# Supporting Information for: Transient X-Ray Absorption Spectral Fingerprints of the $S_{1}$ Dark State in Uracil 

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## Supplementary computational details

By using a utility script of the Columbus package, ${ }^{1-3}$ linear interpolations on Pulay's natural internal coordinates ${ }^{4}$ were performed between four anchor structures $\mathrm{i} \rightarrow \mathrm{ii} \rightarrow \mathrm{iii} \rightarrow \mathrm{iv}\left[\mathrm{S}_{0}\right.$ and $\mathrm{S}_{2}$ minima, and $\mathrm{S}_{2} / \mathrm{S}_{1}$ and $\mathrm{S}_{1} / \mathrm{S}_{0}$ conical intersections (CoIns), respectively], optimized at the state-averaged (SA) complete-active-space self-consistent field (CASSCF) level ${ }^{5,6}$ over 3 states with the $6-31 G^{*}$ basis set, taken from Nachtigallová et al. ${ }^{7,8}$ The natural internal coordinates are combinations of bond stretching, bending, rocking, torsions, out-of-plane wagging, etc., and are especially useful to study the deformation of ring compounds. For uracil (12 atoms), there are 30 such internal coordinates (Table S2). About 100 geometries are interpolated between geometries i-iv to guarantee the smoothness of potential energy surfaces (PESs) and spectra. Valence and core excited states were respectively computed at the SA-CASSCF and SA-RASSCF levels with the aug-cc-pVDZ basis set ${ }^{9}$ by using the Molpro program ${ }^{10,11}$ (RASSCF, restricted-active-space self-consistent field). Here a larger basis set was used since the $6-31 \mathrm{G}^{*}$ basis set is not suitable for core hole calculations, as it does not adequately describe the short-range asymptotic behavior of the core hole states. ${ }^{12}$ This basis set change has negligible influence of the valence state PESs. For valence-state calculation of interpolated structures, initial guess was read from the adjacent geometry to keep the consistency of the active space and smoothness of the PESs. At each geometry, vertical core excitation was always performed after obtaining the converged valence state wavefunction. ${ }^{13}$

The steady-state XAS spectra were first computed. Calibration of core-excited state energies were done by comparing with gas-phase experiment of uracil by Feyer et al. ${ }^{14}$ In the C , N , and O K-edge, energies of the simulated core-excited states were uniformly shifted by $-4.0,-5.1$, and -3.5 eV to match experiment. ${ }^{14}$ All stick spectra were convoluted by Gaussian line shape with fixed half-width-at-half-maximum (HWHM) of 0.3 eV . A broader line width than core hole lifetime ( $\mathrm{C} 1 \mathrm{~s}, 0.06 \mathrm{eV}$; N1s, 0.09 eV ; O1s, $0.13 \mathrm{eV}^{15}$ ) was used in order to represent the vibrionic coupling phenomenologically and to get better agreement with the
steady-state XAS experiment. ${ }^{14}$
The spectra and natural transition orbital (NTO) analysis ${ }^{16,17}$ were simulated by using our in-house code MCNOX. ${ }^{18}$ The NTOs provide a compact picture of a multi-electron state transition into one or a few pairs of particle to hole orbital transitions. For all analysis reported in this work, there is only one dominant pair, showing single-electron feature of the transitions. All graphical molecular orbitals (MOs) were generated using the Gabedit package ${ }^{19}$ at contour threshold of 0.08 .

Additional test calculations were performed to justify the optimized $\mathrm{S}_{1}$ minimum (min $\mathrm{S}_{1}$, labeled as structure v). ${ }^{7}$ Following the same procedure, valence PESs were also computed at interpolated paths connecting $\mathrm{i}-\mathrm{v}$ and $\mathrm{v}-\mathrm{ii}$. As a point test, RASSCF core excitation calculations were carried out only on geometry v and only at the O1s edge. Structures were superimposed to illustrate the structural changes, with graphics generated by VMD. ${ }^{20}$


Figure S1. (a) Optimized minima and minimum-energy conical intersections. Top, top view. Selected bond lengths are labeled in $\AA$, where significant structural changes (around $\mathrm{C}_{5}$ ) are colored in magenta. Bottom, side view. Major deformations are indicated by arrows. (b) Superposition of the four anchor geometries (top, top view; bottom, side view). The deformation is global while the biggest changes happen in the deformation of the $\mathrm{C}_{5}-\mathrm{H}_{5}$ bond with respect to the molecular plane. (c) Illustration for the i-ii, ii-iii, iii-iv interpolations. Major deformations are indicated by arrows.


