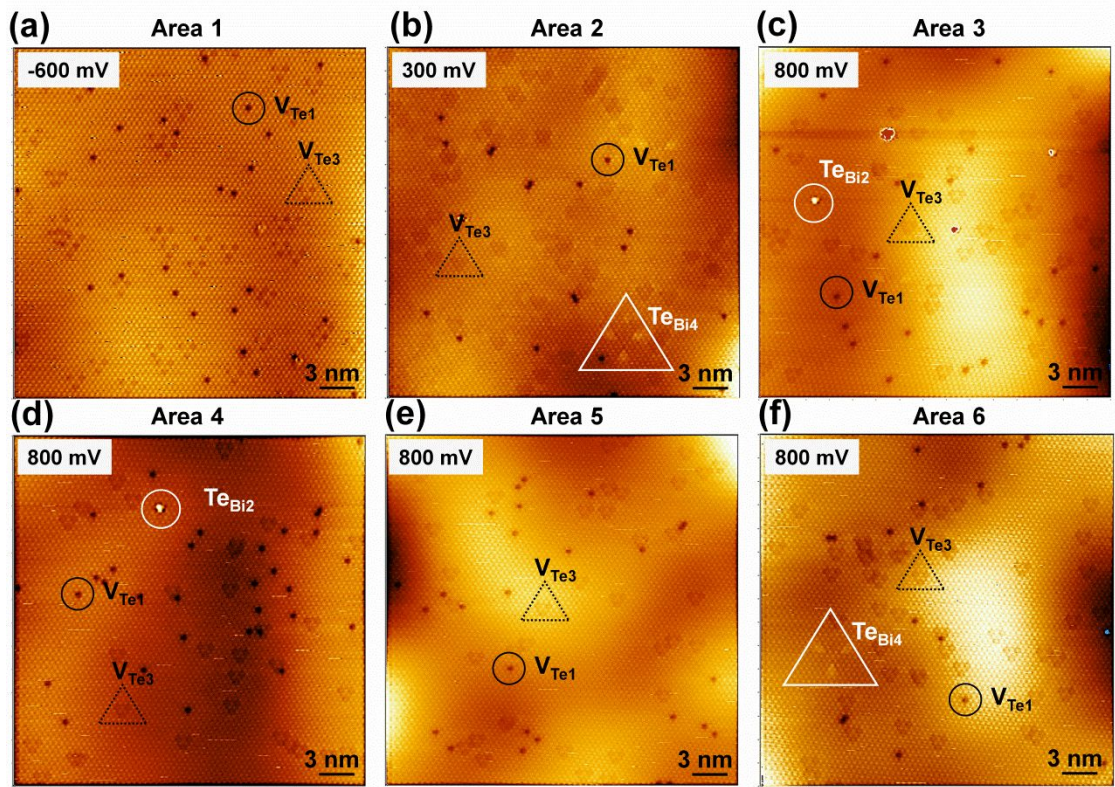
In Bi<sub>2</sub>Te<sub>3</sub>, the STM surface morphology of intrinsic point defects is the projection of the local electron density of states of the defects onto the outermost Te layer,<sup>1</sup> which is intimately related to the type of defects, their locations in the QL, the in-plane threefold symmetry, and the coordination characteristics of Bi and Te. The Te vacancy on the first atomic layer, denoted as V<sub>Te1</sub>, appears as a vacant site with a significant depression at positive or negative sample bias. Similar to V<sub>Te1</sub>, the Te vacancy at the third atomic layer (V<sub>Te3</sub>) forms a dark triangular shape morphology. As for the STM topography on a flat surface, the donor-like point defect will exhibit a bright protrusion while the acceptor-like point defect will present a dark depression at

the negative sample bias.<sup>2,3</sup> The opposite image contrast will be observed at the positive sample bias. Thus, the antisites  $\text{Te}_{\text{Bi}2}$  and  $\text{Te}_{\text{Bi}4}$ , as Te occupies the sites of Bi at the second and fourth atomic layer, show a bright triangular and a bright trefoil shape structure, respectively, at positive sample bias, and show dark triangular regions at negative sample bias, verifying that they are electron donors. In addition, the antisites  $\text{Bi}_{\text{Te}}$  appear chiefly at the first and fifth atomic layer, marked as  $\text{Bi}_{\text{Te}1}$  and  $\text{Bi}_{\text{Te}5}$ , respectively. The  $\text{Bi}_{\text{Te}1}$  defect always shows a bright spot at negative or positive sample bias<sup>2,3</sup> owing to the much larger atomic radius of Bi than Te. The  $\text{Bi}_{\text{Te}5}$  defect exhibits a depressed and a bright double-triangle morphology at the positive and negative bias, respectively, implying their acceptor nature.

