

SUPPORTING INFORMATION

Super-Compressible Coaxial Carbon Nanotube@Graphene Arrays with Invariant Viscoelasticity over –100 to 500 °C in Ambient Air

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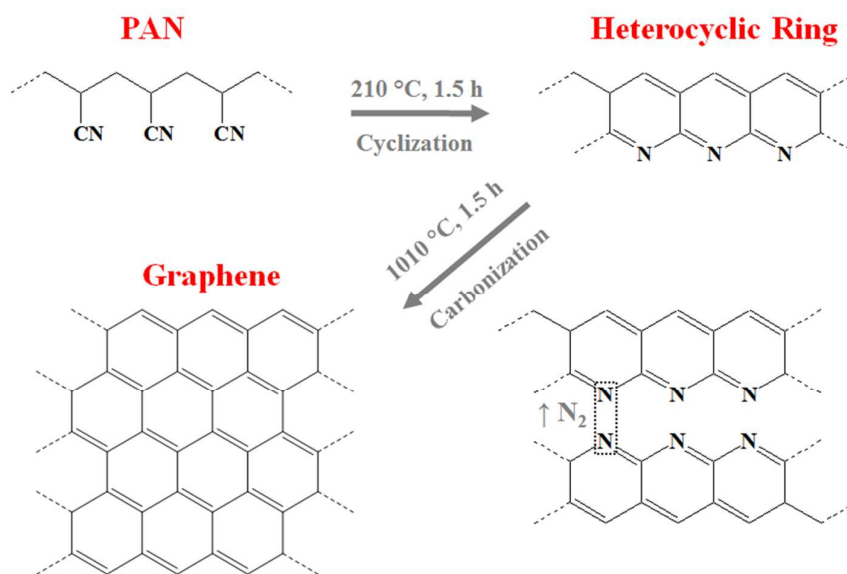
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Scheme S1 Schematic illustration of the synthesis of graphene layers by means of two-step pyrolysis processes.¹

Synthesis of CNT@Gr arrays. Typically, the PAN coated CNT arrays were first heated up to 210 °C and kept at this temperature for 1.5 h in Ar atmosphere, during which the PAN layers can be converted into condensed heterocyclic ring structure (cyclization).^{1,2} Then the resulting samples were further pyrolyzed at 1010 °C for 1.5 h to transform the heterocyclic ring structure into planar graphene layers by splitting off the heterocyclic nitrogen (carbonization).^{1,2} Finally, the conformal graphene layers can be successfully introduced onto the CNTs to form the coaxial CNT@Gr arrays.

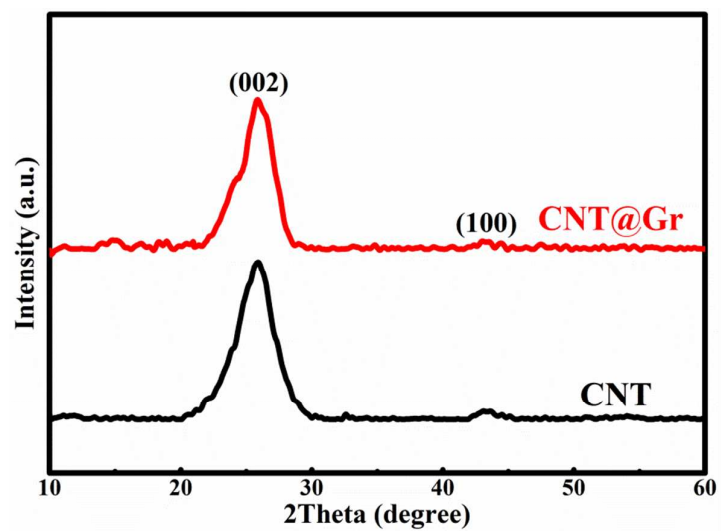


Figure S1 XRD profiles of CNT and CNT@Gr. The characteristic (002) and (100) diffraction peaks indicate that the crystallinity is well preserved for the CNT@Gr.

