

Supporting Information

for

Long Carbon-Based Chains of Interstellar Medium Can Have a Triplet Ground State. Why Is This Important for Astrochemistry

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S1 Additional Computational Details

For the reader's convenience, a comprehensive list of the various quantities discussed in this work is presented in Table S1.

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Table S1: Definition of the various quantities utilized in the present paper.

zero-point energies of the singlet, triplet, cation, anion	$\varepsilon_S, \varepsilon_T, \varepsilon_c, \varepsilon_a$
non-corrected adiabatic singlet-triplet separation	$\Delta_{ad} \equiv \mathcal{E}_T(\mathbf{R}_T) - \mathcal{E}_S(\mathbf{R}_S)$
non-corrected vertical singlet-triplet separation at singlet optimum	$\Delta_S \equiv \mathcal{E}_T(\mathbf{R}_S) - \mathcal{E}_S(\mathbf{R}_S)$
non-corrected vertical singlet-triplet separation at triplet optimum	$\Delta_T \equiv \mathcal{E}_T(\mathbf{R}_T) - \mathcal{E}_S(\mathbf{R}_T)$
corrected adiabatic singlet-triplet separation, <i>etc</i>	$\Delta_{ad}^{corr} \equiv \mathcal{E}_T(\mathbf{R}_T) + \varepsilon_T - [\mathcal{E}_S(\mathbf{R}_S) + \varepsilon_S], \text{ etc} \quad (\text{S1})$
non-corrected adiabatic ionization energy wrt singlet	$IP_S^{ad,non} \equiv \mathcal{E}_c(\mathbf{R}_c) - \mathcal{E}_S(\mathbf{R}_S)$
non-corrected vertical ionization energy at singlet optimum	$IP_S^{vert,non} \equiv \mathcal{E}_c(\mathbf{R}_S) - \mathcal{E}_S(\mathbf{R}_S)$
non-corrected adiabatic ionization energy wrt triplet	$IP_T^{ad,non} \equiv \mathcal{E}_c(\mathbf{R}_c) - \mathcal{E}_T(\mathbf{R}_T)$
non-corrected vertical ionization energy at triplet optimum	$IP_T^{vert,non} \equiv \mathcal{E}_c(\mathbf{R}_T) - \mathcal{E}_T(\mathbf{R}_T)$
corrected vertical ionization energy at singlet optimum, <i>etc</i>	$IP_S^{vert,corr} \equiv \mathcal{E}_c(\mathbf{R}_c) + \varepsilon_c - [\mathcal{E}_S(\mathbf{R}_S) + \varepsilon_S], \text{ etc}$
non-corrected adiabatic electroaffinity wrt singlet	$EA_S^{ad,non} \equiv \mathcal{E}_S(\mathbf{R}_S) - \mathcal{E}_a(\mathbf{R}_a)$
non-corrected vertical electroaffinity at singlet optimum	$EA_S^{vert,non} \equiv \mathcal{E}_S(\mathbf{R}_S) - \mathcal{E}_a(\mathbf{R}_S)$
non-corrected adiabatic electroaffinity wrt triplet	$EA_T^{ad,non} \equiv \mathcal{E}_T(\mathbf{R}_T) - \mathcal{E}_a(\mathbf{R}_a)$
non-corrected vertical electroaffinity at triplet optimum	$EA_T^{vert,non} \equiv \mathcal{E}_T(\mathbf{R}_T) - \mathcal{E}_a(\mathbf{R}_T)$
corrected vertical electroaffinity energy at triplet optimum, <i>etc</i>	$EA_T^{vert,corr} \equiv \mathcal{E}_T(\mathbf{R}_T) + \varepsilon_t - [\mathcal{E}_a(\mathbf{R}_T) + \varepsilon_T], \text{ etc}$
cation's reorganization energy wrt singlet	$\lambda_S^c \equiv \mathcal{E}_c(\mathbf{R}_S) - \mathcal{E}_c(\mathbf{R}_c)$
cation's reorganization energy wrt triplet	$\lambda_T^c \equiv \mathcal{E}_c(\mathbf{R}_T) - \mathcal{E}_c(\mathbf{R}_c)$
anion's reorganization energy wrt singlet	$\lambda_S^a \equiv \mathcal{E}_a(\mathbf{R}_S) - \mathcal{E}_a(\mathbf{R}_a)$
anion's reorganization energy wrt triplet	$\lambda_T^a \equiv \mathcal{E}_a(\mathbf{R}_T) - \mathcal{E}_a(\mathbf{R}_a)$
non-corrected chemical hardness of the singlet isomer	$\eta_S^{vert,non} \equiv IP_S^{vert,non} - EA_S^{vert,non}$
non-corrected chemical hardness of the triplet isomer	$\eta_T^{vert,non} \equiv IP_T^{vert,non} - EA_T^{vert,non}$
non-corrected "combined" hardness of the singlet isomer	$\eta_S^{ad,non} \equiv IP_S^{ad,non} - EA_S^{ad,non}$
non-corrected "combined" hardness of the triplet isomer	$\eta_T^{ad,non} \equiv IP_T^{ad,non} - EA_T^{ad,non}$
corrected chemical hardness of the singlet/triplet isomer	$\eta_{S/T}^{vert,corr} \equiv IP_{S/T}^{vert,corr} - EA_{S/T}^{vert,corr}$
corrected "combined" hardness of the singlet/triplet isomer	$\eta_{S/T}^{ad,corr} \equiv IP_{S/T}^{ad,corr} - EA_{S/T}^{ad,corr}$