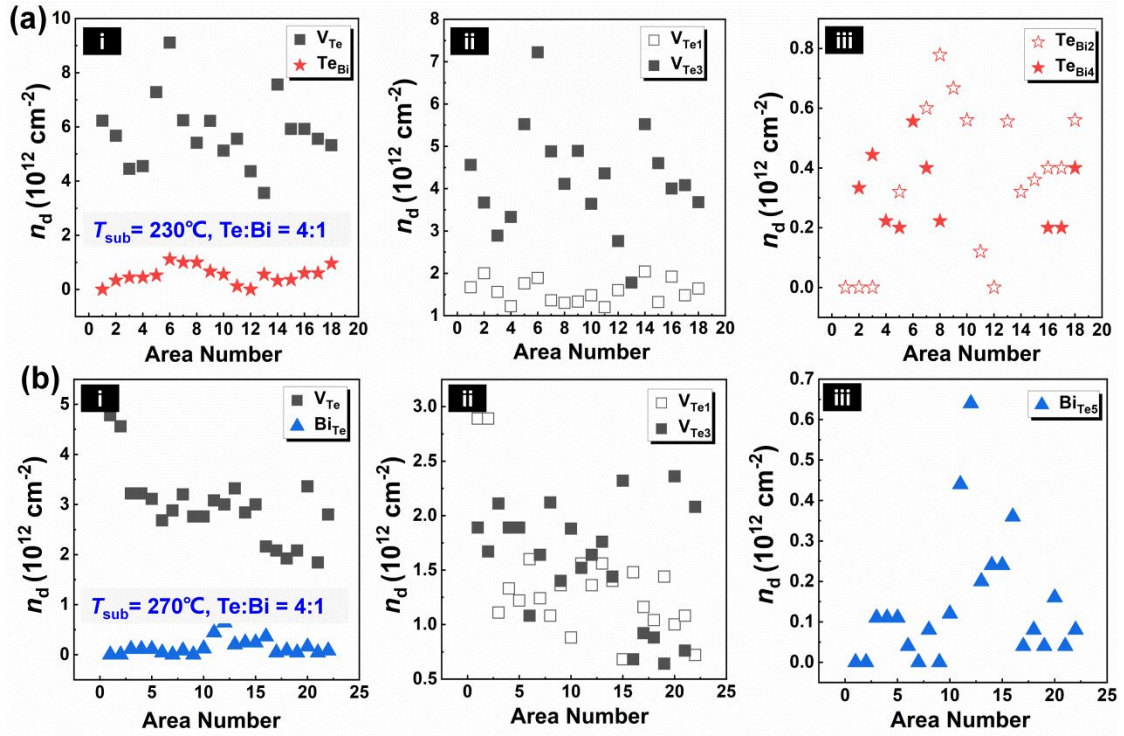


**Figure S4.** STM images ( $50 \text{ nm} \times 50 \text{ nm}$ ) of  $\text{Bi}_2\text{Te}_3$  films grown at  $R = 4/1$  and  $T_{\text{sub}} = 210 \text{ }^\circ\text{C}$  (a),  $230 \text{ }^\circ\text{C}$  (b),  $250 \text{ }^\circ\text{C}$  (c),  $300 \text{ }^\circ\text{C}$  (d),  $340 \text{ }^\circ\text{C}$  (e), and  $380 \text{ }^\circ\text{C}$  (f). Higher  $T_{\text{sub}}$  will result in a more serious re-evaporation of Te, and thus obviously facilitate the formation of vacancies  $V_{\text{Te}}$  and  $\text{Bi}_{\text{Te}}$  antisites as well as their *in-situ* transformation  $V_{\text{Te}} \rightarrow \text{Bi}_{\text{Te}}$ .



