

# **Supporting Information**

## **New Theoretical Insights into the Crystal-Field Splitting and Transition Mechanism for Nd<sup>3+</sup> Doped Y<sub>3</sub>Al<sub>5</sub>O<sub>12</sub>**

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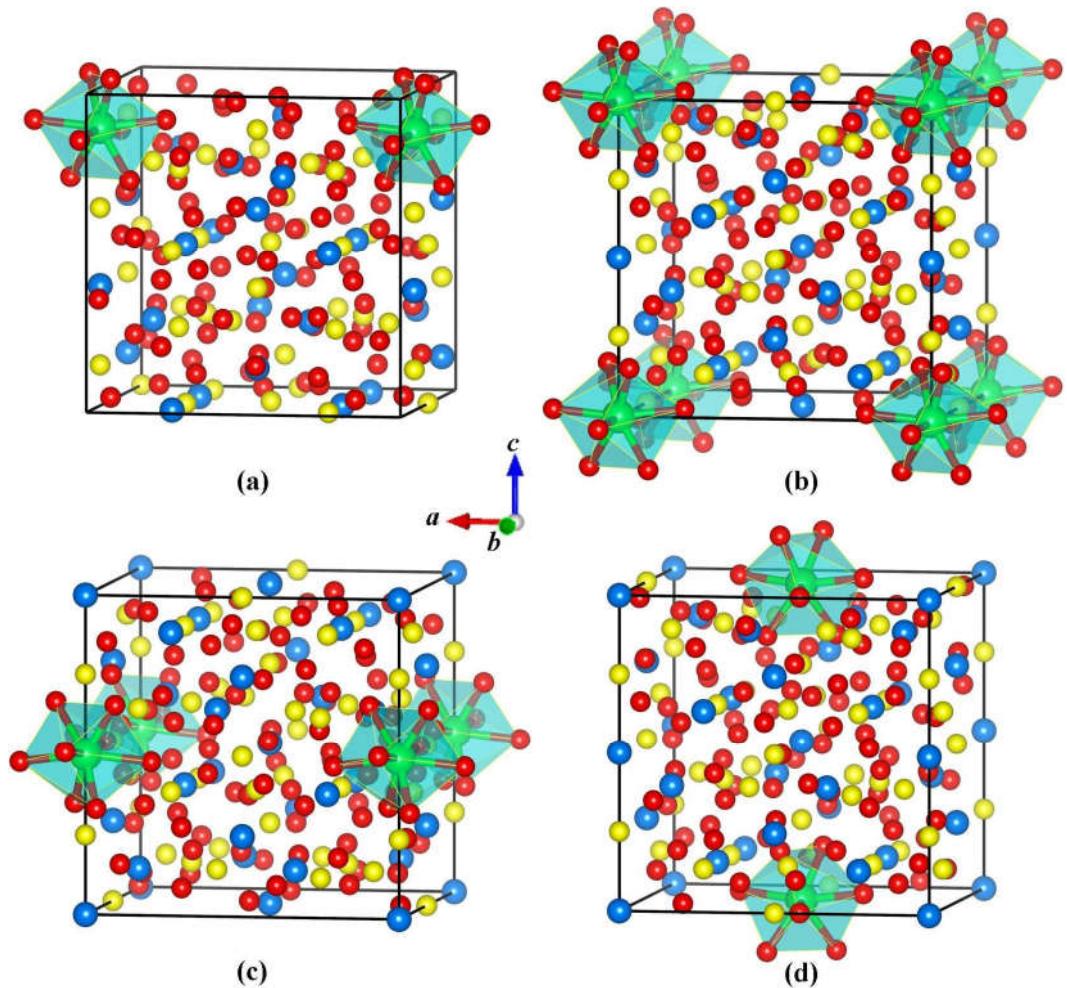
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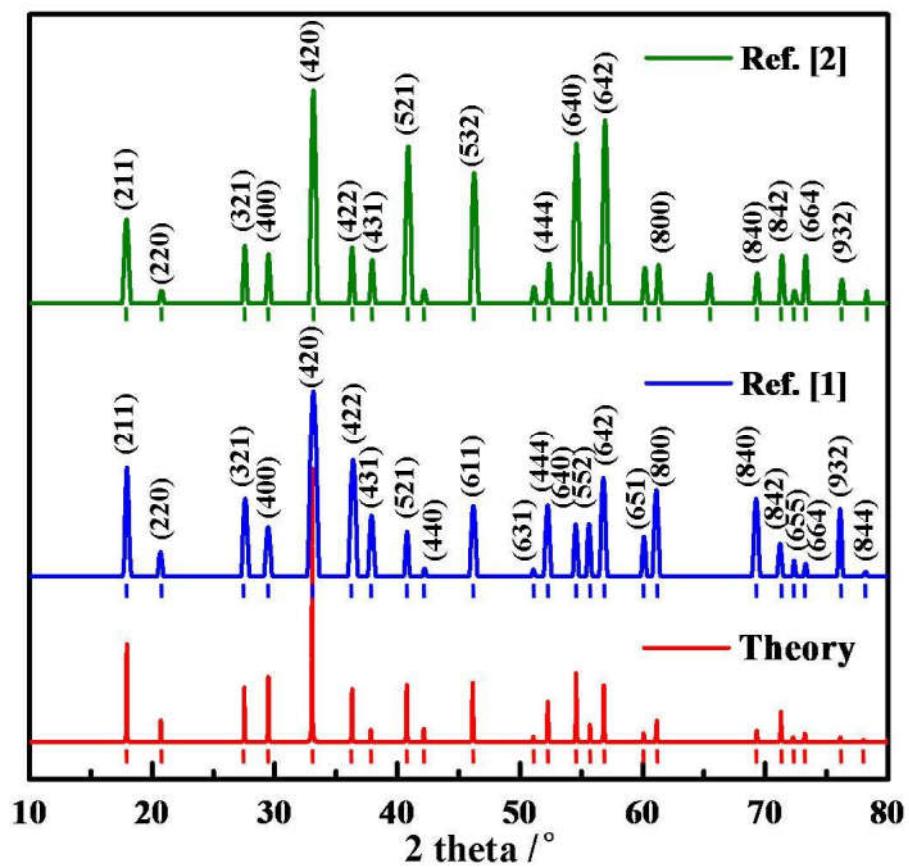
