## Supporting Information

## Macromolecular [2]Rotaxanes Linked with Polystyrene: Nanoscale Film Morphology and Properties

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## Grazing Incidence X-Ray Scattering (GIXS) Data Analysis

The intensity of GIXS (IGIXs) from structures in a thin film can be expressed by the scattering formula derived recently: ${ }^{1-4}$

$$
I_{\mathrm{GIXS}}\left(\alpha_{\mathrm{f}}, 2 \theta_{\mathrm{f}}\right) \cong \frac{1}{16 \pi^{2}} \cdot \frac{1-e^{-2 \operatorname{Im}\left(q_{z}\right) t}}{2 \operatorname{Im}\left(q_{z}\right)} \cdot\left[\begin{array}{l}
\left|T_{\mathrm{i}} T_{\mathrm{f}}\right|^{2} I_{1}\left(q_{\|}, \operatorname{Re}\left(q_{1, z}\right)\right)+  \tag{1}\\
\left|T_{\mathrm{i}} R_{\mathrm{f}}\right|^{2} I_{1}\left(q_{\|}, \operatorname{Re}\left(q_{2, z}\right)\right)+ \\
\left|T_{\mathrm{f}} R_{\mathrm{i}}\right|^{2} I_{1}\left(q_{\|}, \operatorname{Re}\left(q_{3, z}\right)\right)+ \\
\left|R_{\mathrm{i}} R_{\mathrm{f}}\right|^{2} I_{1}\left(q_{\|}, \operatorname{Re}\left(q_{4, z}\right)\right)
\end{array}\right]
$$

where $\alpha_{\mathrm{f}}$ and $2 \theta_{\mathrm{f}}$ are the out-of-plane and in-plane exit angle of the out-going X-ray beam respectively, $\operatorname{Im}\left(q_{z}\right)=\left|\operatorname{Im}\left(\mathrm{k}_{z, \mathrm{f}}\right)\right|+\left|\operatorname{Im}\left(\mathrm{k}_{z, \mathrm{i}}\right)\right|, \operatorname{Re}(x)$ is the real part of $x, t$ is the film thickness, $R_{\mathrm{i}}$ and $T_{\mathrm{i}}$ are the reflected and transmitted amplitudes of the incoming X-ray beam respectively, and $R_{\mathrm{f}}$ and $T_{\mathrm{f}}$ are the reflected and transmitted amplitudes of the outgoing X-ray beam respectively. In addition, $q_{\|}=\sqrt{q_{x}^{2}+q_{y}^{2}}, q_{1, z}=\mathrm{k}_{z, \mathrm{f}}-\mathrm{k}_{z, \mathrm{i}}, q_{2, z}=-\mathrm{k}_{z, \mathrm{f}}-\mathrm{k}_{z, \mathrm{i}}, q_{3, z}=\mathrm{k}_{z, \mathrm{f}}+\mathrm{k}_{z, \mathrm{i}}$, and $q_{4, \mathrm{z}}=-\mathrm{k}_{z, \mathrm{f}}+\mathrm{k}_{z, \mathrm{i}} ;$ here, $\mathrm{k}_{z, \mathrm{i}}$ is the $z$-component of the wave vector of the incoming X-ray beam, which is given by $\mathrm{k}_{z, \mathrm{i}}=\mathrm{k}_{\mathrm{o}} \sqrt{n_{\mathrm{R}}^{2}-\cos ^{2} \alpha_{\mathrm{i}}}$, and $\mathrm{k}_{z, \mathrm{f}}$ is the $z$-component of the wave vector of the outgoing X-ray beam, which is given by $\mathrm{k}_{z, \mathrm{f}}=\mathrm{k}_{\mathrm{o}} \sqrt{n_{\mathrm{R}}^{2}-\cos ^{2} \alpha_{\mathrm{f}}}$, where $\mathrm{k}_{\mathrm{o}}=2 \pi / \lambda, \lambda$ is the wavelength of the X-ray beam, $n_{\mathrm{R}}$ is the refractive index of the film given by $n_{\mathrm{R}}=1-\delta+i \zeta$ with dispersion $\delta$ and absorption $\zeta$, and $\alpha_{\mathrm{i}}$ is the out-of-plane grazing incident angle of the incoming X-ray beam. $q_{x}, q_{y}$, and $q_{z}$ are the components of the scattering vector $\mathbf{q} \cdot I_{1}$ is the scattering intensity of the structure in the film, which can be calculated kinematically.