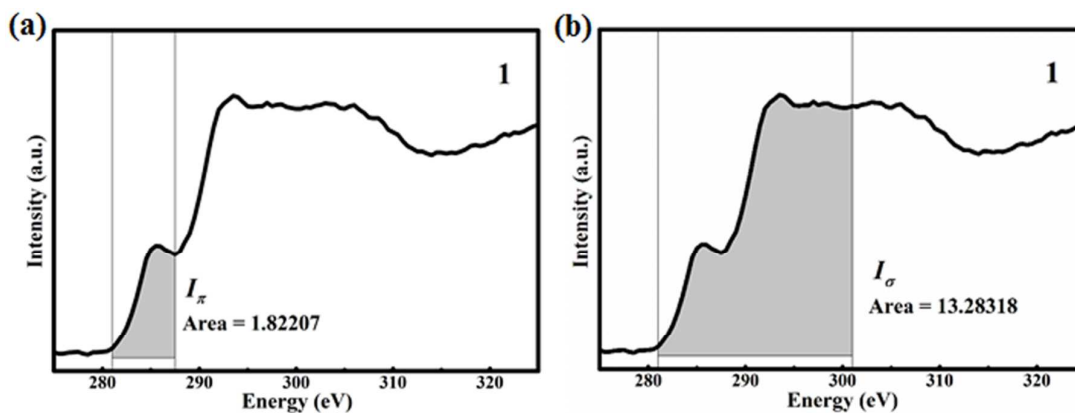


Figure S2 1s to π^* (I_π , shadow area in (a)) and total (I_σ , shadow area in (b)) peak intensities at the C K-edge for point 1.

Calculation of sp^2 / sp^3 ratio. Take the EELS spectrum of point 1 as an example, the 1s to π^* peak intensity at the C K-edge (background was subtracted prior acquiring the spectrum) was measured using an energy window of 5 eV (I_π , as shown in Figure S2a).³ The total C K-edge intensity was then measured in a window of up to 20 eV extending from the edge onset and covering the π^* peak and majority of the 1s to σ^* peak (I_σ , as shown in Figure S2b).³ Next, the sp^2 fraction x was calculated out according to $3x / (4 - x) = (I_\pi / I_\sigma) / (I_\pi^R / I_\sigma^R)$,^{3,4} where I_π^R and I_σ^R are the 1s to π^* and total peak intensities of the 100% sp^2 carbon reference material, respectively. As a result, the sp^2 fractions for **points 1, 2 and 3** were extracted to be 0.75, 0.76 and 0.73, corresponding to **sp^2 / sp^3 ratios of 3, 3.17 and 2.70**, respectively.⁵

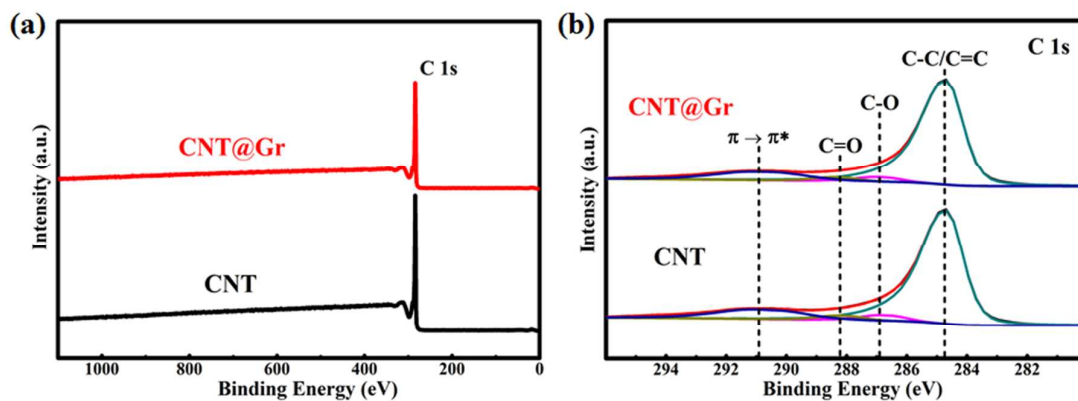


Figure S3 XPS survey and high resolution C 1s spectra of CNT and CNT@Gr, respectively.

