# Supplemental Material for "Spin-Orbit Coupling Determined Topological Phase: Topological Insulator and Quadratic Dirac

## Semimetals"

Lu Tian, Ying Liu\*, Weizhen Meng, Xiaoming Zhang, Xuefang Dai, and Guodong Liu\*

School of Materials Science and Engineering, Hebei University of Technology, Tianjin 300130, China. E-mail: ying liu@hebut.edu.cn; gdliu1978@126.com

#### A. Surface Band Structure for the Triple Degeneracy

To illustrate the topology of triple degenerate point in  $Y_3XC$  (X=Ga, Tl), we plot the surface spectrum on (001) surface. Obviously, there exists surfaces states starting from the triple degenerate point (see in Fig. S1).

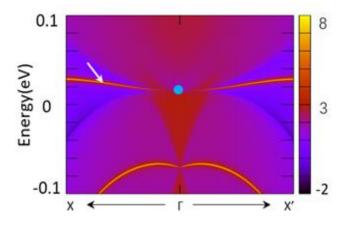


Fig.S1 Surface spectrum for the triple degenerate in Y<sub>3</sub>GaC on (001) surface. The white arrow points to the surface states, cyan dot marks the triple degenerate point.

### **B.** Topological Phase Transition Induced by a Zeeman Field

In this SM, we use the model to study topological phase transition under a Zeenman field. We have mentioned that the effective model we have derived Eq. (2) is based on states of j = 3/2, then the Zeeman field can be written in the form of

 $\mathcal{H}_{\mathrm{Z}} = -\boldsymbol{B} \cdot \boldsymbol{J},$ 

with **J** as the vector of the J = 3/2 spin matrices.

Under a Zeeman field, the time reversal symmetry is broken, the quadratic Dirac point is expected to be transformed into pairs of Weyl points. To demonstrate it, we apply a Zeeman field along high symmetric path, namely, [100] direction.

In this case,  $\mathcal{H}_{Z} = -MJ_{x}$ , the total Hamiltonian takes the form of  $\mathcal{H} = \mathcal{H}_{soc} + \mathcal{H}_{Z}$ . We first search band crossings along  $k_{x}$  axis, then two pairs (W1, W2) of Weyl points are expected to appear. For concreteness, the eigenenergies for  $\mathcal{H}$  along  $k_{x}$ are given by

$$E_{1,2} = -\alpha k_x^2 \pm \frac{B}{2}, \quad E_{3,4} = \alpha k_x^2 \pm \frac{3B}{2}.$$

Accordingly, if  $\alpha > 0$ ,  $E_{1,2}$  will cross with  $E_4$  at  $K_1 = \left(\pm \sqrt{\frac{B}{\alpha}}, 0, 0\right), K_2 =$ 

 $\left(\pm\sqrt{\frac{B}{2\alpha}},0,0\right)$ , respectively. Taking suitable parameters, we present the energy

spectrum under a Zeeman field around the quadratic Dirac point in Fig.S1.

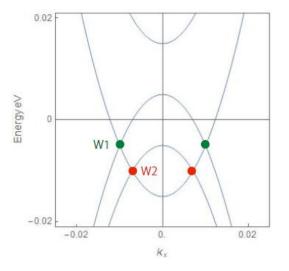


Fig.S2 Band structure under a Zeeman field along [100] direction. In the calculation, we take M = 0, B = 10 meV,  $\alpha = 10$  eV  $\cdot$  Å,  $\gamma = 17$  eV  $\cdot$  Å.

#### **Effective Hamiltonian for Weyl Points**

For the first pair of Weyl points which locate at  $K_1 = \left(\pm \sqrt{\frac{B}{\alpha}}, 0, 0\right)$ , one can derive the

degenerate eigenstates at these two points, specifically,

$$|\psi_1\rangle = \frac{1}{\sqrt{2}}(-1,0,0,1)^T, \ |\psi_2\rangle = \frac{1}{\sqrt{2}}(0,-1,1,0)^T.$$

Then the effective Hamiltonian around  $K_1$  is given by

$$\mathcal{H}_{1} = \begin{pmatrix} \langle \psi_{1} | \mathcal{H}_{soc} | \psi_{1} \rangle & \langle \psi_{1} | \mathcal{H}_{soc} | \psi_{2} \rangle \\ \langle \psi_{2} | \mathcal{H}_{soc} | \psi_{1} \rangle & \langle \psi_{2} | \mathcal{H}_{soc} | \psi_{2} \rangle \end{pmatrix}.$$

Then we expand the model at  $K_1$ , after a unitary transformation, we have the Hamiltonian takes the form as

$$\mathcal{H}_{1}(\boldsymbol{q}) = v_{x}q_{x}\sigma_{z} + [(\gamma_{+}q_{-}^{2} + \gamma_{-}q_{+}^{2})\sigma_{+} + h.c.],$$

with  $\gamma_{\pm} = \frac{1}{8} (\sqrt{3}\alpha \pm \gamma), q_{\pm} = q_y \pm iq_z, \sigma_{\pm} = \sigma_x \pm i\sigma_y$ . The chirality of the Weyl point is determined by sgn  $(|\gamma_+| - |\gamma_-|) = \pm 2$ .

The effective Hamiltonian expanded from  $K_2 = \left(\pm \sqrt{\frac{B}{2\alpha}}, 0, 0\right)$  can be derived in the

same way. Here, the degenerate states take the form of

$$\psi_3 = \frac{1}{2\sqrt{2}} (-2, -\sqrt{3}, 0, 1)^T, \qquad \psi_4 = \frac{1}{2\sqrt{2}} (0, 1, 2, \sqrt{3})^T.$$

Such that the effective Hamiltonian for the band crossing around  $K_2$  is given by

$$\mathcal{H}_{2} = \begin{pmatrix} \langle \psi_{3} | \mathcal{H}_{soc} | \psi_{3} \rangle & \langle \psi_{3} | \mathcal{H}_{soc} | \psi_{4} \rangle \\ \langle \psi_{4} | \mathcal{H}_{soc} | \psi_{3} \rangle & \langle \psi_{4} | \mathcal{H}_{soc} | \psi_{4} \rangle \end{pmatrix}.$$

To derive the main feature at K2, we also expand it at this point, thus

$$\mathcal{H}_2(\boldsymbol{q}) = v_x q_x \sigma_z + v_y q_y \sigma_y - v_z q_z \sigma_x.$$

Here,  $v_y = v_z = \gamma \sqrt{\frac{M}{\alpha}}$ ,  $v_x = 2\sqrt{M\alpha}$ . This model is a typical single Weyl point.