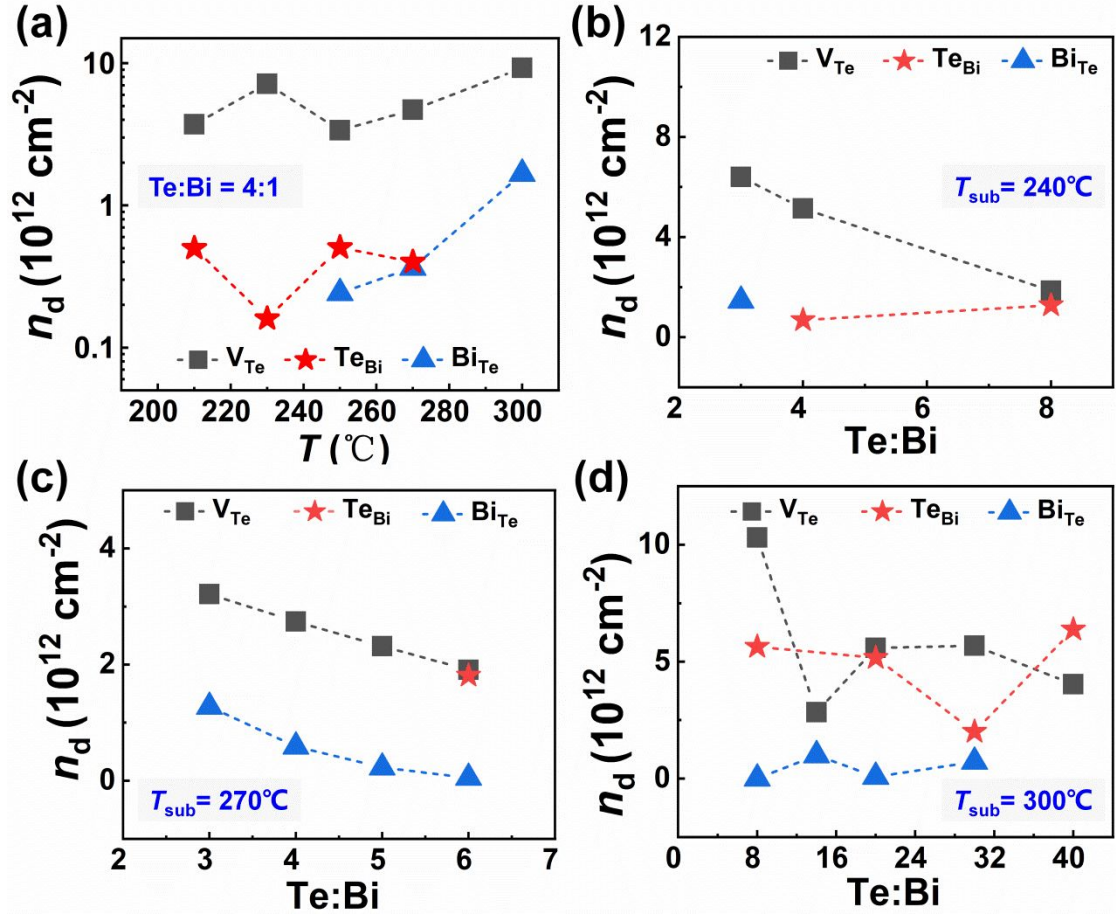


**Figure S5.** STM images from different areas of the  $\text{Bi}_2\text{Te}_3$  film grown at  $T_{\text{sub}} = 230\text{ }^\circ\text{C}$  and  $R = 4/1$ . Six large-area ( $30\text{ nm} \times 30\text{ nm}$ ) STM images show similar features of the point defects, which contain dominant vacancies  $V_{\text{Te}}$  and also a minority of  $\text{Te}_{\text{Bi}}$  antisites.