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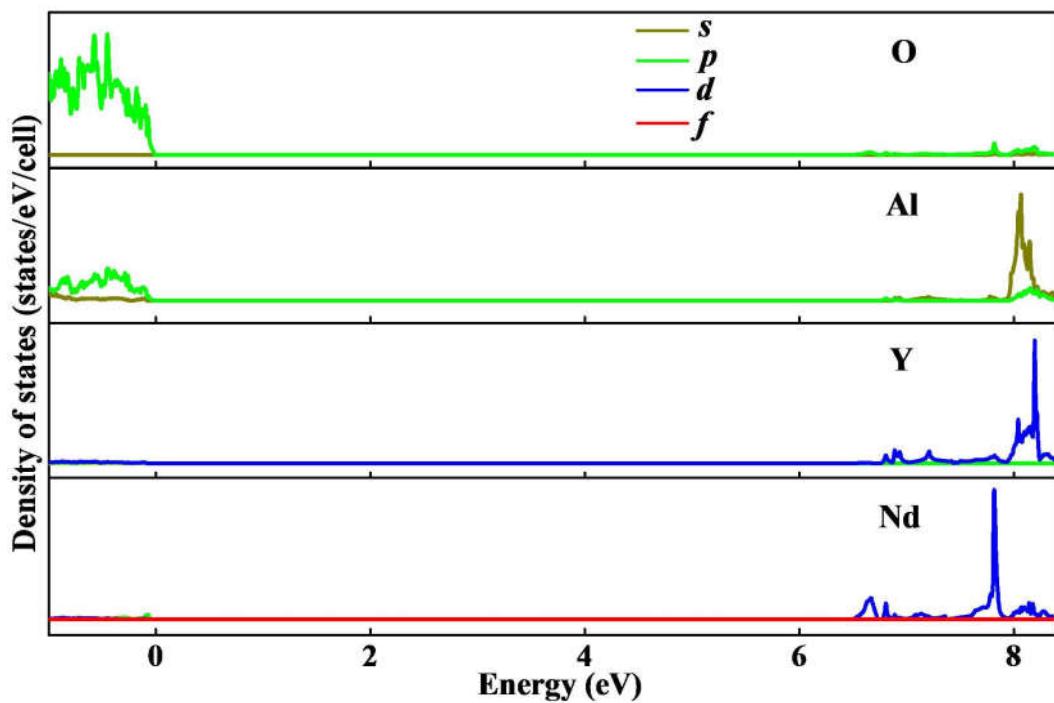
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**Figure S1.** Coordination structures the optimized metastable (a), (b), (c) and (d) for Nd:YAG. The red, yellow, blue and green spheres represent O, Al, Y and Nd atoms, respectively.