In eq $1, I_{1}$ is the scattered intensity from morphological structures in the film and, thus, can be expressed by the following equation: ${ }^{1-6}$

$$
\begin{equation*}
I_{1}(\mathbf{q})=P(\mathbf{q}) \cdot S(\mathbf{q}) \tag{2}
\end{equation*}
$$

where $P(\mathbf{q})$ is the form factor that describes the shape, size, and orientation of scatterers in a thin film, and $S(\mathbf{q})$ is the structure factor which provides information on the relative positions of the group of scatterers, such as the crystal lattice parameters, orientation, dimension, and symmetry in an ordered structure and the interdistance of domains.

The measured GIXS (i.e., GISAXS) patterns inform that lamellar structures were formed in nanoscale films of PVL-b-PS, PVL-rot-PS-F and PVL-rot-PS-M. To analyze the GISAXS data, those lamellar structures may be appropriately handled by using a three layer model that is composed of a dense layer, a less dense layer, and their interfacial layer (Figure S15). For the lamellar structure consisted of three layers, the form factor $P(\mathbf{q})$ can be expressed by the following equation:

$$
P(\mathbf{q})=\left[\begin{array}{l}
4 L_{x} L_{y} H_{\text {inner }}\left(\rho_{\text {inner }}-\rho_{\text {outer }}\right) \cdot \frac{\sin \left(q_{x} L_{x}\right)}{q_{x} L_{x}} \cdot \frac{\sin \left(q_{y} L_{y}\right)}{q_{y} L_{y}} \cdot \frac{\sin \left(q_{z} H_{\text {inner }}\right)}{q_{z} H_{\text {inner }}}+  \tag{3}\\
4 L_{x} L_{y} H_{\text {outer }}\left(\rho_{\text {outer }}-\rho_{\text {matrix }}\right) \cdot \frac{\sin \left(q_{x} L_{x}\right)}{q_{x} L_{x}} \cdot \frac{\sin \left(q_{y} L_{y}\right)}{q_{y} L_{y}} \cdot \frac{\sin \left(q_{z} H_{\text {outer }}\right)}{q_{z} H_{\text {outer }}}
\end{array}\right]^{2}
$$

where $L_{x}$ and $L_{y}$ are the length and width of lamellar structure respectively, $H_{\text {inner }}$ and $H_{o u t e r}$ are the inner and outer layer heights respectively, and $\rho_{\text {inner }}$ and $\rho_{\text {outer }}$ are the relative electron densities of the inner and outer layers respectively. $\rho_{\text {matrix }}$ is the relative electron density of the matrix layer. For the lamellar structure, $H_{\text {inner }}$ corresponds to the crystal layer thickness $\left(=d_{\mathrm{PVL}}\right)$ having a relatively electron density $\rho_{\text {inner }}\left(=\rho_{\text {PVL }}\right)$. The interfacial layer thickness $\left(=d_{i}\right)$ having a relatively electron density $\rho_{\text {outer }}\left(=\rho_{\mathrm{i}}\right)$ is estimated from $H_{\text {inner }}$ and $H_{\text {outer }}: d_{\mathrm{i}}=\left(H_{\text {outer }}-H_{\text {innerer }}\right) / 2$. The amorphous PS layer thickness $d_{\mathrm{PS}}$ having the relatively electron density $\rho_{\text {matrix }}\left(=\rho_{\mathrm{PS}}\right.$ ) is obtained from the long period $D_{\mathrm{L}}$ that is extracted from the structure factor $S(\mathbf{q}): d_{\mathrm{PS}}=\left(D_{L}-d_{\mathrm{PVL}}-2 d_{\mathrm{i}}\right)$.