Figure S2. Evolution of valence excitations. (a) Recapture of potential energy surface of $\mathrm{S}_{0}-\mathrm{S}_{2}$ from Figure 1. (b-c) Evolution of state character as interpreted by NTOs from $\mathrm{S}_{0}$. The two reaction coordinates are defined in panel a according to state character. Note the state index change near geometry iii. Contour isovalue $=0.08$ is used.


Figure S3. Computed vertical potential energy surfaces for 20 lowest $N_{1} 1$ s and $N_{3} 1$ s core excited states along the linear-interpolated internal coordinates. The lowest state of each set is labeled by $e_{1}$. All energies are uniformly shifted by -4.45 eV according to calibration to steady-state XAS experiment. ${ }^{14}$


Figure S4. Core excitations from a valence-excited state. (a) Physical picture. Top: The UV pump initiates a valence excitation and leaves a hole. Bottom: The X-ray probe at time $\tau$ generates multiple core excitations, amongst the lowest core excitation $\left(e_{1}\right)$ fills the valence hole (golden arrow). Gray dashed arrow denotes an ordinary higher excitation. (b-c) Computed NTOs ${ }^{16,17}$ of uracil. Top: NTOs of valence transitions from $\mathrm{S}_{0}$ to (b) $\pi \pi^{*}$ state $\mathrm{S}_{2}$ and (c) $n \pi^{*}$ state $\mathrm{S}_{1}$. Bottom: NTOs of core excitations from (b) $\mathrm{S}_{2}$ and (c) $\mathrm{S}_{1}$ to the lowest O1s core excitation state from $\mathrm{O}_{8}, e_{1}\left(\mathrm{O}_{8}\right)$. Calculations are on the ground state geometry (i.e., $\tau=0$ ).


Figure S5. Steady-state XAS spectra. Simulated (a) C1s, (b) N1s, and (c) O1s XAS spectra of uracil at the ground state geometry by RASSCF compared with gas-phase experiment by Fayer et al. ${ }^{14}$ Theoretical spectra are uniformly shifted by $-4.0,-4.45$, and -3.5 eV to match experiment. Each major peak is assigned as transition from 1s orbital to the lowest core excited state $\left(e_{1}\right)$ of the manifold, and is interpreted by natural transition orbitals ${ }^{16,17}$ as single-electron "hole" $\rightarrow$ "particle" transition within a NTO pair. The particle orbital is shown in the inset. Corresponding 1s hole orbital is always localized on the core-excited atom (the atom is labeled).


Figure S6. (a) Interpretation of O1s absorptions from dark state $S_{1}$ at geometries i, ii, and iii. Top, total spectra (top, thick green lines); bottom, atom-specific contributions. Major peaks are interpreted by NTOs with final orbitals depicted. Transitions to the $\mathrm{O}_{8}$ lone pair orbitals are especially shaded with red color. (b) Schematic illustration of TXAS and the O1s core excitation from $S_{1}$. The two lowest core excited states, $e_{1}\left(O_{8}\right)$ and $e_{1}\left(O_{7}\right)$ have large oscillator strengths.


Figure S7. Same as Figure S6a to interpret the N1s absorptions from dark state $\mathrm{S}_{1}$ at geometries i-iii.


Figure S8. Same as Figure S6a to interpret the C1s absorptions from dark state $\mathrm{S}_{1}$ at geometries i-iii.


Figure S9. Same as Figure S6a to interpret the O1s absorptions from bright $\pi \pi^{*}$ state $\mathrm{S}_{2}$ at geometries i-iii. Transitions filling the valence hole are especially shaded with blue color.


Figure S10. Same as Figure S6a to interpret the N1s absorptions from bright $\pi \pi^{*}$ state $\mathrm{S}_{2}$ at geometries i-iii. Transitions filling the valence hole are especially shaded with blue color.


Figure S11. Same as Figure S6a to interpret the C1s absorptions from bright $\pi \pi^{*}$ state $\mathrm{S}_{2}$ at geometries i-iii. Transitions filling the valence hole are especially shaded with blue color.


Figure S12. Same as Figure S6a to interpret the O1s absorptions from $\mathrm{S}_{0}$ state at geometries i-iv. At geometry iv, transition to the $2 \mathrm{p}\left(\mathrm{C}_{5}\right)$ orbital is especially shaded with gray color.


Figure S13. Same as Figure S6a to interpret the N1s absorptions from $\mathrm{S}_{0}$ state at geometries i-iv. At geometry iv, transitions to the $2 \mathrm{p}\left(\mathrm{C}_{5}\right)$ orbitals are especially shaded with gray color.


Figure S14. Same as Figure S6a to interpret the C1s absorptions from $\mathrm{S}_{0}$ state at geometries i-iv. At geometry iv, transitions to the $2 \mathrm{p}\left(\mathrm{C}_{5}\right)$ orbital are especially shaded with gray color.