S2 Restricted Open-Shell *vs.* Unrestricted Calculations

As noted in Sec. 2 of the main text, spin contamination may in principle affect properties obtained *via* unrestricted (U...) methods. To be on the safe side, calculations for open shell cases reported in the main text were carried out using restricted open-shell (RO...) methods.

In fact, as illustrated by the examples presented here, the impact of spin contamination is insignificant for the specific cases envisaged. For the triplet HC₁₀N, UB3LYP/6-311++g(3df, 3pd) calculations yielded $\langle \mathbf{S}^2 \rangle = 2.1457$ before annihilation of the first spin contaminant and $\langle \mathbf{S}^2 \rangle = 2.0136$ after annihilation. This corresponds to spin values $S = 1.0478$ and $S = 1.00045$, respectively. In view of the negligible deviation from the exact value $S = 1$, it is not surprising that ROB3LYP- and UB3LYP-based properties do not notably differ from each other, as shown by the values collected in Tables S2 and S3.

Table S2: Bond lengths (d) of a triplet HC₁₀N linear chain obtained from restricted open shell and unrestricted DFT/B3LYP/6-311++g(3df, 3pd) calculations. The carbon atoms in the triplet HC₁₀N chain are numbered consecutively (C₁, C₂, C₃, ...) starting from the H atom to the N atom.

	HC ₁	C ₁ C ₂	C ₂ C ₃	C ₃ C ₄	C ₄ C ₅	C ₅ C ₆	C ₆ C ₇	C ₇ C ₈	C ₈ C ₉	C ₉ C ₁₀	C ₁₀ N
ROB3LYP	1.0619	1.2143	1.3342	1.2418	1.2983	1.2734	1.2636	1.3081	1.2333	1.3470	1.1630
UB3LYP	1.0620	1.2166	1.3318	1.2459	1.2967	1.2748	1.2661	1.3057	1.2370	1.3452	1.1643

Table S3: Wiberg bond indices ($\mathcal{N}$) of a triplet HC₁₀N linear chain obtained from restricted open shell and unrestricted DFT/B3LYP/6-311++g(3df, 3pd) calculations. The carbon atoms in the triplet HC₁₀N chain are numbered consecutively (C₁, C₂, C₃, ...) starting from the H atom to the N atom.

	HC ₁	C ₁ C ₂	C ₂ C ₃	C ₃ C ₄	C ₄ C ₅	C ₅ C ₆	C ₆ C ₇	C ₇ C ₈	C ₈ C ₉	C ₉ C ₁₀	C ₁₀ N
UB3LYP	0.9287	2.5102	1.3648	2.0392	1.5793	1.7464	1.8290	1.5154	2.1374	1.2889	2.6138
ROB3LYP	0.9282	2.5411	1.3512	2.0812	1.5654	1.7562	1.8494	1.4985	2.1792	1.2801	2.6315

S3 Properties of the Triplet HC_{11}N

For completeness, bond lengths and bond orders of the triplet HC_{11}N chain (*i.e.*, the metastable HC_{11}N isomer) are depicted in Figs. S1 and S2. These results are particularly important because they suggest that a triplet state enforces a cumulenic structure even in cases where all electrons can form pairs according to the standard rules of valence.