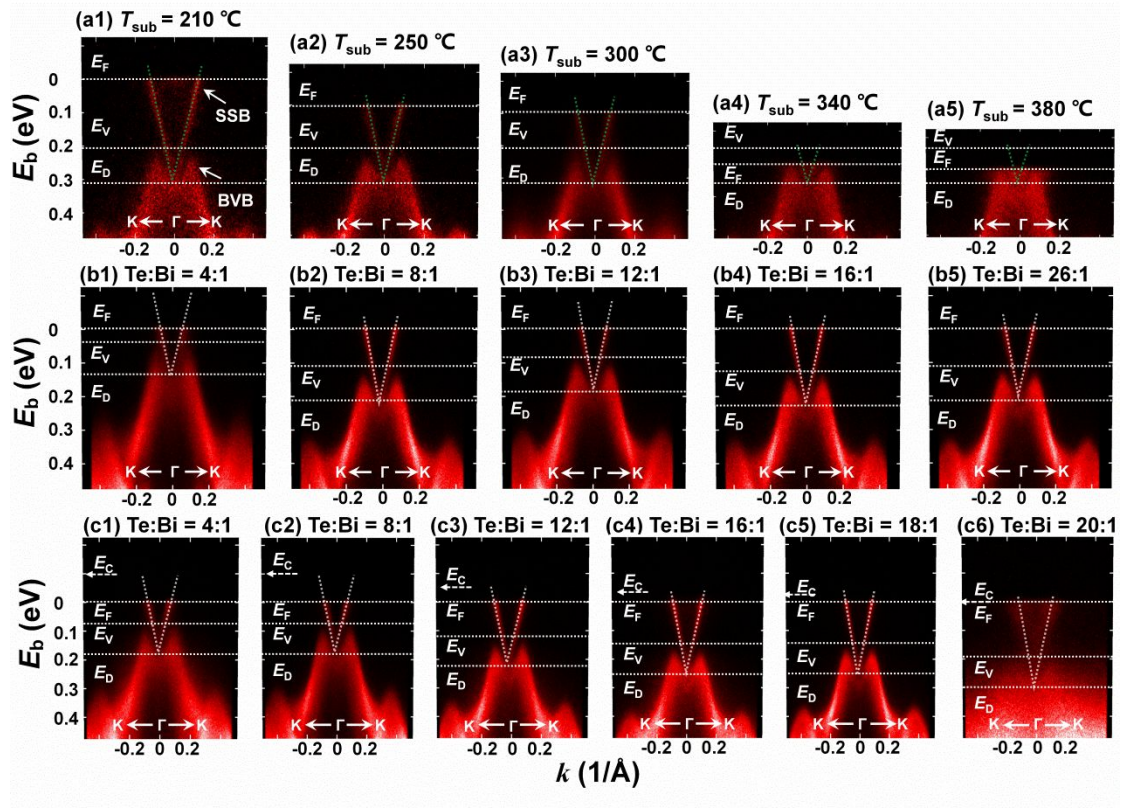


**Figure S6.** Statistical analysis of the density of surface point defects in different areas of  $\text{Bi}_2\text{Te}_3$  films grown at (a)  $T_{\text{sub}} = 230^\circ\text{C}$  and  $R = 4/1$  as well as (b)  $T_{\text{sub}} = 270^\circ\text{C}$  and  $R = 4/1$





**Figure S7.** The summary of the surface defect density for  $\text{Bi}_2\text{Te}_3$  films grown at different  $T_{\text{sub}}$  and  $R$ :  $R = 4/1$  and  $T_{\text{sub}} = 210 \sim 300^\circ\text{C}$  (a);  $T_{\text{sub}} = 240^\circ\text{C}$  and  $R = 3/1 \sim 8/1$  (b);  $T_{\text{sub}} = 270^\circ\text{C}$  and  $R = 3/1 \sim 6/1$  (c);  $T_{\text{sub}} = 300^\circ\text{C}$  and  $R = 8/1 \sim 40/1$  (d). The data are summarized based on the statistical analysis in various measurement areas, as shown in Figs. S4 and S5.