**Figure S2.** The simulated complete XRD patterns of Nd:YAG compared with experimental data.



**Figure S3.** The calculated partial density of states for Nd:YAG.

**Table S1.** Lattice constants  $a$ ,  $b$  and  $c$ , unit-cell volume, relative energies for the ground state and metastable Nd:YAG crystals.

	Space group	$a$ (Å)	$b$ (Å)	$c$ (Å)	$V$ (Å $^3$ )	$\Delta E$ ( $\times 10^{-2}$ meV)
Nd:YAG	$C_{222}$	12.1145	12.1145	12.1145	1777.94	0
Isomer (a)	$P_1$	12.1146	12.1146	12.1220	1779.07	0.42
Isomer (b)	$C_{222}$	12.1145	12.1145	12.1222	1779.07	0.47
Isomer (c)	$C_{222}$	12.1146	12.1146	12.1223	1779.11	1.25
Isomer (d)	$P_1$	12.1145	12.1145	12.1222	1779.06	18.76

**Table S2.** The experimental and calculated Stark energy levels (all in cm<sup>-1</sup>) of Nd<sup>3+</sup> in YAG.

<sup>2S+1</sup> L <sub>J</sub>	State	Present work			Other <sup>[3]</sup>			Present work			Other <sup>[3]</sup>		
		E <sub>obs</sub> <sup>[3]</sup>	E <sub>calc</sub>	ΔE	E <sub>calc</sub>	ΔE	<sup>2S+1</sup> L <sub>J</sub>	State	E <sub>obs</sub> <sup>[3]</sup>	E <sub>calc</sub>	ΔE	E <sub>calc</sub>	ΔE
<sup>4</sup> I <sub>9/2</sub>	1	0	-12	-12	-4	-4		87	—	21505	—	21491	—
	2	130	130	0	142	12		88	21522	21524	2	21514	-8
	3	199	197	-2	205	6		89	21593	21612	19	21608	15
	4	308	311	3	328	20		90	21661	21674	13	21664	3
	5	857	850	-7	870	13		91	21697	21700	3	21709	12
<sup>4</sup> I <sub>11/2</sub>	6	2002	2004	2	1990	-12		92	21767	21766	-1	21754	-13
	7	2029	2028	-1	2017	-12		93	21791	21781	-10	21782	-9
	8	2110	2109	-1	2102	-8		94	—	21844	—	21851	—
	9	2147	2148	1	2140	-7		95	21872	21863	-9	21866	-6
	10	2468	2465	-3	2465	-3		96	21906	21914	8	21906	0
<sup>4</sup> I <sub>13/2</sub>	11	2521	2523	2	2524	3		97	22036	22063	27	22043	7
	12	3922	3928	6	3908	-14	<sup>2</sup> P <sub>1/2</sub>	98	23155	23122	-33	23135	-20
	13	3930	3932	2	3916	-14	<sup>2</sup> D(1) <sub>5/2</sub>	99	23674	23654	-20	23685	11
	14	4032	4036	4	4026	-6		100	23764	23757	-7	23764	0
	15	4047	4054	7	4041	-6		101	23849	23848	-1	23838	-11
<sup>4</sup> I <sub>15/2</sub>	16	4435	4430	-5	4421	-14	<sup>2</sup> P <sub>3/2</sub>	102	25994	26003	9	26026	32
	17	4442	4441	-1	4441	-1		103	—	26078	—	26085	—
	18	4498	4506	8	4504	6	<sup>4</sup> D <sub>3/2</sub>	104	27571	27562	-9	27550	-21
	19	5758	5757	-1	5756	-2	and	105	27670	27677	7	27683	13
	20	5814	5815	1	5804	-10	<sup>4</sup> D <sub>5/2</sub>	106	27809	27819	10	27818	9
	21	5936	5942	6	5939	3		107	28183	28192	9	28182	-1
	22	5970	5966	-4	5971	1		108	28263	28268	5	28272	9
	23	6570	6555	-15	6557	-13	<sup>4</sup> D <sub>1/2</sub>	109	28359	28355	-4	28374	15