Figure S15. Same as Figure S6a to interpret the (a) C1s, (b) N1s and (c) O1s absorptions from $S_{1}$ state at geometry iv.


Figure S16. Additional justification for the $S_{1}$ minima (min $S_{1}$, denoted as geometry v). (a) Optimized structure. Top, top view; bottom, side view. Selected bond lengths are labeled in $\AA$. (b) Schematic illustration of the structural difference of $\min S_{0}$, min $S_{1}$ and $\min S_{2}$ in terms of the root-mean-squared distances (rmsd). (c) Superimposed structures of min $S_{1}$ with min $S_{0}$ (left) or min $S_{2}$ (right). Top, top view; bottom, side view. (d) CASSCF valence PESs along another linear interpolated internal coordinates (LIIC) connecting min $S_{0}$ and min $S_{2}$ via min $S_{1}$. (f) Comparison of simulated O1s absorptions from dark state $S_{1}$ at the three minima. Top, total spectra (thick lines); bottom, atom-specific contributions from $\mathrm{O}_{8}$ and $\mathrm{O}_{7}$ (thin lines). The three total spectra are similar, which shows that excited state absorptions from the dark $n \pi^{*}$ state $S_{1}$ is less sensitive to structures.


Figure S17. CASSCF $(10,8)$ active space. Active space for uracil at the optimized ground state ( $\min \mathrm{S}_{0}$ ) at the CASSCF level (10 electrons in 8 orbitals) averaged over 3 lowest states $\mathrm{S}_{0}-\mathrm{S}_{2}$.

Table S1. Cartesian coordinates (aligned, in $\AA$ ) of anchor structures used.

| i | $\left(\min \mathrm{S}_{0}\right)$ |  |  | ii | $\left(\mathrm{min} \mathrm{S}_{2}\right)$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| C | -0.880497 | -0.001460 | -0.920817 | C | -0.871166 | -0.005059 | -0.914306 |
| C | 0.228394 | -0.001460 | 1.305006 | C | 0.282036 | -0.000618 | 1.232143 |
| C | 1.514814 | 0.004984 | 0.608334 | C | 1.519029 | -0.078069 | 0.650061 |
| C | 1.535060 | 0.007608 | -0.736754 | C | 1.587378 | 0.027097 | -0.830498 |
| N | -0.880497 | -0.001460 | 0.453828 | N | -0.840575 | 0.018865 | 0.473339 |
| N | 0.375309 | 0.005359 | -1.481593 | N | 0.363057 | 0.068080 | -1.482237 |
| O | -1.882652 | -0.006002 | -1.577641 | O | -1.904816 | -0.051370 | -1.524876 |
| O | 0.091176 | -0.006405 | 2.497617 | O | -0.033837 | 0.022239 | 2.544352 |
| H | -1.779645 | -0.009010 | 0.885316 | H | -1.722118 | 0.056455 | 0.939648 |
| H | 2.411511 | 0.006173 | 1.192763 | H | 2.406213 | -0.139102 | 1.242198 |
| H | 2.445087 | 0.011685 | -1.302851 | H | 2.366382 | 0.592757 | -1.312882 |
| H | 0.404548 | 0.005077 | -2.475286 | H | 0.344754 | 0.126085 | -2.477125 |
| iii | $\left(\mathrm{CoIn} \mathrm{S}_{2} / \mathrm{S}_{1}\right.$ ) |  |  | iv | $\left(\mathrm{CoIn} \mathrm{S} \mathrm{S}_{1} / \mathrm{S}_{0}\right)$ |  |  |
| C | -0.866570 | -0.002161 | -0.856336 | C | -0.876830 | -0.032274 | -0.890777 |
| C | 0.281459 | 0.049116 | 1.270136 | C | 0.260095 | -0.006308 | 1.325951 |
| C | 1.497393 | -0.126604 | 0.638339 | C | 1.483556 | -0.381581 | 0.569033 |
| C | 1.496880 | 0.228724 | -0.818814 | C | 1.480099 | 0.244513 | -0.735659 |
| N | -0.789089 | 0.482805 | 0.423286 | N | -0.881208 | 0.073641 | 0.463095 |
| N | 0.398806 | -0.176331 | -1.464873 | N | 0.402979 | 0.101935 | -1.493465 |
| O | -1.867453 | -0.200016 | -1.482408 | O | -1.846732 | -0.119441 | -1.590055 |
| O | -0.050320 | -0.254368 | 2.438650 | O | 0.079147 | 0.120681 | 2.499856 |
| H | -1.676263 | 0.580480 | 0.870603 | H | -1.781977 | 0.163051 | 0.885102 |
| H | 2.056878 | -1.009866 | 0.901164 | H | 1.359404 | -1.452714 | 0.387921 |
| H | 2.358797 | 0.505085 | -1.391471 | H | 2.295373 | 0.801784 | -1.167509 |
| H | 0.359028 | -0.319060 | -2.455497 | H | 0.373389 | 0.352904 | -2.462873 |
| v | $\left(\operatorname{min~} \mathrm{S}_{1}\right.$ ) |  |  |  |  |  |  |
| C | -0.881535 | -0.001501 | -0.918584 |  |  |  |  |
| C | 0.276456 | -0.001371 | 1.219919 |  |  |  |  |
| C | 1.511011 | 0.005005 | 0.638855 |  |  |  |  |
| C | 1.573661 | 0.008277 | -0.767696 |  |  |  |  |
| N | -0.890662 | -0.004492 | 0.459251 |  |  |  |  |
| N | 0.364974 | 0.004794 | -1.484217 |  |  |  |  |
| O | -1.896018 | -0.004340 | -1.560426 |  |  |  |  |
| O | 0.043220 | -0.005207 | 2.560877 |  |  |  |  |
| H | -1.785967 | -0.009166 | 0.891655 |  |  |  |  |
| H | 2.400459 | 0.007324 | 1.235752 |  |  |  |  |
| H | 2.476629 | 0.013264 | -1.336159 |  |  |  |  |
| H | 0.370751 | 0.006975 | -2.477739 |  |  |  |  |

