Supporting Information

Multiple Intersystem Crossing Processes in Ge-doped Silica Glass: Emission Mechanism of Twofold Coordinated Ge atoms

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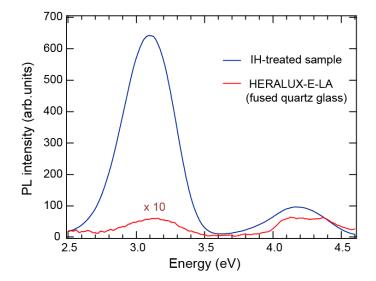


Figure S1. Room-temperature PL spectra of a Ge-doped silica sample prepared by IH method and a commercial fused quartz sample (HERALUX-E-LA) measured under $S_0 \rightarrow S_1$ excitation at 5 eV.

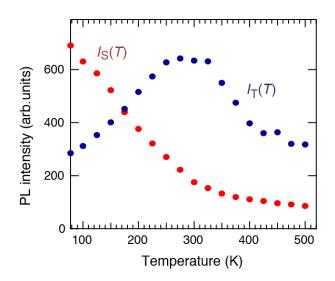


Figure S2. Temperature dependence of the PL signals of a commercial fused quartz sample measured under 5-eV excitation. PL intensities of the $S_1 \rightarrow S_0(I_S(T))$ and the $T_1 \rightarrow S_0(I_T(T))$ emission bands peaking at ~4.2 and ~3.1 eV, respectively, are shown as a function of temperature *T*.

Detailed derivation of Eqs. (5)-(8)

According to the emission scheme shown in Fig. 1(a), the kinetic equations for the respective energy levels can be written as follows:

$$\frac{dN_0}{dt} = -\alpha N_0 + k_{\rm S}^{\rm R} N_1 + \left(k_{\rm T}^{\rm R} + k_{\rm T}^{\rm NR}\right) N_2 \tag{S1}$$

$$\frac{dN_1}{dt} = \alpha N_0 - (k_{\rm S}^{\rm R} + k_{\rm ISC})N_1 \tag{S2}$$

$$\frac{dN_2}{dt} = k_{\rm ISC} N_1 - (k_{\rm T}^{\rm R} + k_{\rm T}^{\rm NR}) N_2$$
(S3)

where N_0 , N_1 , and N_2 are the population of the S₀, S₁ and T_n states, respectively,

When we treat the above equations under the steady state approximation $dN_0/dt=dN_1/dt=dN_2/dt=0$, the quantum efficiency η of the $S_1 \rightarrow S_0(\eta_S)$ and $T_n \rightarrow S_0(\eta_T)$ emission can be represented by

$$\eta_{\rm S} = \frac{k_{\rm S}^{\rm R} N_1}{\alpha N_0} = \frac{k_{\rm S}^{\rm R}}{k_{\rm S}^{\rm R} + k_{\rm ISC}} \tag{S4}$$

$$\eta_{\rm T} = \frac{k_1 N_2}{\alpha N_0} = \frac{k_{\rm T}^{\rm R} k_{\rm ISC}}{(k_{\rm T}^{\rm R} + k_{\rm T}^{\rm NR})(k_{\rm S}^{\rm R} + k_{\rm ISC})} \qquad .$$
(S5)

Under normal excitation conditions, the emission intensity is linearly correlated with η . Accordingly, the resulting emission intensity of the $S_1 \rightarrow S_0$ (I_S) and $T_n \rightarrow S_0$ (I_T) processes are

$$I_{\rm S} = K_1 \frac{k_{\rm S}^{\rm R}}{k_{\rm S}^{\rm R} + k_{\rm ISC}} , \qquad (S6) \text{ or } (5)$$
$$I_{\rm T} = K_2 \frac{k_{\rm T}^{\rm R} k_{\rm ISC}}{(k_{\rm T}^{\rm R} + k_{\rm T}^{\rm NR})(k_{\rm S}^{\rm R} + k_{\rm ISC})} \qquad (S7) \text{ or } (6)$$

where K_1 and K_2 are the scaling (temperature independent) constants depending on a given excitation intensity and the instrumental conditions. Using the functional form of $k_{\rm ISC}$ [see Eq. (4) in the text], Eq. (S6) can be modified into the following form:

$$I_{\rm S}(T) = K_1 \frac{k_{\rm S}^{\rm R}}{(k_{\rm S}^{\rm R} + k_{\rm ISC}^{(1)} + A_{\rm ISC}^{(2)} \exp(-E_{\rm ISC}^{(2)}/k_B T) + A_{\rm ISC}^{(3)} \exp(-E_{\rm ISC}^{(3)}/k_B T))}$$
$$= K_1 \frac{1}{(1 + (k_{\rm ISC}^{(1)}/k_{\rm S}^{\rm R}) + (A_{\rm ISC}^{(2)}/k_{\rm S}^{\rm R})\exp(-E_{\rm ISC}^{(2)}/k_B T) + (A_{\rm ISC}^{(3)}/k_{\rm S}^{\rm R})\exp(-E_{\rm ISC}^{(3)}/k_B T))}$$

$$= \frac{C_1}{1 + C_2 \exp\left(-\frac{E_{\rm ISC}^{(2)}}{k_B T}\right) + C_3 \exp\left(-\frac{E_{\rm ISC}^{(3)}}{k_B T}\right)} \quad , \qquad (S8) \text{ or } (7)$$

where

$$\begin{split} C_{1} &= \frac{K_{1}k_{\rm S}^{\rm R}}{k_{\rm S}^{\rm R} + k_{\rm ISC}^{(1)}} \ , \\ C_{2} &= \frac{A_{\rm ISC}^{(2)}}{k_{\rm S}^{\rm R} + k_{\rm ISC}^{(1)}} \ , \\ C_{3} &= \frac{A_{\rm ISC}^{(3)}}{k_{\rm S}^{\rm R} + k_{\rm ISC}^{(1)}}, \end{split}$$

Similarly, Eq. (S7) can be modified into the following form using the functional form of $k_{\rm T}^{\rm NR}$ [see Eq. (3) in the text] and $k_{\rm ISC}$ [see Eq. (4) in the text]:

$$I_{\rm T}(T) = K_2 \frac{k_{\rm T}^{\rm R}(k_{\rm ISC}^{(1)} + A_{\rm ISC}^{(2)} \exp(-E_{\rm ISC}^{(2)}/k_B T) + A_{\rm ISC}^{(3)} \exp(-E_{\rm ISC}^{(3)}/k_B T))}{\left(k_{\rm T}^{\rm R} + A_{\rm T}^{\rm NR1} \exp(-\frac{E_{\rm T}^{\rm NR1}}{k_B T}) + A_{\rm T}^{\rm NR2} \exp(-E_{\rm T}^{\rm NR2}/k_B T)\right) \left(k_{\rm S}^{\rm R} + k_{\rm ISC}^{(1)} + A_{\rm ISC}^{(2)} \exp(-E_{\rm ISC}^{(2)}/k_B T) + A_{\rm ISC}^{(3)} \exp(-E_{\rm ISC}^{(3)}/k_B T)\right)}$$

$$=\frac{C_4(C_5+C_2\exp\left(-\frac{E_{\rm ISC}^{(2)}}{k_BT}\right)+C_3\exp\left(-\frac{E_{\rm ISC}^{(3)}}{k_BT}\right))}{(1+C_6\exp\left(-\frac{E_{\rm T}^{\rm NR1}}{k_BT}\right)+C_7\exp\left(-\frac{E_{\rm T}^{\rm NR2}}{k_BT}\right))(1+C_2\exp\left(-\frac{E_{\rm ISC}^{(2)}}{k_BT}\right)+C_3\exp\left(-\frac{E_{\rm ISC}^{(3)}}{k_BT}\right))},$$
 (S9) or (8)

where

$$C_4=K_2,$$

$$C_{5} = \frac{k_{\rm ISC}^{(1)}}{k_{\rm S}^{\rm R} + k_{\rm ISC}^{(1)}},$$
$$C_{6} = \frac{A_{\rm T}^{\rm NR_{1}}}{k_{\rm T}^{\rm R}},$$
$$C_{7} = \frac{A_{\rm T}^{\rm NR_{2}}}{k_{\rm T}^{\rm R}}.$$

From the fits of the data shown in Fig. 2(a) to Eqs. (7) and (8), we can obtain the values of activation energies, i.e., $E_{\rm T}^{\rm NR1}$, $E_{\rm T}^{\rm NR2}$, $E_{\rm ISC}^{(2)}$, $E_{\rm ISC}^{(3)}$, along with the parameters related to the frequency factors, i.e., C_i (*i*=1–7). The fitted values are shown in Fig. 2(a). It should be noted, however, that the values of C_i (*i*=1–7) do not represent the absolute values of the frequency factors but are the relative values among these frequency factors.