For the form factor, all structural parameters are further assumed to follow a Gaussian distribution $G(A)$ :

$$
\begin{equation*}
G(A)=\frac{1}{\sqrt{2 \pi} \sigma_{A}} \exp \left[-\frac{(A-\bar{A})^{2}}{2 \sigma_{A}{ }^{2}}\right] \tag{4}
\end{equation*}
$$

where $A$ can be one of the parameters, $\bar{A}$ is the mean value, and $\sigma_{A}$ is the standard deviation of $A$ from $\bar{A}$.

For a paracrystalline lattice consisting of the three layers described above, the structure factor $S(\mathbf{q})$ (the so-called interference function or lattice factor) can be determined from the Fourier transform of a complete set of lattice points. ${ }^{1-4,7,8}$ In a paracrystal with distortion of the second kind, the positions of the lattice points can only be described with a positional distribution function (i.e., $g$-factor). In the simple case where the autocorrelation function of the crystal lattice is given by the convolution product of the distributions of the lattice points along three axes, and the distribution function is a Gaussian, $S(\mathbf{q})$ can be expressed by the following equation: ${ }^{7}$

$$
\begin{gather*}
S(\mathbf{q})=\prod_{k=1}^{3} Z_{k}(\mathbf{q})  \tag{5}\\
Z_{k}(\mathbf{q})=1+\frac{F_{k}(\mathbf{q})}{1-F_{k}(\mathbf{q})}+\frac{F_{k}^{*}(\mathbf{q})}{1-F_{k}^{*}(\mathbf{q})}  \tag{6}\\
F_{k}(\mathbf{q})=\left|F_{k}(\mathbf{q})\right| e^{-i \mathbf{q} \cdot a_{k}}  \tag{7}\\
\left|F_{k}(\mathbf{q})\right|=\exp \left[-\frac{1}{2}\left(q_{1}^{2} g_{1}^{2}+q_{2}^{2} g_{2}^{2}+q_{3}^{2} g_{3}^{2}\right)\right] . \tag{8}
\end{gather*}
$$

Here $g_{1}, g_{2}$, and $g_{3}(=g)$ are the components of the $g$-factor defined as

$$
\begin{align*}
g_{1} & =\Delta \mathbf{a}_{1} / \mathbf{a}_{1}  \tag{9a}\\
g_{2} & =\Delta \mathbf{a}_{2} / \mathbf{a}_{2}  \tag{9b}\\
g_{3} & =\Delta \mathbf{a}_{3} / \mathbf{a}_{3} \tag{9c}
\end{align*}
$$

where $\mathbf{a}_{\mathbf{k}}$ is the component of the fundamental vector $\mathbf{a}$ of the domain structure and $\Delta \mathbf{a}_{\mathbf{k}}$ is the displacement of the vector $\mathbf{a}_{\mathbf{k}}$. And $q_{1}, q_{2}$, and $q_{3}$ are the components of the scattering vector $\mathbf{q}$.

For the lamellar structure composed of three layers (Figure S15), the components of $\mathbf{q}$ in Equation 8 are defined by

$$
\begin{align*}
q_{1} & =\mathbf{a}_{1} \cdot \mathbf{q}_{x}=\left|d_{x} \times \mathbf{q}_{x}\right|  \tag{10a}\\
q_{2} & =\mathbf{a}_{2} \cdot \mathbf{q}_{y}=\left|d_{y} \times \mathbf{q}_{y}\right|  \tag{10b}\\
q_{3} & =\mathbf{a}_{3} \cdot \mathbf{q}_{z}=\left|L \times \mathbf{q}_{z}\right| \tag{10c}
\end{align*}
$$

where $d_{x}$ and $d_{y}$ are the lattice dimension parameters (i.e., $d$-spacing values) along the $q_{x}$ - and $q_{y}$ direction respectively, and $L$ is the long period along the $q_{z}$-direction.