	24	6583	6585	2	6596	13	$^2\text{I}_{11/2}$	110	28580	28611	31	28601	21
	25	6639	6641	2	6659	20		111	28800	28822	22	28819	19
	26	6734	6737	3	6758	24		112	28930	28956	26	28957	27
$^4\text{F}_{3/2}$	27	11427	11426	-1	11421	-6		113	29140	29139	-1	29124	-16
	28	11512	11500	-12	11493	-19		114	29270	29287	17	29258	-12
$^4\text{F}_{5/2}$	29	12370	12354	-16	12359	-11		115	29360	29382	22	29366	6
$^4\text{F}_{5/2}$ and $^2\text{H}(2)_{9/2}$	30	12432	12437	5	12436	4	$^2\text{L}_{15/2},$ $^4\text{D}_{7/2},$ and	116	29715	29705	-10	29712	-3
	31	12519	12500	-19	12451	-68	$^2\text{I}_{13/2}$	117	29876	29865	-11	29867	-9
	32	12575	12590	15	12610	35		118	29880	29881	1	29881	1
	33	12607	12640	33	12650	43		119	29920	29905	-15	29917	-3
	34	12623	12646	23	12693	70		120	29953	29953	0	29952	-1
	35	12819	12827	8	12811	-8		121	—	30054	—	30048	—
	36	12840	12844	4	12856	16		122	30070	30064	-6	30057	-13
$^4\text{F}_{7/2}$	37	13363	13355	-8	13364	1		123	30140	30135	-5	30121	-19
$^4\text{F}_{7/2}$ and $^4\text{S}_{3/2}$	38	13433	13426	-7	13440	7		124	30160	30140	-20	30147	-13
	39	13563	13555	-8	13555	-8		125	30190	30193	3	30198	8
	40	13572	13574	2	13572	0		126	30230	30234	4	30233	3
	41	13596	13591	-5	13596	0		127	30289	30270	-19	30283	-6
	42	13633	13632	-1	13646	13		128	30360	30340	-20	30350	-10
$^4\text{F}_{9/2}$	43	14626	14634	8	14643	17		129	30400	30387	-13	30389	-11
	44	14678	14675	-3	14685	7		130	30464	30453	-11	30442	-22
	45	14793	14788	-5	14779	-14		131	30495	30504	9	30495	0
	46	14819	14825	6	14815	-4		132	30547	30552	5	30550	3
	47	14916	14924	8	14914	-2		133	30611	30602	-9	30593	-18
$^2\text{H}(2)_{11/2}$	48	15838	15836	-2	15898	60		134	30620	30620	0	30611	-9
	49	15870	15900	30	15920	50	$^2\text{L}_{17/2}$	135	—	31237	—	31247	—
	50	—	15920	—	15946	—		136	—	31406	—	31393	—
	51	15957	15943	-14	15956	-1		137	31440	31441	1	31459	19
	52	16103	16079	-24	16037	-66		138	31530	31503	-27	31525	-5