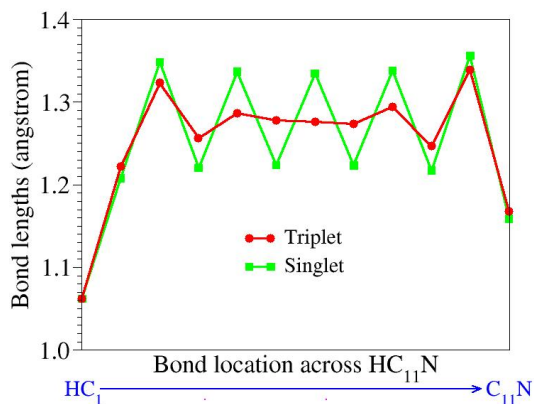


Figure S1: Bond lengths of triplet and singlet HC_{11}N chains. The carbon atoms in chains are numbered consecutively ($\text{C}_1, \text{C}_2, \text{C}_3, \dots$) starting from the H atom to the N atom. Notice that the triplet HC_{11}N chain is characterized by a cumulenic structure, in contrast to the polyynic character of the singlet HC_{11}N chain (*cf.* Fig. S2 and Fig. 2 in the main text).

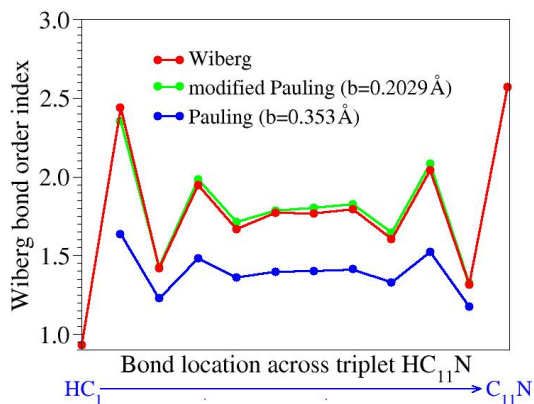


Figure S2: Wiberg bond indices ($\mathcal{N}$) of a triplet HC_{11}N linear chain. The carbon atoms in chains are numbered consecutively ($\text{C}_1, \text{C}_2, \text{C}_3, \dots$) starting from the H atom to the N atom. Notice that the triplet HC_{11}N chain is characterized by a cumulenic structure, in contrast to the polyynic character of the singlet HC_{11}N chain (*cf.* Fig. S1 and Fig. 2 in the main text).

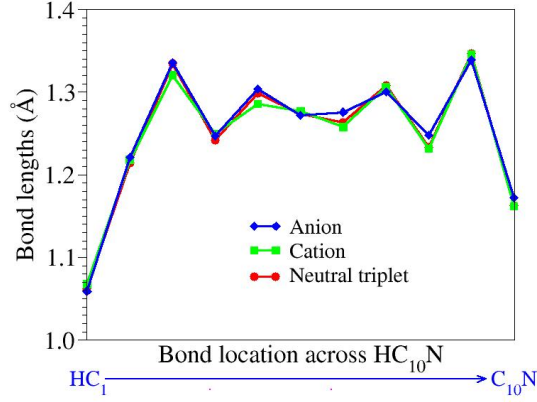


Figure S3: Bond lengths of various HC_{10}N linear chains: neutral triplet, cation, and anion (labels T, +, and - in the main text). The carbon atoms in chains are numbered consecutively ($\text{C}_1, \text{C}_2, \text{C}_3, \dots$) starting from the H atom to the N atom. The small differences in the geometries of these charge species (chain lengths $L_{n.t.} = 13.7390 \text{ \AA}$, $L_+ = 13.7217 \text{ \AA}$, $L_- = 13.7745 \text{ \AA}$) reflect themselves in the very small reorganization energies λ indicated in Tables 7 and 8 of the main text.

S4 Additional Data for Electric Dipole and Quadrupole Momentum of HC_{10}N Chains

To demonstrate the robustness of the estimates for electric dipole D and quadrupole momentum $\mathbf{Q}$ presented in Table 5 of the main text, values obtained by using various basis sets are collected in Table S4

S5 Additional Vibrational Spectra of HC_{10}N Chains

Fig. S4 depicts additional infrared spectra. All numerical data underlying this figure and the related Fig. 5 of the main text are collected in Tables S5 and S6 where the $\text{C}_{\infty v}$ -symmetries ($a_1 = \sigma_+$, $a_2 = \sigma_-$, $e_1 = \pi$) are indicated as obtained from GAUSSIAN 16¹ calculations.

Table S4: Electric dipoles $\mathbf{D}$ (debye) and quadrupole moments $\mathbf{Q}$ (debye-angstrom) of singlet ($S = 0$) and triplet ($S = 1$) HC_{10}N chains in field-independent basis ($D_x = D_y = Q_{i \neq j} = 0$, molecular z -axis from H (positively charged) to N (negatively charged)). Results obtained *via* DFT/B3LYP and CCSD(T) using the basis sets indicated in the second column. (Dipole oriented from H (positive) N (negative)).