**Figure S8.** The ARPES electronic band structures along the  $\Gamma$ -K direction measured at 10 K for  $\text{Bi}_2\text{Te}_3$  films grown at various  $T_{\text{sub}}$  and  $R$ . (a1-a5):  $R = 4$  and  $T_{\text{sub}} = 210$  °C (a1), 250 °C (a2), 300 °C (a3), 340 °C (a4), and 380 °C (a5); (b1-b5):  $T_{\text{sub}} = 300$  °C and  $R = 4:1$  (b1), 8:1 (b2), 12:1 (b3), 16:1 (b4), and 26:1 (b5);  $T_{\text{sub}} = 250$  °C and  $R = 4:1$  (c1), 8:1 (c2), 12:1 (c3), 16:1 (c4), 18:1 (c5), and 20:1 (c6). The  $E_D$ ,  $E_V$ ,  $E_C$  and  $E_F$  are the energy positions of the Dirac point, the valence band maximum, the conduction band minimum and the Fermi level, while BVB and SSB represent the bulk valence band and the surface state band, respectively. With increasing the  $T_{\text{sub}}$  from 210 °C to 380 °C, an obvious downshift of the  $E_F$ (ARPES) is observed for  $\text{Bi}_2\text{Te}_3$  films, due to the increased density of  $p$ - $\text{Bi}_{\text{Te}}$  defects. In addition, with increasing the  $R$  value, the  $E_F$ (ARPES) of  $\text{Bi}_2\text{Te}_3$  films gradually moves towards the conduction band due to the increased density of  $n$ - $\text{Te}_{\text{Bi}}$  defects.

**Table S1:** Room temperature physical parameters of Bi<sub>2</sub>Te<sub>3</sub> films grown at different  $T_{\text{sub}}$  and  $R$ . ( $d$  stands for the film thickness,  $D$  is the grain size,  $n_{\text{H}}$  is the electron density,  $\mu$  is the carrier mobility,  $\sigma$  is the electrical conductivity,  $S$  is the Seebeck coefficient,  $PF$  is the power factor,  $m^*$  stands for carrier effective mass and  $l$  is the mean free path of electrons) The positive  $S$  is found in the Bi<sub>2</sub>Te<sub>3</sub> films grown at  $R = 4/1$  as well as at  $T_{\text{sub}} = 340$  °C and 380 °C, indicating the  $p$ -Bi<sub>Te</sub> defects dominate the electronic transport. The grain size  $D$  and the mean free path  $l$  are calculated similarly as in references 4–6.

T °C	Te: Bi	$d$ nm	$D$ nm	$n_{\text{H}}$ $10^{19} \text{ cm}^{-3}$	$\mu$ $\text{cm}^2 \text{ V}^{-1} \text{ s}^{-1}$	$\sigma$ $10^4 \text{ S m}^{-1}$	$S$ $\mu\text{V K}^{-1}$	$PF$ $\text{mW m}^{-1} \text{ K}^{-2}$	$m^*$ $m_e$	$l$ nm
210	4:1	138	81	4.0	55	3.5	-210	1.5	1.2	4.3
250	4:1	125	112	5.9	30	2.9	-127	0.5	1.0	2.2
300	4:1	122	112	6.4	25	2.3	-149	0.5	1.1	1.9
340	4:1	108	85	9.6	36	5.6	3	0.0006	\	\
380	4:1	98	84	47.2	7	5.3	9	0.004	\	\
210	4:1	60	81	8.6	71	9.7	-182	3.23	1.8	6.8
210	6:1	175	132	4.8	32	2.3	-104	0.25	0.7	1.8
230	4:1	67	97	4.5	40	2.9	-137	0.54	0.9	2.7
230	6:1	53	104	5.5	53	4.7	-190	1.70	1.4	4.5
230	8:1	66	85	8.3	63	8.4	-186	2.93	1.8	6.1
230	10:1	67	85	8.9	66	9.4	-210	4.14	2.1	6.9
230	12:1	66	161	4.9	53	4.1	-198	1.61	1.3	4.4
240	4:1	84	85	4.7	43	3.2	-125	0.49	0.8	2.8
240	8:1	85	91	4.0	67	4.3	-175	1.32	1.1	4.9
240	12:1	87	145	4.0	67	4.3	-216	2.02	1.3	5.4
240	14:1	94	130	5.1	59	4.8	-220	2.33	1.5	5.2
240	16:1	95	112	7.3	82	9.5	-189	3.39	1.6	7.5
240	16:1	210	97	5.4	90	7.7	-209	3.38	1.5	7.9
250	4:1	62	66	3.9	78	4.8	-160	1.23	1.0	5.3
250	8:1	63	85	6.2	63	6.2	-165	1.70	1.1	5.0
250	12:1	67	90	7.0	70	7.9	-163	2.10	1.4	6.0
250	16:1	68	104	8.2	87	9.8	-183	3.26	1.7	8.2
250	18:1	68	85	11.1	81	14.4	-180	4.66	2.1	8.4
250	20:1	71	66	15.0	46	11.0	-161	2.86	2.3	5.0