	53	16119	16102	-17	16067	-52		139	31570	31573	3	31567	-3
<sup>4</sup> G <sub>5/2</sub>	54	16849	16852	3	16848	-1		140	31585	31598	13	31583	-2
	55	16992	16979	-13	16978	-14		141	31665	31652	-13	31667	2
	56	17047	17057	10	17054	7		142	—	31776	—	31795	—
<sup>2</sup> G <sub>7/2</sub>	57	17241	17226	-15	17220	-21		143	—	31903	—	31885	—
	58	17268	17275	7	17275	7	<sup>2</sup> H(1) <sub>9/2</sub>	144	32613	32614	1	32621	8
	59	17322	17296	-26	17302	-20		145	32663	32663	0	32662	0
	60	17575	17565	-10	17594	19		146	32745	32725	-20	32730	-15
<sup>4</sup> G <sub>7/2</sub>	61	18723	18719	-4	18709	-14		147	32802	32805	3	32801	-1
	62	18822	18830	8	18831	9		148	32840	32822	-18	32835	-5
	63	18843	18855	12	18863	20	<sup>2</sup> D(2) <sub>3/2</sub>	149	32980	32970	10	32966	-14
	64	18986	18968	-18	18987	1		150	33045	33051	6	33056	11
<sup>2</sup> K <sub>13/2</sub> and	65	19154	19136	-18	19133	-21	<sup>2</sup> D(2) <sub>5/2</sub>	151	33693	33724	31	33699	6
<sup>4</sup> G <sub>9/2</sub>	66	—	19272	—	19262	—	and	152	—	33800	—	33783	—
	67	19294	19319	25	19317	23	<sup>2</sup> H(1) <sub>11/2</sub>	153	33840	33837	-3	33830	-10
	68	—	19429	—	19421	—		154	34050	34080	30	34071	21
	69	19470	19452	-18	19448	-22		155	34110	34122	12	34123	13
	70	—	19509	—	19499	—		156	34210	34179	-31	34170	-40
	71	19543	19547	4	19555	12		157	34260	34255	-5	34256	-4
	72	19596	19586	-10	19593	-3		158	34290	34327	37	34333	43
	73	19651	19666	15	19660	9		159	—	34457	—	34475	—
	74	19814	19833	19	19823	9	<sup>2</sup> F(2) <sub>5/2</sub>	160	37789	37763	-26	37799	10
	75	—	19890	—	19870	—		161	37900	37939	39	37920	20
	76	20048	20045	-3	20026	-22		162	38065	38051	-14	38043	-22
<sup>2</sup> G(1) <sub>9/2</sub> ,	77	20730	20716	-14	20719	-11	<sup>2</sup> F(2) <sub>7/2</sub>	163	39202	39208	6	39232	30
<sup>2</sup> D(1) <sub>3/2</sub> ,	78	20773	20781	8	20797	24		164	—	39255	—	39251	—
<sup>4</sup> G <sub>11/2</sub> , and	79	20790	20787	-3	20819	29		165	39330	39316	-14	39312	-18
<sup>2</sup> K <sub>15/2</sub>	80	20803	20802	-1	20833	30		166	—	39565	—	39553	—
	81	20962	20954	-8	20954	-8	<sup>2</sup> G(2) <sub>9/2</sub>	167	—	47000	—	47000	—

82	—	21020	—	21011	—	168	—	47018	—	47010	—
83	21029	21056	27	21057	28	169	47200	47191	-9	47170	-30
84	21080	21076	-4	21072	-8	170	—	47266	—	47247	—
85	21159	21144	15	21141	-18	171	—	47325	—	47266	—
86	21162	21150	-12	21152	-10						

**Table S3.** Calculated wavelengths ( $\lambda$ ), ED ( $A_{ED}$ ) and MD ( $A_{MD}$ ) radiative decay rates, branching ratios ( $\beta$ ) and radiative lifetimes ( $\tau$ ) for spontaneous emission transitions between the first 11 excited states in Nd:YAG. Available theoretical and experimental results are also listed for comparison.

Transition		$\lambda$ (nm)		$A_{total}$ ( $s^{-1}$ )		$\beta$		$\tau$ ( $\mu s$ )	
		Present	Other	$A_{ED}$ ( $s^{-1}$ )	$A_{MD}$ ( $s^{-1}$ )	Present	Other	Present	Other
$^4I_{11/2}$	$^4I_{9/2}$	5359		13.4	1.6	15.0		1	66541
$^4I_{13/2}$	$^4I_{9/2}$	2609		37.6	0	37.6		0.69	18429
	$^4I_{11/2}$	5085		14.4	2.2	16.6		0.31	
$^4I_{15/2}$	$^4I_{9/2}$	1707		11.5	0	11.5		0.18	15795
	$^4I_{11/2}$	2504		34.2	0	34.2		0.54	
	$^4I_{13/2}$	4933		16.0	1.6	17.6		0.28	
$^4F_{3/2}$	$^4I_{9/2}$	897	885 <sup>[4]</sup> 900 <sup>[5]</sup>	1254.2	0	1254.2	1420 <sup>[7]</sup>	0.34	0.32 <sup>[7]</sup>
	$^4I_{11/2}$	1077	1073 <sup>[6]</sup>	1962.0	0	1962.0	1940 <sup>[7]</sup>	0.54	0.54 <sup>[7]</sup>
	$^4I_{13/2}$	1367	1350 <sup>[7]</sup> 1340 <sup>[4]</sup>	415.8	0	415.8	473 <sup>[7]</sup>	0.11	0.14 <sup>[7]</sup>
	$^4I_{15/2}$	1891	1850 <sup>[5]</sup> 1800 <sup>[7]</sup>	20.9	0	20.9	15 <sup>[7]</sup>	0.01	0.003 <sup>[7]</sup>
$^4F_{5/2}$	$^4I_{9/2}$	822	809 <sup>[4]</sup> 808 <sup>[8]</sup>	3192.3	0	3192.3		0.66	207
	$^4I_{11/2}$	970		476.7	0	476.7		0.10	
	$^4I_{13/2}$	1200		976.2	0	976.2		0.20	
	$^4I_{15/2}$	1584		195.3	0	195.3		0.04	
$^2H(2)_{9/2}$	$^4I_{9/2}$	811		488.8	6.9	495.7		0.58	1177
	$^4I_{11/2}$	955		67.9	4.2	72.1		0.09	