Spin	Method	D_z	Q_{xx}	Q_{yy}	Q_{zz}
$S = 0$	B3LYP, 6-311++g(3df, 3pd)	-7.0103	-59.7616	-63.1138	-67.7443
	B3LYP, cc-pvTZ	-6.9045	-59.5597	-62.8507	-65.8263
	B3LYP, aug-cc-pvTZ	-6.9985	-63.0100	-59.7976	-67.6938
	CCSD(T), 6-311++g(3df, 3pd)	-7.0202	-64.4555	-61.1693	-67.0326
	CCSD(T), cc-pVTZ	-6.9619	-64.2501	-61.0084	-65.8328
	CCSD(T), aug-cc-pVTZ	-7.0078	-64.3279	-61.1934	-66.9847
$S = 1$	RO-B3LYP, 6-311++g(3df, 3pd)	-6.8902	-61.3902	-61.3902	-67.9786
	RO-B3LYP, cc-pvTZ	-6.7830	-61.1659	-61.1659	-66.1385
	RO-B3LYP, aug-cc-pVTZ	-6.8777	-61.3555	-61.3555	-67.9323
	RO-CCSD(T), 6-311++g(3df, 3pd)	-6.9117	-62.8235	-62.8235	-67.9001
	RO-CCSD(T), cc-pVTZ	-6.8525	-62.6438	-62.6438	-66.7272
	RO-CCSD(T), aug-cc-pVTZ	-6.8976	-62.7706	-62.7706	-67.8580

Table S5: Unscaled phonon frequencies of the molecular vibrations with significant infrared (IR) intensities obtained *via* DFT/ROB3LYP/6-311++g(3df, 3pd) calculations for the triplet HC_{10}N chain.

$\nu \text{ cm}^{-1}$	IR intensity (km/mol)	Symmetry
91.02	3.5134	π
91.02	3.5134	π
254.14	3.2010	π
254.14	3.2010	π
478.38	5.6058	π
478.38	5.6058	π
585.81	6.5847	π
585.81	6.5847	π
637.66	17.8412	π
637.66	17.8412	π
663.25	19.3101	π
663.25	19.3101	π
709.80	8.3849	σ_+
2075.25	22.2267	σ_+
2105.85	122.2943	σ_+
2258.32	169.4757	σ_+
3460.05	257.9971	σ_+

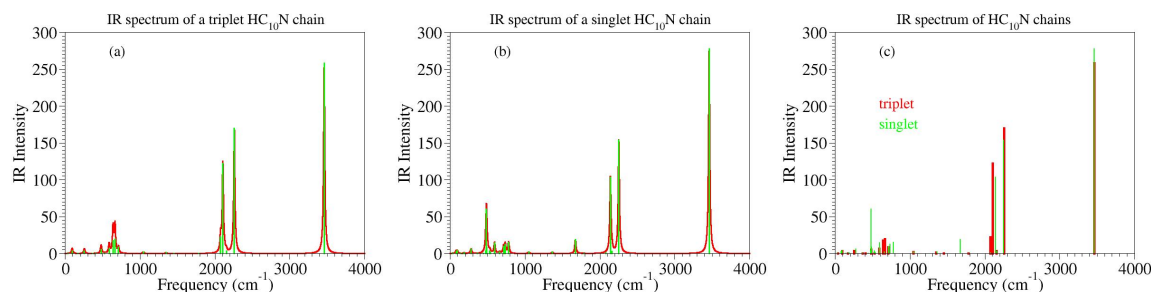


Figure S4: Additional infrared spectra of the triplet and singlet HC_{10}N isomers. The spectra depicted here were convoluted by using obtained by using Lorentzian convolutions of the spectral lines (green spikes) calculated *via* DFT/B3LYP/6-311++g(3df, 3pd) (halfwidth $hw = 10 \text{ cm}^{-1}$).

References

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Table S6: Unscaled phonon frequencies of the molecular vibrations with significant infrared (IR) intensities obtained *via* DFT/RB3LYP/6-311++g(3df, 3pd) calculations for the singlet HC₁₀N chain.

$\nu \text{ cm}^{-1}$	IR intensity (km/mol)	Symmetry
77.71	2.8889	?a
91.04	3.7212	?a
277.56	6.2471	?a
470.59	3.5720	?a
481.68	60.1046	?a
488.98	8.5384	?a
523.90	3.3852	?a
591.51	14.5577	?a
712.92	9.4541	σ_+
733.86	12.5985	?a
779.22	15.2197	?a
1671.86	18.7659	σ_+
2139.68	103.0968	σ_+
2252.12	153.6603	σ_+
3459.40	277.1618	σ_+