Table S2. Sample Columbus output (file icoordtyp) for 30 natural internal coordinates of uracil at geometry i. The atomic orders are consistent with Table S1.

| 1 | $\mathrm{C} 3-\mathrm{C} 2$ | STRE | $\mathrm{sc}=.9200$ |
| ---: | :--- | ---: | ---: |
| 2 | $\mathrm{C} 4=\mathrm{C} 3$ | STRE | $\mathrm{sc}=.8600$ |
| 3 | $\mathrm{~N} 5-\mathrm{C} 1$ | STRE | $\mathrm{sc}=.9000$ |
| 4 | $\mathrm{~N} 5-\mathrm{C} 2$ | STRE | $\mathrm{sc}=.9000$ |
| 5 | $\mathrm{~N} 6-\mathrm{C} 1$ | STRE | $\mathrm{sc}=.9000$ |
| 6 | $\mathrm{~N} 6-\mathrm{C} 4$ | STRE | $\mathrm{sc}=.9000$ |
| 7 | $\mathrm{O} 7=\mathrm{C} 1$ | STRE | $\mathrm{sc}=.8260$ |
| 8 | $\mathrm{O} 8=\mathrm{C} 2$ | STRE | $\mathrm{sc}=.8260$ |
| 9 | H9-N5 | STRE | $\mathrm{sc}=.9000$ |
| 10 | H10-C3 | STRE | $\mathrm{sc}=.9000$ |
| 11 | H11-C4 | STRE | $\mathrm{sc}=.9000$ |
| 12 | H12-N6 | STRE | $\mathrm{sc}=.9000$ |
| 13 | 6 -membered ring | BEND | $\mathrm{sc}=.7900$ |
| 14 | $6-$-membered ring | BEND | $\mathrm{sc}=.7900$ |
| 15 | $6-$ membered ring | BEND | $\mathrm{sc}=.7900$ |
| 16 | $6-$-membered ring | TORS | $\mathrm{sc}=.9600$ |
| 17 | $6-\mathrm{membered} \mathrm{ring}$ | TORS | $\mathrm{sc}=.9600$ |
| 18 | $6-$ membered ring | TORS | $\mathrm{sc}=.9600$ |
| 19 | ring sXY C1-O7 | ROCK | $\mathrm{sc}=.8000$ |
| 20 | ring sXY C1-O7 | OUT | $\mathrm{sc}=.7000$ |
| 21 | ring sXY C2-O8 | ROCK | $\mathrm{sc}=.8000$ |
| 22 | ring sXY C2-O8 | OUT | $\mathrm{sc}=.7000$ |
| 23 | ring sXY C3-H10 | ROCK | $\mathrm{sc}=.8000$ |
| 24 | ring sXY C3-H10 | OUT | $\mathrm{sc}=.7000$ |
| 25 | ring sXY C4-H11 | ROCK | $\mathrm{sc}=.8000$ |
| 26 | ring sXY C4-H11 | OUT | $\mathrm{sc}=.7000$ |
| 27 | ring sXY N5-H9 | ROCK | $\mathrm{sc}=.8000$ |
| 28 | ring sXY N5-H9 | OUT | $\mathrm{sc}=.7000$ |
| 29 | ring sXY N6-H12 | ROCK | $\mathrm{sc}=.8000$ |
| 30 | ring sXY N6-H12 | OUT | $\mathrm{sc}=.7000$ |
|  |  |  |  |

## Supplementary References and Notes

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