	$^4I_{13/2}$	1176	167.7	0	167.7	0.20	
	$^4I_{15/2}$	1544	113.5	0	113.5	0.13	
$^4F_{7/2}$	$^4I_{9/2}$	762	2572.7	0.1	2572.8	0.50	193
	$^4I_{11/2}$	889	1562.3	0	1562.3	0.30	
	$^4I_{13/2}$	1077	366.2	0	366.2	0.07	
	$^4I_{15/2}$	1378	677.6	0	677.6	0.13	
$^4S_{3/2}$	$^4I_{9/2}$	755	2776.3	0	2776.3	0.43	156
	$^4I_{11/2}$	879	1566.3	0	1566.3	0.24	
	$^4I_{13/2}$	1062	1389.1	0	1389.1	0.22	
	$^4I_{15/2}$	1354	674.4	0	674.4	0.11	
$^4F_{9/2}$	$^4I_{9/2}$	694	264.1	1.9	266.0	0.06	237
	$^4I_{11/2}$	798	1539.3	1.7	1541.0	0.37	
	$^4I_{13/2}$	946	1483.9	0	1483.9	0.35	
	$^4I_{15/2}$	1171	912.9	0	912.9	0.22	
$^2H(2)_{11/2}$	$^4I_{9/2}$	638	79.0	0	79.0	0.28	3592
	$^4I_{11/2}$	725	54.6	2.6	57.2	0.21	
	$^4I_{13/2}$	845	35.7	5.7	41.4	0.15	
	$^4I_{15/2}$	1020	73.0	0	73.0	0.26	
	$^4F_{3/2}$	2215	1.4	0	1.4	0.01	
	$^4F_{5/2}$	2865	2.0	0	2.0	0.01	
	$^2H(2)_{9/2}$	3007	17.6	1.9	19.5	0.07	
	$^4F_{7/2}$	3930	3.9	0	3.9	0.01	
$^4G_{5/2}$	$^4I_{9/2}$	597	4651.8	0	4651.8	0.58	124
	$^4I_{11/2}$	672	2783.4	0	2783.4	0.35	
	$^4I_{13/2}$	775	488.8	0	488.8	0.06	
	$^4I_{15/2}$	919	20.7	0	20.7	0.01	

**Table S4.** The calculated statevectors of the ground state and the first 11 excited states for Nd<sup>3+</sup> in YAG.

State $^{2s+1}L_J$	Statevector $ ^{2s+1}L_J\rangle$
$^4I_{9/2}$	0.984 $ ^4I_{9/2}\rangle - 0.168  ^2H(2)_{9/2}\rangle + 0.057  ^2H(1)_{9/2}\rangle$
$^4I_{11/2}$	-0.995 $ ^4I_{11/2}\rangle + 0.096  ^2H(2)_{11/2}\rangle - 0.036  ^2H(1)_{11/2}\rangle$
$^4I_{13/2}$	-0.998 $ ^4I_{13/2}\rangle - 0.064  ^2K_{13/2}\rangle + 0.024  ^2I_{13/2}\rangle$
$^4I_{15/2}$	-0.993 $ ^4I_{15/2}\rangle - 0.115  ^2K_{15/2}\rangle + 0.009  ^2L_{15/2}\rangle$
$^4F_{3/2}$	-0.970 $ ^4F_{3/2}\rangle - 0.224  ^2D(1)_{3/2}\rangle + 0.062  ^2D(2)_{3/2}\rangle$
$^4F_{5/2}$	-0.988 $ ^4F_{5/2}\rangle - 0.147  ^2D(1)_{5/2}\rangle + 0.033  ^2F(2)_{5/2}\rangle$
$^2H(2)_{9/2}$	-0.731 $ ^2H(2)_{9/2}\rangle + 0.385  ^4F_{9/2}\rangle - 0.340  ^2G(1)_{9/2}\rangle$
$^4F_{7/2}$	-0.965 $ ^4F_{7/2}\rangle + 0.199  ^2G(1)_{7/2}\rangle - 0.158  ^2G(2)_{7/2}\rangle$
$^4S_{3/2}$	-0.971 $ ^4S_{3/2}\rangle - 0.222  ^2P_{3/2}\rangle - 0.073  ^4F_{3/2}\rangle$
$^4F_{9/2}$	0.862 $ ^4F_{9/2}\rangle + 0.454  ^2H(2)_{9/2}\rangle - 0.156  ^2H(1)_{9/2}\rangle$
$^2H(2)_{11/2}$	0.894 $ ^2H(2)_{11/2}\rangle - 0.354  ^2H(1)_{11/2}\rangle - 0.251  ^4G_{11/2}\rangle$
$^4G_{5/2}$	-0.993 $ ^4G_{5/2}\rangle - 0.083  ^2F(1)_{5/2}\rangle - 0.082  ^2F(2)_{5/2}\rangle$