Moreover, for a structure with a given orientation in a film, its fundamental vectors can be rotated and transformed by a rotation matrix. When the structure of the film is randomly oriented in the plane of the film but uniaxially oriented out of plane, the peak position vector $\mathbf{q}_{\mathbf{c}}$ of a certain reciprocal lattice point $\mathbf{c}^{*}$ in the sample reciprocal lattice is given by

$$
\begin{align*}
\mathbf{q}_{\mathbf{c}} & =\mathbf{R} \cdot \mathbf{c}^{*} \\
& \equiv\left(q_{\mathrm{c}, \mathrm{x}}, q_{\mathrm{c}, \mathrm{y}}, q_{\mathrm{c}, \mathrm{z}}\right) \tag{11}
\end{align*}
$$

where $\mathbf{R}$ is a $3 \times 3$ matrix to decide the preferred orientation of the structure in the film, and $q_{\mathrm{c}, \mathrm{x}}$, $q_{c, y}$, and $q_{c, z}$ are the $x, y, z$ components of the peak position vector $\mathbf{q}_{\mathrm{c}}$, respectively. Using Equation 11, every peak position can be obtained. Because of cylindrical symmetry, the Debye-Scherrer ring composed of the in-plane randomly oriented $\mathbf{c}^{*}$ cuts an Ewald sphere at two positions in its top hemisphere: $q_{\|}=q_{\mathrm{c}, \mid} \equiv \pm \sqrt{q_{\mathrm{c}, \mathrm{x}}^{2}+q_{\mathrm{c}, \mathrm{y}}^{2}}$ with $q_{\mathrm{z}}=q_{\mathrm{c}, \mathrm{z}}$. Thus diffraction patterns with cylindrical symmetry are easily calculated in the $\mathbf{q}$-space. It is then convenient to determine the preferred orientation of known structures and further to analyze anisotropic X-ray scattering patterns. However, since $\mathbf{q}$-space is distorted in GIXS by refraction and reflection effects, the relation between the detector plane expressed as the Cartesian coordinate defined by two perpendicular axes (i.e., by $2 \theta_{\mathrm{f}}$ and $\alpha_{\mathrm{f}}$ ) and the reciprocal lattice points is needed. The two wave vectors $\mathrm{k}_{\mathrm{z}, \mathrm{i}}$ and $\mathrm{k}_{\mathrm{z}, \mathrm{f}}$ are corrected for refraction as $\mathrm{k}_{\mathrm{z}, \mathrm{i}}=\mathrm{k}_{\mathrm{o}} \sqrt{n_{\mathrm{R}}^{2}-\cos ^{2} \alpha_{\mathrm{i}}}$ and $\mathrm{k}_{\mathrm{z}, \mathrm{f}}=\mathrm{k}_{\mathrm{o}} \sqrt{n_{\mathrm{R}}^{2}-\cos ^{2} \alpha_{\mathrm{f}}}$ respectively. Therefore, the two sets of diffractions that result from the incoming and outgoing Xray beams, and denoted by $q_{1}$ and $q_{3}$ respectively, are given at the exit angles by the following expression:

$$
\begin{equation*}
\alpha_{\mathrm{f}}=\arccos \left(\sqrt{n_{\mathrm{R}}^{2}-\left(\frac{q_{\mathrm{c}, \mathrm{z}}}{\mathrm{k}_{\mathrm{o}}} \pm \sqrt{n_{\mathrm{R}}^{2}-\cos ^{2} \alpha_{\mathrm{i}}}\right)^{2}}\right) \tag{12}
\end{equation*}
$$