XPS Analysis. XPS survey spectra show that no other elements except carbon (graphitic C 1s peak at 284.68 eV)⁶ can be detected for both the CNT and CNT@Gr. Meanwhile, high resolution C 1s spectrum of the CNT@Gr displays no noticeable shift or change in the peaks assigned to sp^2 and sp^3 hybridized states as compared to those of the initial CNT.⁷

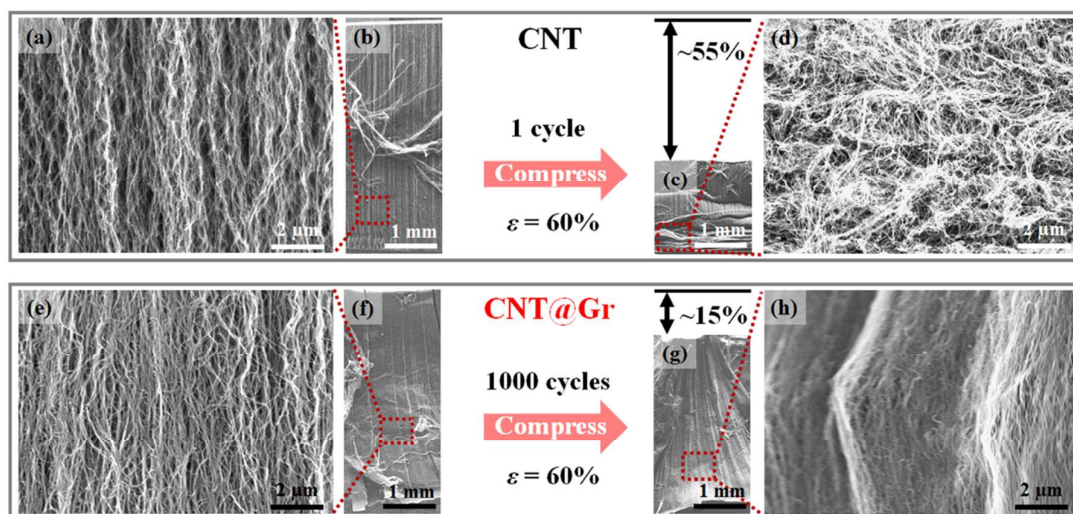


Figure S4 Cross-sectional SEM images of CNT (a–d) and CNT@Gr (e–h) arrays before and after 1 and 1000 compression cycles at strain (ϵ) = 60%, respectively. (a, d, e, h) are the zoom-in views of (b, c, f, g), respectively.

