# Supporting Information for: Fabrication of 

 Supramolecular Chirality from Achiral Molecules at the Liquid/Liquid Interface Studied by Second Harmonic GenerationLu Lin, ${ }^{\dagger, \ddagger}$ Zhen Zhang, ${ }^{\ddagger}$ Yuan Guo, ${ }^{*, \ddagger, \pi}$ and Minghua Liu ${ }^{*, \ddagger}$<br>$\dagger$ National Center for Nanoscience and Technology, Beijing 100190, P. R. China $\ddagger$ Beijing National Laboratory for Molecular Sciences, Institute of Chemistry, Chinese Academy of Sciences, Beijing 100190, P. R. China<br>【University of Chinese Academy of Sciences, Beijing 100049, P. R. China

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## Fresnel Coefficients

In eq $1, a_{i}(i=1-7)$ denote the Fresnel coefficients, which can be expressed as:

$$
\begin{align*}
& a_{1}=L_{y y}(2 \omega) L_{z z}(\omega) L_{y y}(\omega) \sin \beta \\
& a_{2}=-2 L_{x x}(2 \omega) L_{z z}(\omega) L_{x x}(\omega) \cos ^{2} \beta \sin \beta \\
& a_{3}=L_{z z}(2 \omega) L_{x x}(\omega) L_{x x}(\omega) \cos ^{2} \beta \sin \beta \\
& a_{4}=L_{z z}(2 \omega) L_{z z}(\omega) L_{z z}(\omega) \sin ^{3} \beta  \tag{1}\\
& a_{5}=L_{z z}(2 \omega) L_{y y}(\omega) L_{y y}(\omega) \sin \beta \\
& a_{6}=2 L_{x x}(2 \omega) L_{z z}(\omega) L_{y y}(\omega) \cos \beta \sin \beta \\
& a_{7}=2 L_{y y}(2 \omega) L_{z z}(\omega) L_{x x}(\omega) \cos \beta \sin \beta
\end{align*}
$$

with

$$
\begin{align*}
L_{x x}\left(\omega_{i}\right) & =\frac{2 n_{1}\left(\omega_{i}\right) \cos \gamma}{n_{1}\left(\omega_{i}\right) \cos \gamma+n_{2}\left(\omega_{i}\right) \cos \beta} \\
L_{y y}\left(\omega_{i}\right) & =\frac{2 n_{1}\left(\omega_{i}\right) \cos \beta}{n_{1}\left(\omega_{i}\right) \cos \beta+n_{2}\left(\omega_{i}\right) \cos \gamma}  \tag{2}\\
L_{z z}\left(\omega_{i}\right) & =\frac{2 n_{2}\left(\omega_{i}\right) \cos \beta}{n_{1}\left(\omega_{i}\right) \cos \gamma+n_{2}\left(\omega_{i}\right) \cos \beta}\left(\frac{n_{1}\left(\omega_{i}\right)}{n^{\prime}\left(\omega_{i}\right)}\right)^{2}
\end{align*}
$$

In the above equations, $n_{1}$ and $n_{2}$ are the refractive indices of water and organic phase, $n^{\prime}$ is the refractive index of the interfacial layer, $\beta$ is the incident angle, and $\gamma$ is the refracted angle. Here we choose $n_{1}(\omega)=1.329, n_{1}(2 \omega)=1.339$ and $n_{2}(\omega)=n_{2}(2 \omega)=1.445$. According to Zhuang et al. (Phys. Rev. B 1999, 59, 12632), $n^{\prime}=n_{1} \sqrt{\frac{n_{2}^{2}\left(n_{2}^{2}+5\right)}{4 n_{2}^{2}+2}}$.

## Calculation of the Orientational Angle

The tensor elements $\chi_{z z z}, \chi_{z x x}, \chi_{x x z}$ and $\chi_{y x z}$ can be obtained by fitting the SHG curves. The fitting results are listed in Table S1.