where $q_{\mathrm{c}, \mathrm{Z}} / \mathrm{k}_{\mathrm{o}}>\sqrt{n_{\mathrm{R}}^{2}-\cos ^{2} \alpha_{\mathrm{i}}}$. In Equation 12, the positive sign denotes diffractions produced by the outgoing X-ray beam, and the negative sign denotes diffractions produced by the incoming X-ray beam. The in-plane incidence angle $2 \theta_{\mathrm{i}}$ is usually zero, so the in-plane exit angle $2 \theta_{\mathrm{f}}$ can be expressed as follow:

$$
\begin{equation*}
2 \theta_{\mathrm{f}}=\arccos \left[\frac{\cos ^{2} \alpha_{\mathrm{i}}+\cos ^{2} \alpha_{\mathrm{f}}-\left(\frac{q_{\mathrm{c}, \mid}}{\mathrm{k}_{\mathrm{o}}}\right)^{2}}{2 \cos \alpha_{\mathrm{i}} \cos \alpha_{\mathrm{f}}}\right] \tag{13}
\end{equation*}
$$

Therefore, diffraction spots detected on the detector plane in GIXS measurements can be directly compared to those derived using Equations 11-13 from an appropriate model and thus analyzed in terms of the model.

To obtain information on the orientation of the paracrystal lattice of the phase separated micro domain structures from GIXS data, the distribution of the orientation vector $\boldsymbol{n}$ is given by a function $D(\varphi)$, where $\varphi$ is the polar angle between the $\boldsymbol{n}$ vector and the out-of-plane of the film (see an example in Figure S15); for example, $\varphi$ is zero when the $\boldsymbol{n}$ vector in the film is oriented normal to the film plane. To calculate the 2D GIXS patterns, $D(\varphi)$ should be represented by a numerical function. In relation to the distribution of the lattice orientation, $D(\varphi)$ can generally be considered as a Gaussian distribution:

$$
\begin{equation*}
D(\varphi)=\frac{1}{\sqrt{2 \pi} \sigma_{\varphi}} \exp \left[-\frac{(\varphi-\bar{\varphi})^{2}}{2 \sigma_{\varphi}^{2}}\right] \tag{14}
\end{equation*}
$$

where $\bar{\varphi}$ and $\sigma_{\varphi}$ are the mean angle and standard deviation of $\varphi$ from $\bar{\varphi}$, respectively. The observed scattering intensity $I_{\text {GIXS }, \varphi}(\mathbf{q})$ is obtained by integrating $I_{\text {GIXS }}(\mathbf{q})$ over possible orientations of the lattice:

$$
\begin{equation*}
I_{\mathrm{GIXS}, \varphi}(\mathbf{q})=\int_{-\pi}^{\pi} I_{\mathrm{GIXS}}(\mathbf{q}) D(\varphi) \mathrm{d} \varphi \tag{15}
\end{equation*}
$$

The second order orientation factor $O_{\mathrm{s}}$ can be defined as the following equation: ${ }^{7-9}$

$$
\begin{equation*}
O_{\mathrm{s}}=\int D(\varphi) \frac{\left(3 \cos ^{2} \varphi-1\right)}{2} \mathrm{~d} \varphi \tag{21}
\end{equation*}
$$

When $D(\varphi)$ is strongly peaked around $\varphi=0^{\circ}$ (i.e. vertical alignment), $\cos \varphi=1$ and $O_{\mathrm{s}}=1$. On the other hand, when $\varphi=90^{\circ}, \cos \varphi=0$ and $O_{s}=-0.5$. If the orientation is entirely random, $<\cos ^{2} \varphi>=$ $1 / 3$ and $O_{\mathrm{s}}=0$. Thus, $O_{\mathrm{s}}$ is a measure of the orientation of nanostructures.