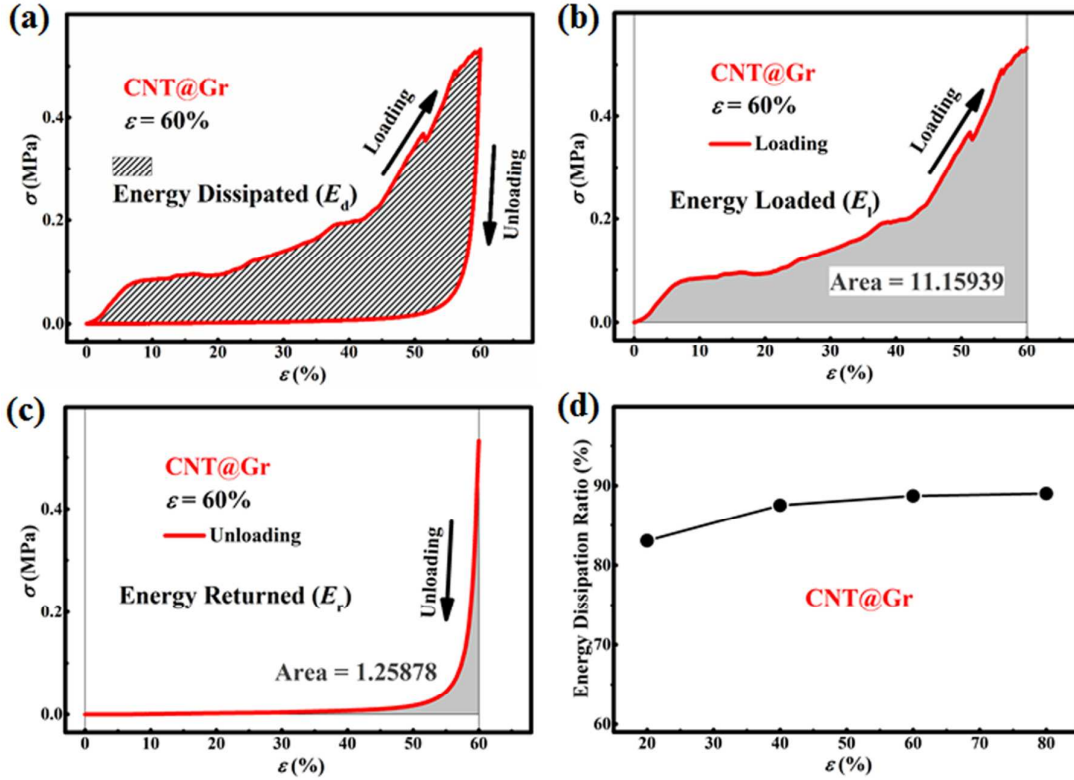


Figure S5 Energy dissipated (E_d , shadow area in (a)), energy loaded (E_l , shadow area in (b)) and energy returned (E_r , shadow area in (c)) for the first compression cycle at $\varepsilon = 60\%$. (d) Energy dissipation ratios of CNT@Gr arrays at various applied ε .

Calculation of energy dissipation ratio. Take the first compression cycle at $\varepsilon = 60\%$ for the CNT@Gr arrays as an example, the energy dissipation ratio (E_d / E_l) is extracted from the corresponding stress (σ) vs. ε curves,⁸ where E_d is the energy dissipated during the loading-unloading cycle (*i.e.* the area of the hysteresis loop, Figure S4a), E_l is the energy loaded during the compression loading process (*i.e.* the area under the loading σ vs. ε curve, which is integrated to be 11.15939, Figure S4b). Meanwhile, $E_d = E_l - E_r$, where E_r is the energy returned (*i.e.* the area under the unloading σ vs. ε curve, which is integrated to be 1.25878, Figure S4c). As a result, the energy dissipation ratio of the CNT@Gr arrays for the first compression cycle at $\varepsilon = 60\%$ could be calculated to be $E_d / E_l = (E_l - E_r) / E_l = (11.15939 - 1.25878) / 11.15939 \approx 88.72\%$.

