Supporting information

Magnetized MXene microspheres with multi-scale magnetic coupling and enhanced polarized interfaces for distinct microwave absorption via a spray-drying method

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Figure S1. SEM images of preparation M/F composites at lower temperature (a, b), at higher suspension concentration (c, d).



Figure S2. SEM image of accordion-like MXene after the ultrasonic cell disrupting.



Figure S3. SEM image of M/F composites.



Figure S4. SEM images of M/F-m sample at lower magnification (a) and higher magnification (b).



Figure S5. HRTEM image of MXene nanosheets (a); corresponding areas (a1, a2 and a3) and corresponding strain maps (b1, b2 and b3) of the MXene nanosheets. The color variation indicated the strain values ranging from -0.7 to 0.7.



Figure S6. ε'' value versus frequency of M/F composite.



Figure S7. Dielectric loss tangent (a) and magnetic loss tangent (b) versus frequency of four different samples.



Figure S8. The cyclical variation of magnetic moments of the individual Fe₃O₄ nanosphere under an alternating magnetic field obtained from computational micromagnetic simulation (a, b and c). The size of single Fe₃O₄ nanosphere is 250nm. The frequency of external magnetic field is 2 GHz.

MXene-based composites,	Fe ₃ O ₄ -base	ed composi	ites and	the as-pre	pared M/F
composites.					
Absorber	RL _{min}	Matching	EAB	Thickness	refs
	(dB)	frequency	(GHz)	(mm)	
		(GHz)			
MXene/amorphous	-48.4	11.6	2.8	1.85	S1
carbon/TiO ₂					
MXene/ZnO	-26.3	17.4	1.4	4	S2
MXene/Ni _{0.5} Zn _{0.5} Fe ₂ O ₄	-42.5	13.5	3	6.5	S3
MXene/PVB/Ba ₃ Co ₂ Fe ₂₄ O ₄	-46.3	5.8	1.6	2.8	S4
MXene/Ni-modified	-18.2	16.2	6.3	1.5	S5
MXene/Co ₃ O ₄	-34.5	14	6.3	2.0	S6

Table S1. Comparison of microwave absorption performance among the reported

MXene/Fe ₃ O ₄	-50.6	13.5	4.67	2	this work
Fe ₃ O ₄ /C	-40	16	-	1.5	S14
Fe ₃ O ₄	-53.0	2.2	<2	4	S13
MXene/carbonyl iron	-15.52	12.8	8.16	1	S12
MXene/Ni chain	-49.9	11.9	2.1	1.75	S11
MXene/Fe ₃ O ₄ /PANI	-40.3	15.3	5.2	1.9	S10
MXene/TiO ₂ /MoS ₂	~ -16	~ 9.8	2.6	2.5	S9
MXene/CoFe	-36.29	8.56	2.64	2.2	S8
MXene/FeCo	-17.86	-	8.8	1.6	S7

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The basic principle of geometric phase analysis is as follows:

For a perfect crystal, a HRTEM can be described as a Fourier series:

$$I(r) = \sum_{g} H_{g} e^{2\pi i g \cdot r} \quad (1)$$

Where I(r) is the image intensity at the position r, g is the periodicities of the Bragg reflections, the Fourier coefficients H_g can be described as:

$$H(g) = A_g e^{ip_g} (2)$$

Where A_g is the amplitude of the set of sinusoidal lattice fringes g, P_g is the lateral position of the fringes in the original image.

In real image conditions, the Fourier coefficient H_g has conjugate symmetry. Image strength can be expressed as the following real number function:

$$I(r) = A_0 + \sum_{g > 0} 2A_g \cos(2\pi g \cdot r + p_g) (3)$$

When processing the actually captured high-resolution image, the lattice image is subjected to fast Fourier transform processing to obtain an inverted space bitmap. A

specific \pm g direction lattice is selected by a mask to obtain a specific direction stripe information, and then an inverse Fourier transform is performed to obtain a lattice fringe B_g(r) in the specific direction:

$$B_g(r) = 2A_g \cos\left(2\pi g \cdot r + p_g\right)$$
(4)

In order to describe the lattice changes caused by distortion and defects in the material, the amplitude and phase of the lattice fringes should be expressed by the functions $A_g(r)$ and $P_g(r)$ for the position π , which should be written as:

$$B_g(r) = 2A_g(r)\cos\left(2\pi g \cdot r + p_g(r)\right)$$
(5)