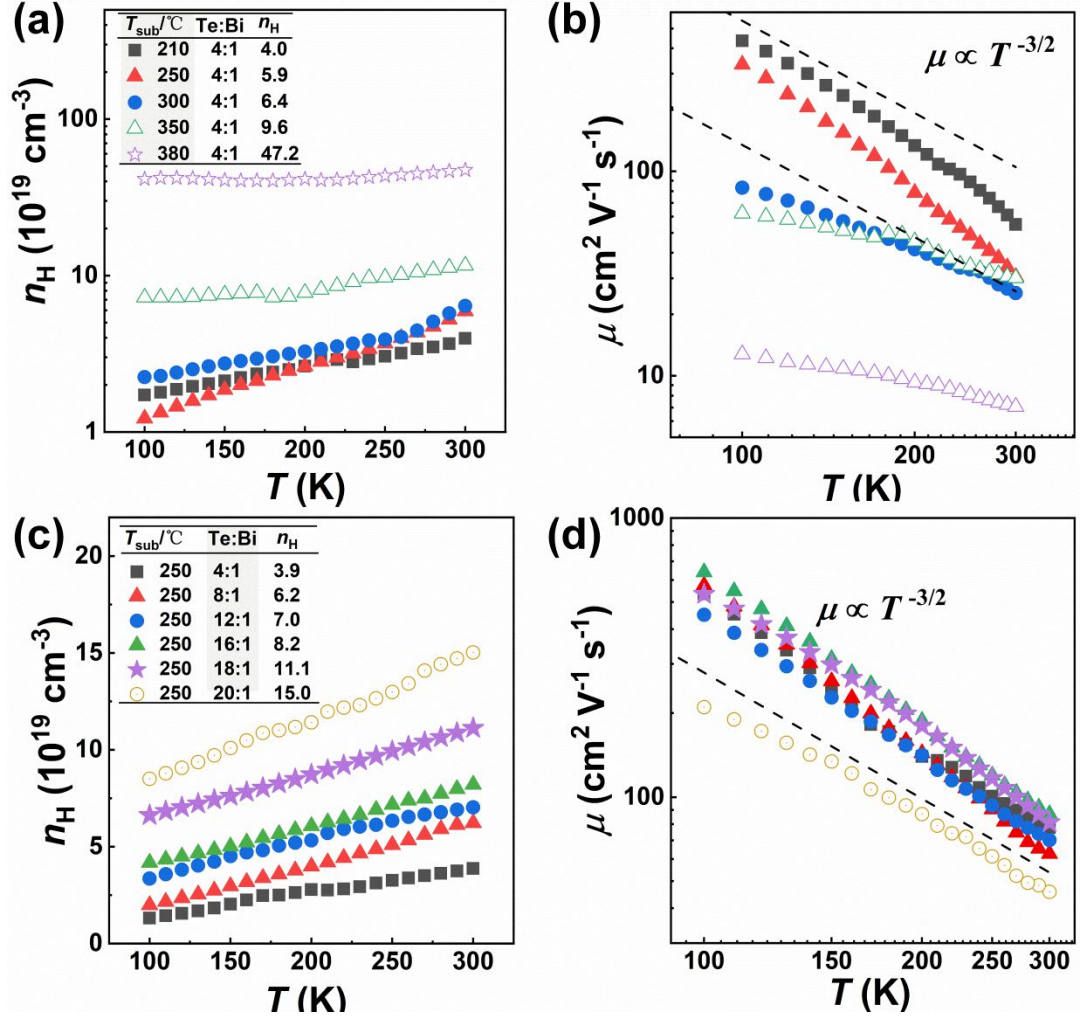
**Table S2:** Room temperature physical parameters of Bi<sub>2</sub>Te<sub>3</sub> films grown at different  $T_{\text{sub}}$  and  $R$ . The positive  $S$  of the Bi<sub>2</sub>Te<sub>3</sub> film grown at  $T_{\text{sub}} = 270$  °C and  $R = 30:1$  is caused by the presence of  $p$ -Te precipitates in the matrix as in references 7-8.