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Figure S1. Thermograms of PVL homopolymer ( $5190 M_{\mathrm{n}, \mathrm{NMR}} ; 3850 M_{\mathrm{n}, \mathrm{NMR}}$ for only PVL part excluded end groups): (a) TGA thermogram; (b) DSC thermograms. All thermograms were measured at a rate of $10.0^{\circ} \mathrm{C} \mathrm{min}^{-1}$ under nitrogen atmosphere.


Figure S2. DSC thermograms of PVL-b-PS measured at two different rates: (a) $5.0^{\circ} \mathrm{C} \mathrm{min}^{-1}$; (b) $2.0^{\circ} \mathrm{C} \mathrm{min}^{-1}$. All measurements were conducted under nitrogen atmosphere.


Figure S3. DSC thermograms of PVL-rot-PS-F measured at two different rates: (a) $5.0^{\circ} \mathrm{C} \mathrm{min}^{-1}$; (b) $2.0^{\circ} \mathrm{C} \mathrm{min}{ }^{-1}$. All measurements were conducted under nitrogen atmosphere.


Figure S4. DSC thermograms of PVL-rot-PS-M measured at $5.0{ }^{\circ} \mathrm{C} \mathrm{min}{ }^{-1}$ under nitrogen atmosphere.


Figure S5. Representative GISAXS data of the THF-annealed films (110-130 nm thick) of PVL-$b$-PS, PVL-rot-PS-F, and PVL-rot-PS-M measured with a sample-to-detector distance (SDD) of 2910 mm using a synchrotron X-ray beam $(\lambda=0.12096 \mathrm{~nm})$. PVL-b-PS: (a) 2D scattering image measured with of $\alpha_{\mathrm{i}}=0.133^{\circ}$; (d) 2D image reconstructed from the structural parameters in Table 3 using the GIXS formula; (g) out-of-plane scattering profile extracted along the meridian line at $2 \theta_{\mathrm{f}}=0.123^{\circ} ;(\mathrm{j})$ in-plane scattering profile extracted along the equatorial line at $\alpha_{\mathrm{f}}=0.198^{\circ}$; (m) azimuthal scattering profile of the first-order lamellar peak at $0.371^{\circ}$. PVL-rot-PS-F: (b) 2D scattering image measured with of $\alpha_{\mathrm{i}}=0.134^{\circ}$; (e) 2D image reconstructed from the structural parameters in Table 3 using the GIXS formula; (h) out-of-plane scattering profile extracted along
the meridian line at $2 \theta_{\mathrm{f}}=0.123^{\circ} ;(\mathrm{k})$ in-plane scattering profile extracted along the equatorial line at $\alpha_{\mathrm{f}}=0.198^{\circ} ;(\mathrm{n})$ azimuthal scattering profile of the first-order lamellar peak at $0.357^{\circ}$. PVL-rot-PS-M: (c) 2D scattering image measured with of $\alpha_{\mathrm{i}}=0.131^{\circ}$; (f) 2D image reconstructed from the structural parameters in Table 3 using the GIXS formula; (i) out-of-plane scattering profile extracted along the meridian line at $2 \theta_{\mathrm{f}}=0.107^{\circ}$; (l) in-plane scattering profile extracted along the equatorial line at $\alpha_{\mathrm{f}}=0.238^{\circ}$; (o) azimuthal scattering profile of the firstorder lamellar peak at $0.520^{\circ}$. In $(g-1)$, the symbols are the measured data and the solid lines were obtained by fitting the data using the GIXS formula of lamellar structure model. In ( $\mathrm{m}-\mathrm{o}$ ), the symbols are the measured data and the sold lines were obtained by fitting the data using the Gaussian functions.