## Appendix – Method and equations for the calculations of transition intensities

The model Hamiltonian for Nd<sup>3+</sup> is defined as<sup>[10]-[11]</sup>

$$\begin{aligned}
 H_f = & E_{AVE} + \sum_{k=2,4,6} F^k f_k + \zeta_{4f} \times \sum_i \vec{l}_i \cdot \vec{s}_i + \alpha L(L+1) + \beta G(G_2) + \gamma G(R_7) \\
 & + \sum_{i=2,3,4,6,7,8} T^i t_i + \sum_{j=0,2,4} M^j m_j + \sum_{k=2,4,6} P^k p_k
 \end{aligned} \tag{A1}$$

where  $E_{AVE}$  represents the barycenter energy of the 4f<sup>3</sup> configuration. The next seven terms represent the Coulomb repulsion, spin-orbit, two-body, three-body, spin-other-orbit and electrostatically correlated spin-orbit interactions. Moreover,  $F^k$  and  $\zeta_{4f}$  are the radial parts of the electrostatic and spin-orbit coupling constant. Two-body and Judd's three-body parameters are represented by  $\alpha$ ,  $\beta$ ,  $\gamma$  and  $T^i$ .  $G(G_2)$  and  $G(R_7)$  represent the eigenvalues of the Casimir's operators for the Lie groups  $G_2$  and  $R_7$ . The remaining parameters,  $M^j$  and  $P^k$ , are used to represent the Marvin integrals and spin-orbit perturbations.

The crystal field interaction  $H_{CF}$  for Nd<sup>3+</sup> in YAG, in the form of Wybourne normalization, can be expressed as<sup>[12]</sup>

$$\begin{aligned}
 H_{CF} = & B_2^0 C_2^0 + B_2^2 (C_2^2 + C_2^{-2}) + B_4^0 C_4^0 + B_4^2 (C_4^2 + C_4^{-2}) + B_4^4 (C_4^4 + C_4^{-4}) \\
 & + B_6^0 C_6^0 + B_6^2 (C_6^2 + C_6^{-2}) + B_6^4 (C_6^4 + C_6^{-4}) + B_6^6 (C_6^6 + C_6^{-6})
 \end{aligned} \tag{A2}$$

where  $C_q^k$  are the normalized spherical-tensor operators and  $B_q^k$  are the crystal field parameters (CFPs). The values of these CFPs can be determined by the least-squares fit to the observed energy levels.<sup>[12]</sup>

The ED ( $A_{ED}$ ) radiative decay rates can be written as<sup>[13]-[14]</sup>

$$A_{ED(SLJ \rightarrow S'L'J')} = \frac{16\pi^3 e^2}{3\varepsilon_0 h c^3} \frac{\nu^3}{(2J+1)} \chi_{ED} \sum_{\lambda=2,4,6} \Omega_{(\lambda)} \left| \langle l^N SLJ | U^{(\lambda)} | l^N S'L'J' \rangle \right|^2 \quad (\text{A3})$$

where  $\nu$  is the transition frequency,  $n$  is the refractive index and  $\chi_{ED}$  is the local-field correction for ED induced transitions with the form of  $(n^2+1)^2/(9n)$  and  $n(n^2+1)^2/9$  for absorption and emission transition, respectively. The Judd-Ofelt intensity parameters  $\Omega_{(\lambda)}$  should be summed over  $\lambda=2,4,6$  for a product with the even-rank reduced matrix elements of the  $U^{(\lambda)}$  tensor operator.