T	Te:B	$d$	$D$	$n_{\text{H}}$	$\mu$	$\sigma$	$S$	$PF$	$m^*$	$l$
°C	i	nm	nm	$10^{19} \text{ cm}^{-3}$	$\text{cm}^2 \text{ V}^{-1} \text{ s}^{-1}$	$10^4 \text{ S m}^{-1}$	$\mu\text{V K}^{-1}$	$\text{mW m}^{-1} \text{ K}^{-2}$	$m_e$	nm
270	4:1	54	76	8.4	31	4.16	-124	0.62	1.2	2.4
270	8:1	60	85	6.9	39	4.36	-147	0.94	1.2	3.2
270	12:1	62	76	7.2	35	4.0	-128	0.67	1.1	2.6
270	16:1	67	85	7.9	31	3.98	-167	1.11	1.5	2.8
270	20:1	60	97	8.0	50	6.4	-194	2.42	1.7	4.8
270	25:1	72	69	6.4	38	3.89	-91	0.30	0.7	2.3
270	30:1	263	85	4.2	15	1.02	15	0.0023	/	/
300	4:1	59	63	12.4	26	5.1	-112	0.64	1.4	2.2
300	8:1	58	104	10.2	21	3.4	-149	0.76	1.6	1.9
300	16:1	61	69	7.0	44	4.9	-135	0.89	1.1	3.3
300	20:1	77	91	5.9	47	4.4	-131	0.75	1.0	3.4
300	26:1	63	63	6.2	49	4.9	-162	1.28	1.3	4.0
340	8:1	59	104	42.6	5	3.40	-31	0.03	0.9	0.3
340	16:1	61	112	8.1	28	3.61	-135	0.66	1.3	2.3
340	20:1	48	85	7.6	41	5.1	-151	1.17	1.3	3.5
340	25:1	62	104	8.4	26	3.56	-118	0.50	1.1	2.0
340	30:1	64	104	8.5	27	3.65	-134	0.66	1.3	2.2
380	8:1	42	50	97.2	13	19.8	-69	0.93	3.4	1.7
380	16:1	58	132	10.7	23	3.99	-115	0.53	1.3	1.9
380	30:1	54	91	13.9	17	3.75	-31	0.04	0.4	0.8
380	40:1	62	81	9.8	23	3.59	-72	0.19	0.8	1.4



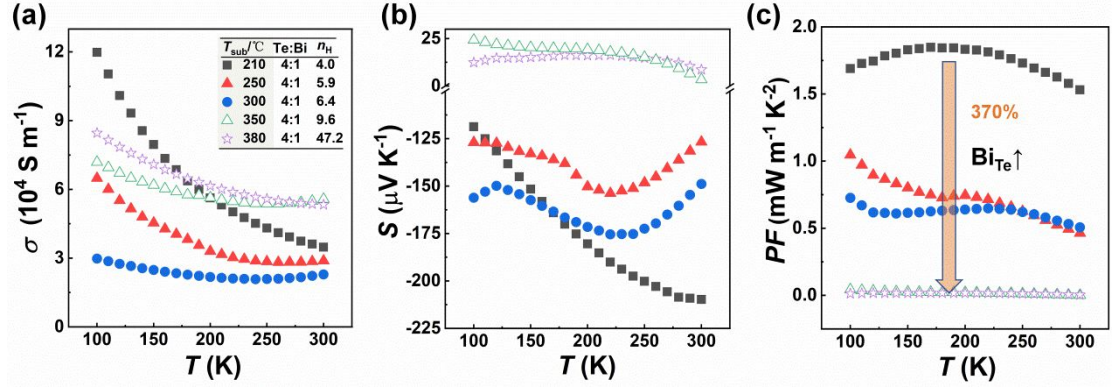
**Table S3:** Room temperature  $\sigma$ ,  $S$  and  $PF$  of  $\text{Bi}_2\text{Te}_3$  films grown at  $T_{\text{sub}} = 250$  °C with different  $R$ , measured by the PPMS and ZEM-3 apparatus. The deviation from two measurements is within  $\pm 10\%$ , majorly due to the measurement uncertainty from two different testing equipments.