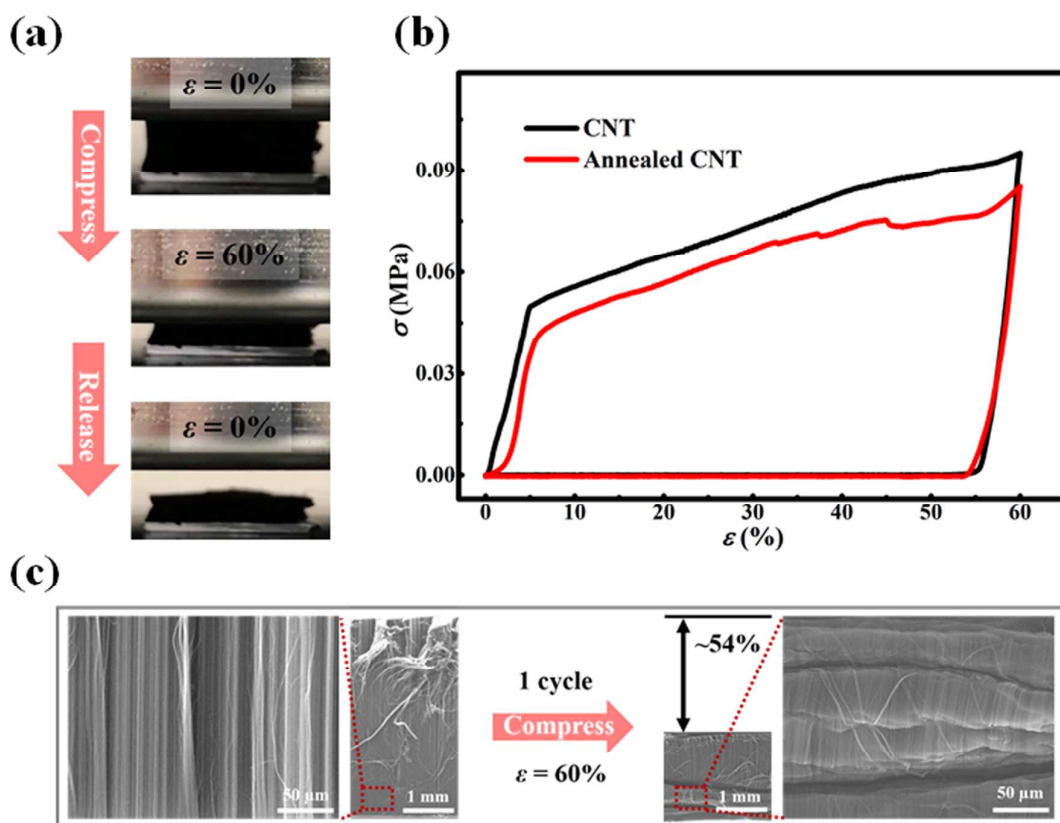


Figure S6 Compressive mechanical responses of annealed CNT arrays. (a) Annealed CNT arrays deform almost plastically upon uniaxial compression at $\varepsilon = 60\%$. (b) Compressive stress σ vs. ε curves at $\varepsilon = 60\%$ of CNT arrays before and after high temperature annealing. (c) Cross-sectional SEM images of the annealed CNT arrays before and after 1 cycle of compression at $\varepsilon = 60\%$.