The MD ( $A_{MD}$ ) radiative decay rates can be written as<sup>[13]-[14]</sup>

$$A_{MD} = \frac{\pi h e^2}{3\varepsilon_0 c^5 m_e^2} \frac{\nu^3}{g} \chi_{MD} \left| \langle l^N \psi | L + g_e S | l^N \psi' \rangle \right|^2 \quad (\text{A4})$$

where  $g_e = 2.00232$  is the gyromagnetic ratio of the electron and  $g$  is the degeneracy of the initial level.  $\chi_{MD}$  is the local-field correction for MD induced transitions with the form of  $n$  and  $n^3$  for the absorption and emission transition, respectively.  $\psi$  and  $\psi'$  are the statevectors for the initial and terminating levels for the  $\psi \rightarrow \psi'$  transition, respectively. For transitions between J-multiplets, the statevector takes the form of  $\psi(SLJ)$  with  $g = (2J+1)$  while for transitions between crystal field levels, it takes the form of  $\psi(SLJ\Gamma_i)$ .

The radiative lifetime can be written as<sup>[13]-[14]</sup>

$$\tau_{SLJ} = \frac{1}{\sum_{S'L'J'} (A_{ED(SLJ \rightarrow S'L'J')} + A_{MD(SLJ \rightarrow S'L'J')})} \quad (\text{A5})$$

The branching ratio can be written as<sup>[13]-[14]</sup>

$$\beta_{(SLJ \rightarrow S'L'J')} = \tau_{SLJ} \times [A_{ED(SLJ \rightarrow S'L'J')} + A_{MD(SLJ \rightarrow S'L'J')}] \quad (\text{A6})$$

The MD absorption oscillator strengths can be written as<sup>[13]-[14]</sup>

$$P_{MD} = \frac{h\nu}{6m_e c^2} \frac{n}{(2J+1)} \left| \langle l^N SLJ | L + g_e S | l^N SLJ' \rangle \right|^2 \quad (\text{A7})$$

## Coordinates of all atoms for the ground state Nd:YAG

Atom	<i>x</i>	<i>y</i>	<i>z</i>
Nd	0.50000	0.50000	0.50000
Y1	0.12479	0.24980	-0.12460
Y2	0.37452	0.24891	-0.62590
Y9	0.00000	0.00000	0.00000
Y10	0.00000	-0.50000	-0.24970
Y12	0.87504	0.25006	-0.62507
Y13	0.37494	-0.24993	-0.12472
Y20	0.00000	-0.50000	-0.75044
Y22	-0.00000	-0.00000	-0.50000
Y23	0.50000	-0.50000	-0.00000
Al1	-0.00008	0.24983	0.12499
Al2	0.50048	-0.24854	-0.37389
Al5	0.00013	-0.24976	-0.37489
Al6	0.49981	0.24969	0.12438
Al17	0.37480	0.24966	-0.12514
Al18	0.12480	0.24963	-0.62515
Al21	0.24996	-0.12498	0.12486
Al23	0.24890	-0.37497	-0.37449
Al25	0.00000	-0.00000	-0.24993
Al26	0.00000	-0.50000	0.00014
Al27	0.50000	-0.00000	-0.49984
Al28	0.50000	-0.50000	-0.75245
O1	0.96966	0.19943	-0.02409
O2	0.47044	-0.29614	-0.52395
O5	0.03024	-0.30032	-0.22569
O6	0.53015	0.19951	-0.72657
O9	0.14896	-0.72004	0.07461

O10	0.64961	-0.21972	-0.42473
O11	0.64957	-0.27943	-0.82532
O12	0.14872	0.21957	-0.32453
O17	0.05031	0.10071	-0.84490
O18	0.55170	-0.39698	-0.33916
O19	0.94980	-0.39918	-0.40513
O20	0.44963	0.10066	0.09463
O49	0.03008	-0.70014	-0.72598
O50	0.53005	-0.19820	-0.22522
O53	0.96986	-0.19956	-0.52415
O54	0.46963	-0.69995	-0.02458
O57	0.85068	0.21985	-0.82461
O58	0.35080	-0.28069	-0.32493
O59	0.35077	-0.21964	0.07426
O60	0.85114	-0.72018	-0.42514
O65	0.94982	-0.60105	0.09467
O66	0.44932	-0.10020	-0.40455
O67	0.05040	-0.10069	-0.34477
O68	0.55039	-0.60136	-0.84615

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