T °C	Te:Bi	From PPMS			From ZEM-3		
		$\sigma$ $10^4 \text{ S m}^{-1}$	$S$ $\mu\text{V K}^{-1}$	$PF$ $\text{mW m}^{-1} \text{ K}^{-2}$	$\sigma$ $10^4 \text{ S m}^{-1}$	$S$ $\mu\text{V K}^{-1}$	$PF$ $\text{mW m}^{-1} \text{ K}^{-2}$
250	4:1	4.8	-160	1.23	4.0	-152	0.92
250	8:1	6.2	-165	1.70	4.2	-164	1.13
250	12:1	7.9	-163	2.10	6.6	-173	1.97
250	16:1	9.8	-183	3.26	8.9	-191	3.25
250	18:1	14.4	-180	4.66	12.8	-186	4.43
250	20:1	11.0	-161	2.86	10.1	-162	2.65

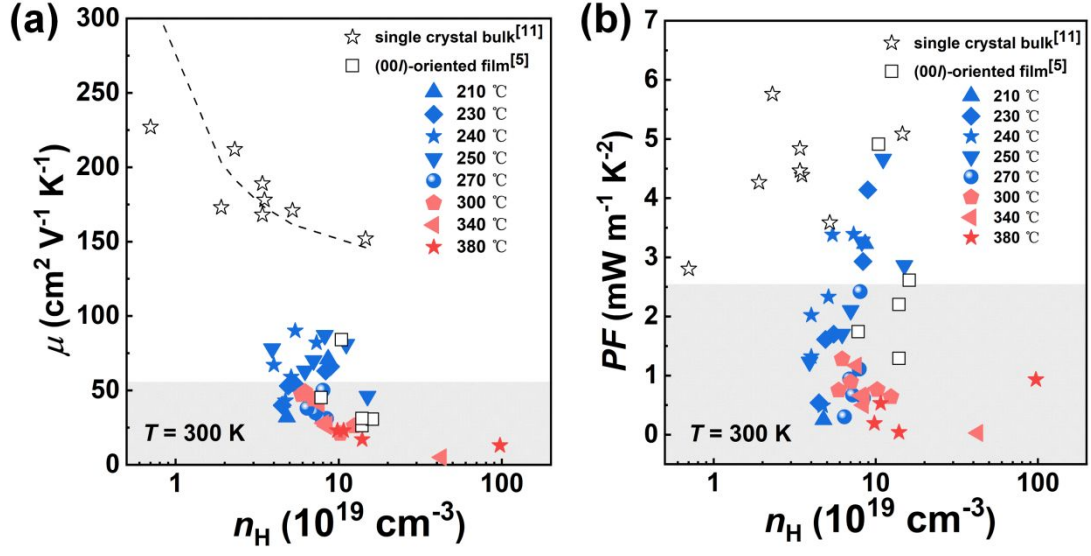


**Figure S9.** (a-b) Temperature dependence of  $n_H$  and  $\mu$  of the  $\text{Bi}_2\text{Te}_3$  films grown at  $R = 4/1$  and  $T_{\text{sub}} = 210 \sim 380^\circ\text{C}$ ; (c-d) Temperature dependence of  $n_H$  and  $\mu$  of the  $\text{Bi}_2\text{Te}_3$  films grown at  $T_{\text{sub}} = 250^\circ\text{C}$  and  $R = 4/1 \sim 20/1$ .

The  $n_H$  rises sharply with the increasing temperature, which suggests the formation of shallow donor levels near the conduction band minimum associated with the intrinsic point defects.<sup>6,9</sup> The mobility  $\mu$  follows the  $T^{-3/2}$  dependence, which suggests that the acoustic phonon scattering becomes the dominant carrier scattering mechanism at room temperature.<sup>10</sup>



**Figure S10.** Temperature dependence of  $\sigma$ ,  $S$  and  $PF$  of the  $\text{Bi}_2\text{Te}_3$  films grown at  $R = 4/1$  and  $T_{\text{sub}} = 210 \sim 380$  °C.



**Figure S11.** (a) The dependence of the carrier mobility  $\mu$  on the carrier concentration  $n_H$  and (b) the dependence of the power factor  $PF$  on the carrier concentration  $n_H$  in  $n$ -Bi<sub>2</sub>Te<sub>3</sub> films at room temperature. The literature data for single crystal bulks and films are included for comparison.



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