Figure S6. Representative GIWAXS patterns of the THF-annealed films (110-130 nm thick) of PVL-b-PS, PVL-rot-PS-F, and PVL-rot-PS-M films measured with $\alpha_{\mathrm{i}}=0.174^{\circ}$ at $\mathrm{SDD}=209$ mm using a synchrotron X-ray beam $(\lambda=0.1209 \mathrm{~nm})$. PVL- $b$-PS: (a) 2D scattering image; (d) 1D scattering profile averaged quadrantly from the 2D image in (a); (g) azimuthal scattering profile of the $\{110\}$ peak. PVL-rot-PS-F: (b) 2D scattering image; (e) 1D scattering profile averaged quadrantly from the 2D image in (b); (h) azimuthal scattering profile of the $\{110\}$ peak. PVL-rot-PS-M: (c) 2D scattering image; 1D scattering profile averaged quadrantly from the 2D image in (c); (i) azimuthal scattering profile of the $\{110\}$ peak.

## Orientations

(Orthorhombic Crystal)



$\mu$ (degree)

Figure S7. Azimuthal scattering profiles and analysis results of the $\{110\}$ peaks of PVL crystals in Figure 5 and orientational crystal domains: (a) as-cast PVL-b-PS film; (b) as-cast PVL-rot-PS-F film; (c) as-cast PVL-rot-PS-M film.

## Orientations

(Orthorhombic Crystal)




Figure S8. Azimuthal scattering profiles and analysis results of the $\{110\}$ peaks of PVL crystals in Figure S6 and orientational crystal domains: (a) THF-annealed PVL-b-PS film; (b) THFannealed PVL-rot-PS-F film; (c) THF-annealed PVL-rot-PS-M film.
(a)


Figure S9. (a) Enlarged $\{110\}$ reflection profile of PVL-b-PS-F film (as-cast film) obtained from the 1D scattering profile in Figure 5d. The symbols represent the measured data; the blue curves represent the scattering profiles obtained by the deconvolution of the $\{110\}$ reflection profile; the red curve represents the sum of the deconvoluted scattering profiles in blue color. (b) Top views of orthorhombic lattices in three different rotational domains which were determined by the deconvolution analysis of the scattering profile in (a); $d_{110}$ is the $d$-spacing of $\{110\}$ reflection. (c) Deconvolution analysis results of the scattering profile in (a); $\omega_{i}$ and $\phi_{i}$ are the rotational angle and relative volume fraction of orthorhombic lattice domain $i$ respectively where the rotational axis is parallel to the $c$-axis of the lattice (which is also parallel to the out-of-plane of the film).


Figure S10. (a) Enlarged $\{110\}$ reflection profile of PVL-rot-PS-F film (as-cast film) obtained from the 1D scattering profile in Figure 5e. The symbols represent the measured data; the blue curves represent the scattering profiles obtained by the deconvolution of the $\{110\}$ reflection profile; the red curve represents the sum of the deconvoluted scattering profiles in blue color. (b) Top views of orthorhombic lattices in three different rotational domains which were determined by the deconvolution analysis of the scattering profile in (a); $d_{110}$ is the $d$-spacing of $\{110\}$ reflection. (c) Deconvolution analysis results of the scattering profile in (a); $\omega_{i}$ and $\phi_{i}$ are the rotational angle and relative volume fraction of orthorhombic lattice domain $i$ respectively where the rotational axis is parallel to the $c$-axis of the lattice (which is also parallel to the out-of-plane of the film).


Figure S11. (a) Enlarged $\{110\}$ reflection profile of PVL-rot-PS-M film (as-cast film) obtained from the 1D scattering profile in Figure 5f. The symbols represent the measured data; the blue curves represent the scattering profiles obtained by the deconvolution of the $\{110\}$ reflection profile; the red curve represents the sum of the deconvoluted scattering profiles in blue color. (b) Top views of orthorhombic lattices in three different rotational domains which were determined by the deconvolution analysis of the scattering profile in (a); $d_{110}$ is the $d$-spacing of $\{110\}$ reflection. (c) Deconvolution analysis results of the scattering profile in (a); $\omega_{i}$ and $\phi_{i}$ are the rotational angle and relative volume fraction of orthorhombic lattice domain $i$ respectively where the rotational axis is parallel to the $c$-axis of the lattice (which is also parallel to the out-of-plane of the film).