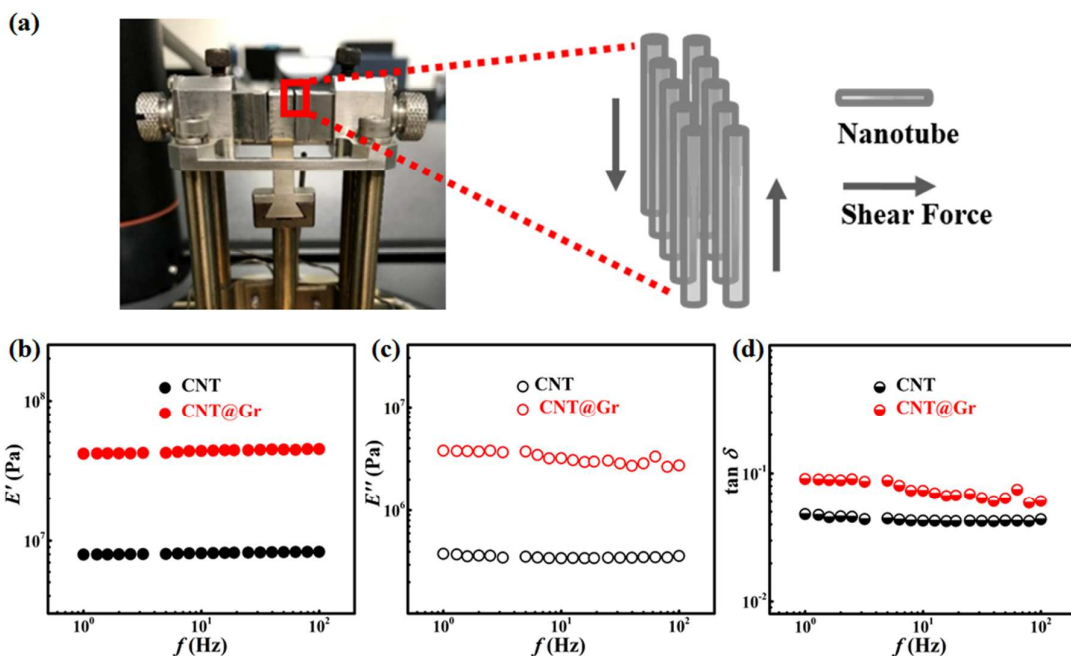


Figure S7 Dynamic shear tests of CNT and CNT@Gr arrays. (a) Shear sandwich clamp was used to measure the longitudinal shear viscoelastic properties of the NT arrays. (b) Storage modulus (E'), (c) loss modulus (E'') and (d) damping ratio ($\tan \delta$) of the CNT and CNT@Gr arrays as functions of applied frequency (f) (1–100 Hz, at $\varepsilon = 2\%$).

Dynamic shear tests. Upon the applied shear loading of $\varepsilon = 2\%$ over the broad f range, the CNT@Gr arrays perform generally stable E' and E'' of ~ 41.7 and ~ 3.8 MPa, respectively, which are ~ 5 – and ~ 10 –fold higher than those of their CNT counterpart (~ 8.0 and ~ 0.38 MPa, respectively). As a result, the CNT@Gr arrays shows ~ 2 –fold increase in the $\tan \delta$ (~ 0.1) as compared to the CNT arrays (~ 0.045). These substantial enhancements in the longitudinal shear viscoelastic properties of the CNT@Gr arrays indicate their reinforced inter-tube interactions because of the enlarged tube diameter and strengthened tube-tube contacts as compared to CNT arrays, which contribute to their enhanced compressive strength and viscoelasticity.

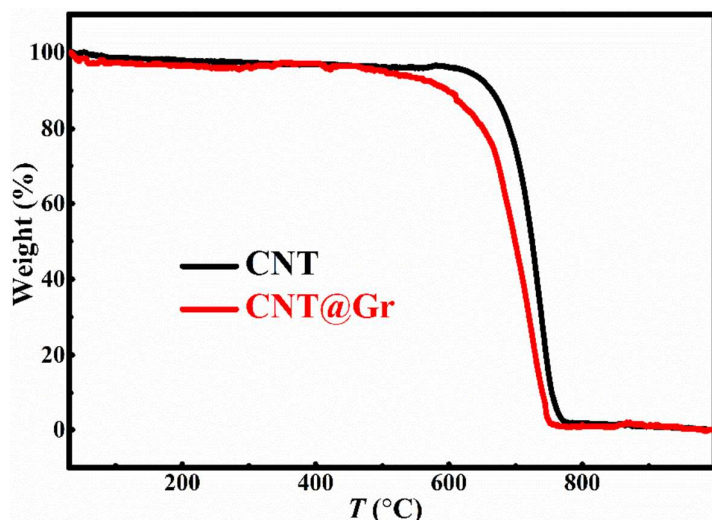


Figure S8 TGA analysis of CNT and CNT@Gr under a constant air flow of 50 mL min⁻¹. With the encapsulation of oxidation resistant graphene layers, the CNT@Gr generally retains the thermal stability of CNT in ambient air.

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