We use $\chi_{z z z}, \chi_{z x x}$ and $\chi_{x x z}$ to calculate the orientational angle. By assuming $\beta_{c a a}=0$,

Table 1: Fitting results of the tensor elements.

| Concentration (mM) | $\chi_{z z z}$ | $\chi_{z x x}$ | $\chi_{x x z}$ | $\chi_{y x z}$ |
| :---: | ---: | ---: | ---: | ---: |
| 0.005 | 123.84 | 25.05 | 24.84 | 0.01 |
| 0.010 | 140.11 | 31.74 | 31.45 | 0.09 |
| 0.025 | 243.33 | 61.98 | 57.85 | 0.91 |
| 0.035 | 312.64 | 87.68 | 81.42 | 2.88 |
| 0.060 | 369.98 | 91.36 | 83.90 | 2.56 |
| 0.075 | 565.86 | 109.30 | 102.86 | 1.88 |
| 0.100 | 682.75 | 122.00 | 113.74 | 1.03 |
| 0.150 | 721.19 | 126.58 | 116.89 | 0.37 |
| 0.300 | 782.09 | 130.60 | 126.99 | 0.32 |

$\chi_{z x x}$ should be equal to $\chi_{x x z}$. As shown in Table S1, the values of $\chi_{z x x}$ are close to that of $\chi_{x x z}$. Therefore, it is reasonable to neglect $\beta_{c a a}$ for calculating the orientational angle. To minimize the error, we suppose:

$$
\begin{equation*}
\frac{\chi_{z z z}}{\chi_{z z z}+\chi_{z x x}+\chi_{x x z}} \approx \frac{\left\langle\cos ^{3} \theta\right\rangle}{\left\langle\cos ^{3} \theta\right\rangle+\left\langle\sin ^{2} \theta \cos \theta\right\rangle} \tag{3}
\end{equation*}
$$

Assuming $\theta$ has a very narrow distribution, we have:

$$
\begin{equation*}
\theta=\arccos \left(\sqrt{\frac{\chi_{z z z}}{\chi_{z z z}+\chi_{z x x}+\chi_{x x z}}}\right) \tag{4}
\end{equation*}
$$

We shall assess the error of the calculation. By assuming a narrow distribution of $\theta$, we have:

$$
\begin{equation*}
\frac{\chi_{z z z}}{\chi_{z z z}+\chi_{z x x}+\chi_{x x z}}=\frac{\cos ^{2} \theta+\frac{\beta_{c a a}}{\beta_{c c c}} \sin ^{2} \theta\left\langle\cos ^{2} \psi\right\rangle}{1+\frac{\beta_{c a a}}{2 \beta_{c c c}}} \tag{5}
\end{equation*}
$$

The ratio of $\beta_{c a a}$ and $\beta_{c c c}$ can be estimated as:

$$
\begin{equation*}
\frac{\beta_{c a a}}{\beta_{c c c}} \approx \frac{\chi_{z x x}-\chi_{x x z}}{\left(\chi_{z x x}+\chi_{x x z}\right) \sin ^{2} \theta} \leq 0.15 \tag{6}
\end{equation*}
$$

At lower concentrations, $\psi$ is randomly distributed. With increasing concentration, $\psi$ approaches zero. It is thus reasonable to suggest $\frac{1}{2} \leq\left\langle\cos ^{2} \psi\right\rangle \leq 1$. Let $\theta=33^{\circ}$, then the
calculated value for $\theta$ lies between $33.4^{\circ}$ and $34.7^{\circ}$. Therefore, the error of the calculated orientation is less than $2^{\circ}$.

## Langmuir Fitting of Adsorption

According to the Langmuir adsorption model, we have

$$
\begin{equation*}
\frac{\chi_{z z z}}{\cos ^{3} \theta} \approx N_{s} \beta_{c c c}=A \frac{K c}{1+K c} \tag{7}
\end{equation*}
$$

where $c$ is the bulk concentration, $A$ and $K$ are constants. The values of $A$ and $K$ can be obtained by fitting. Then the coverage can be calculated as $\frac{\chi_{z z z}}{A \cos ^{3} \theta}$.