Figure S12. (a) Enlarged $\{110\}$ reflection profile of PVL-b-PS film (THF-annealed film) obtained from the 1D scattering profile in Figure S6d. The symbols represent the measured data; the blue curves represent the scattering profiles obtained by the deconvolution of the $\{110\}$ reflection profile; the red curve represents the sum of the deconvoluted scattering profiles in blue color. (b) Top views of orthorhombic lattices in three different rotational domains which were determined by the deconvolution analysis of the scattering profile in (a); $d_{110}$ is the $d$ spacing of $\{110\}$ reflection. (c) Deconvolution analysis results of the scattering profile in (a); $\omega_{i}$ and $\phi_{i}$ are the rotational angle and relative volume fraction of orthorhombic lattice domain $i$ respectively where the rotational axis is parallel to the $c$-axis of the lattice (which is also parallel to the out-of-plane of the film).


Figure S13. (a) Enlarged $\{110\}$ reflection profile of PVL-rot-PS-F film (THF-annealed film) obtained from the 1D scattering profile in Figure S6e. The symbols represent the measured data; the blue curves represent the scattering profiles obtained by the deconvolution of the $\{110\}$ reflection profile; the red curve represents the sum of the deconvoluted scattering profiles in blue color. (b) Top views of orthorhombic lattices in three different rotational domains which were determined by the deconvolution analysis of the scattering profile in (a); $d_{110}$ is the $d$ spacing of $\{110\}$ reflection. (c) Deconvolution analysis results of the scattering profile in (a); $\omega_{i}$ and $\phi_{i}$ are the rotational angle and relative volume fraction of orthorhombic lattice domain $i$ respectively where the rotational axis is parallel to the $c$-axis of the lattice (which is also parallel to the out-of-plane of the film).


Figure S14. (a) Enlarged $\{110\}$ reflection profile of PVL-rot-PS-M film (THF-annealed film) obtained from the 1D scattering profile in Figure S6f. The symbols represent the measured data; the blue curves represent the scattering profiles obtained by the deconvolution of the $\{110\}$ reflection profile; the red curve represents the sum of the deconvoluted scattering profiles in blue color. (b) Top views of orthorhombic lattices in three different rotational domains which were determined by the deconvolution analysis of the scattering profile in (a); $d_{110}$ is the $d$ spacing of $\{110\}$ reflection. (c) Deconvolution analysis results of the scattering profile in (a); $\omega_{i}$ and $\phi_{i}$ are the rotational angle and relative volume fraction of orthorhombic lattice domain $i$ respectively where the rotational axis is parallel to the $c$-axis of the lattice (which is also parallel to the out-of-plane of the film).


Figure S15. A lamellar structure model composed of three layers: (a) 3D representation of lamellar structure where $\boldsymbol{n}_{1}$ is the orientation vector of the structure and $\varphi_{1}$ is the polar angle between the $\boldsymbol{n}_{1}$ vector and the out-of-plane of the film; (b) 2D representation of lamellar structure. The inner and outer layers in (a) correspond to the dense and interfacial layers in (b) respectively; the blue colored layer in (b) corresponds to the less dense layer; (c) The electron density profile along the direction of layer stacks in the lamellar structure where $\rho_{\mathrm{PvL}}$ and $\rho_{\mathrm{PS}}$ are the electron densities of crystalline (dense) and amorphous (less dense) layers respectively. The dimension of the lamellar structure is defined by $D_{\mathrm{L}}$ (long period), $L_{x}, L_{y}, d \mathrm{PVL}\left(=H_{\text {inner }}\right), d_{\mathrm{i}}\left(=\left(H_{\text {outer }}-H_{\text {inner }}\right) / 2\right)$, and $d_{\mathrm{PS}}\left(=D_{\mathrm{L}}-H_{\text {outer }}\right)$.

